

Tracking simulations for cERL-FEL: Detailed comparison of SAD and GPT with theoretical explanations

Demin Zhou, Olga Tanaka
Accelerator laboratory, KEK

Acknowledgements:

N. Iida, R. Kato, T. Miyajima, N. Nakamura, K. Oide,
M. Shimada, Y. Tanimoto

cERL-FEL meeting, Jan. 20, 2021

Outline

- **Simulation of CSR, LSC and TSC effects using SAD and GPT**
 - * Use magnet layout [Daihon20200607](#)
 - * Use initial beam of ideal 3D Gaussian distribution
- **Scaling laws of momentum spread increase due to SC**
- **Summary**

1. Introduction

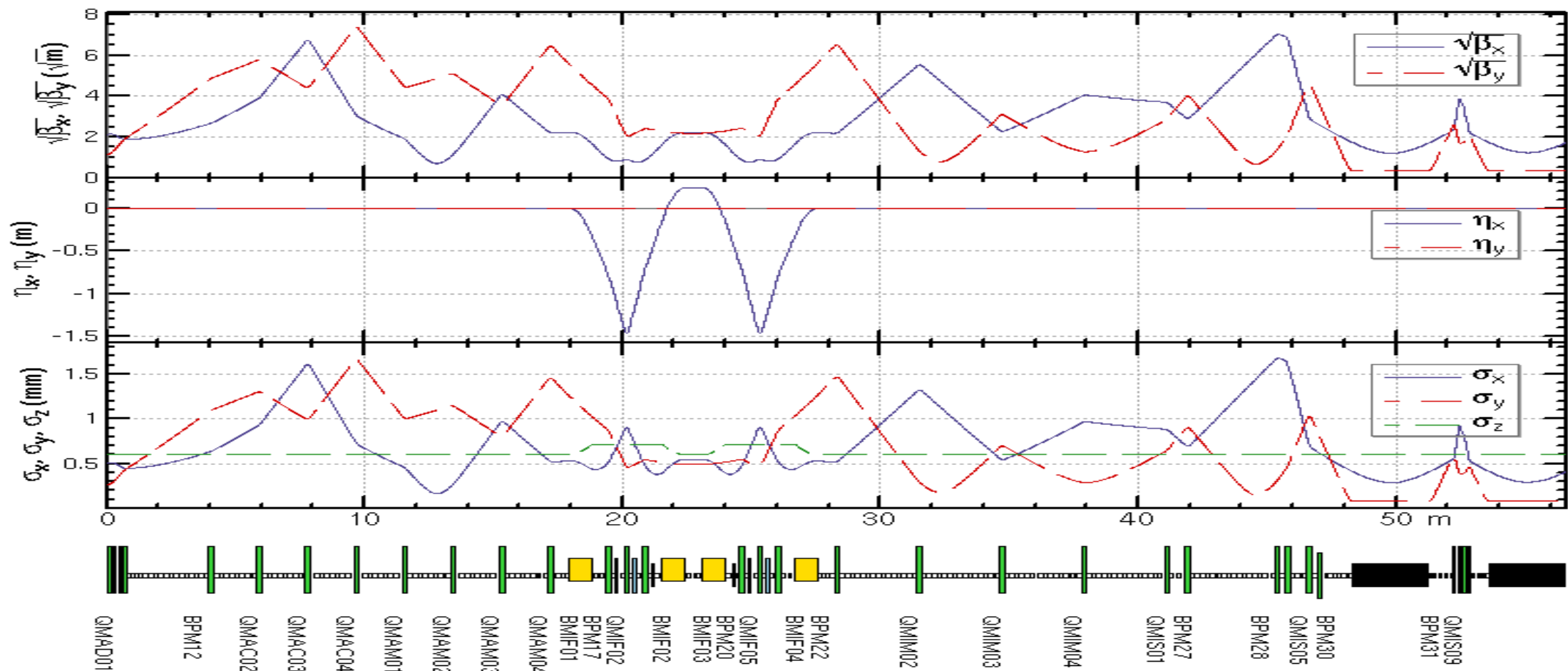
Layout:

Start from POINTA2 (50 mm before POINTA)

Beam energy 17.5 MeV

Arc $R_{56} = 0.0$ m

```
POINTA2  =(AX=-0.86752 BX=4.459129 AY=-0.59393 BY=1.273629
           AZ=0. BZ=1.0006922855944562
           EMITX=5.6888827897707966e-08 EMITY=5.0746293858979176e-08
           EMITZ=3.5975094959999995e-07)
```



2. Tracking simulations

On-axis and average LSC models (Ref: <https://arxiv.org/abs/2101.04369>)

* On-axis LSC model

$$Z_{\parallel}(k) = \frac{iZ_0k}{4\pi\beta^2\gamma^2} e^{\xi_{\sigma}^2/2} \Gamma(0, \xi_{\sigma}^2/2) \quad \xi_{\sigma} = \frac{k\sigma}{\beta\gamma} \quad \sigma = \frac{1}{2}(\sigma_x + \sigma_y)$$

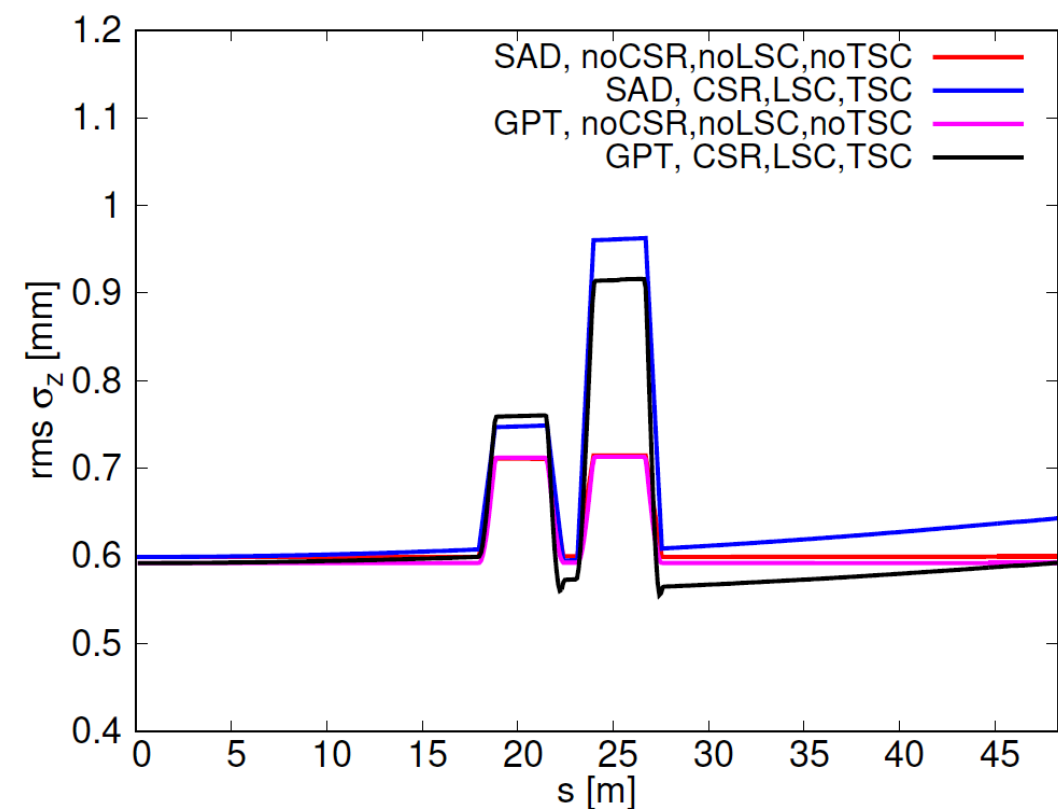
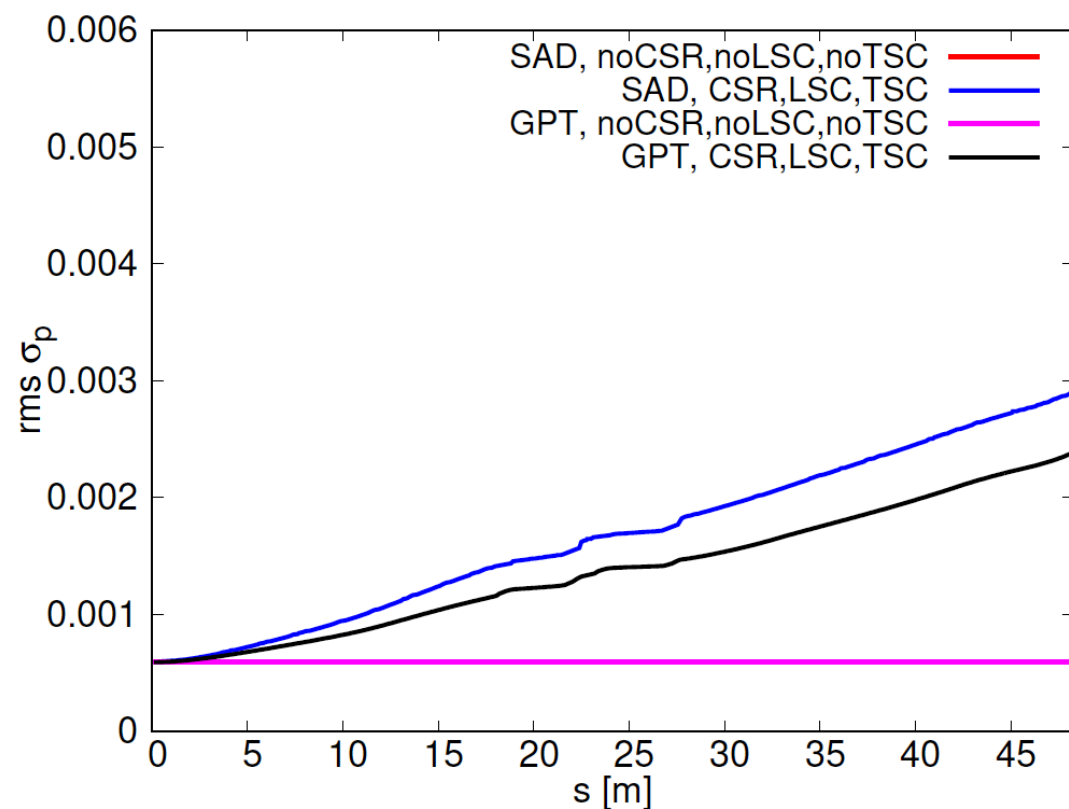
* Average LSC model (Average over the transverse beam density)

$$Z_{\parallel}(k) = \frac{iZ_0k}{4\pi\beta^2\gamma^2} e^{\xi_{\sigma}^2} \Gamma(0, \xi_{\sigma}^2)$$

2. Tracking simulations

rms energy spread and bunch length with **on-axis** LSC model

- * Clear discrepancy in energy spread => 3D LSC model is necessary. Currently SAD simulation use 1D LSC model (same as that used in ELEGANT)
- * Clear discrepancy in bunch length => Result of difference in energy spread (energy spread translated to bunch length through the arc)

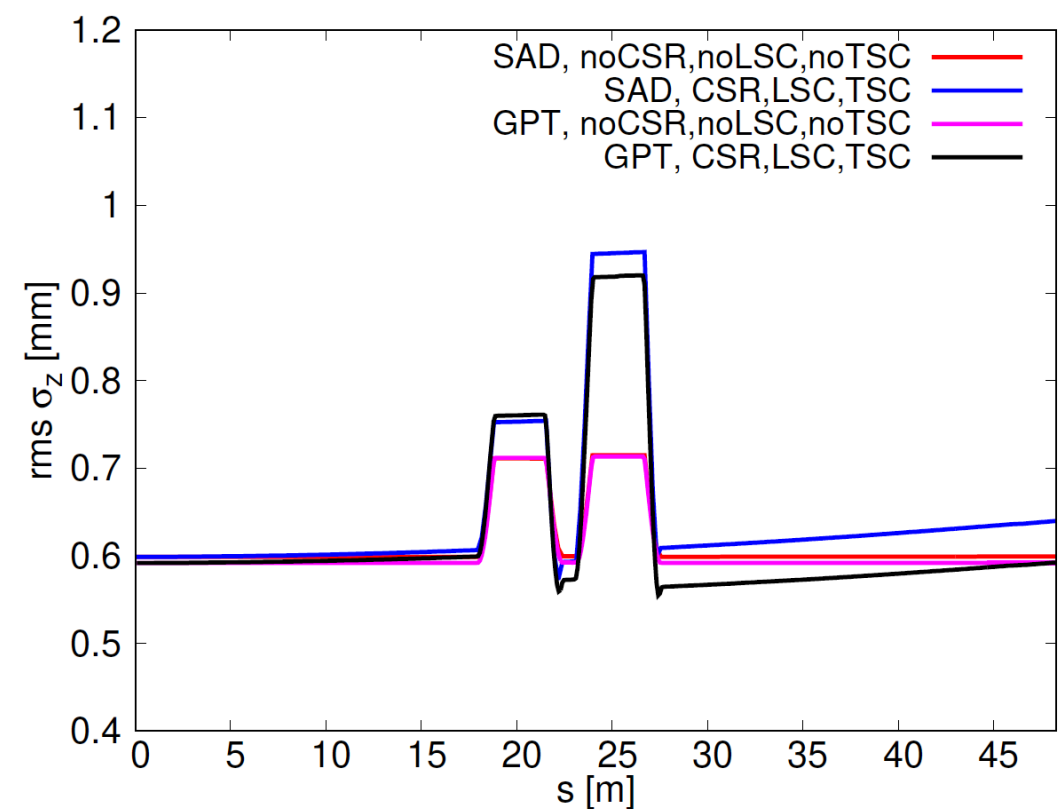
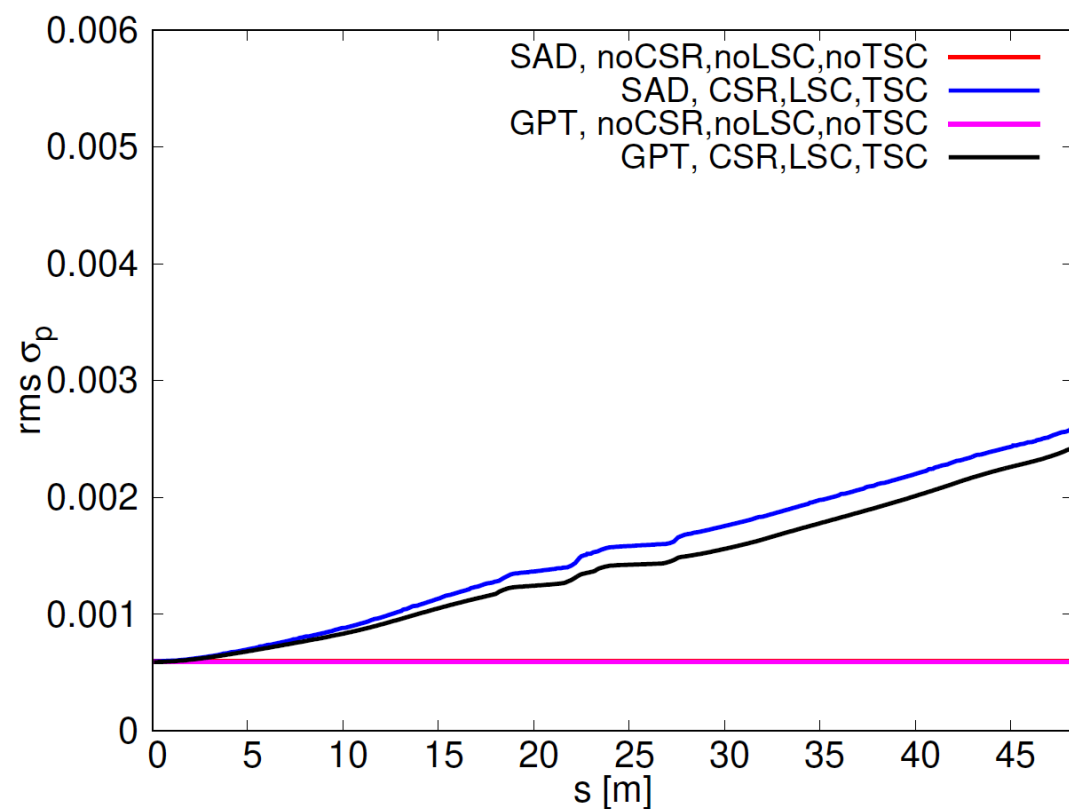


2. Tracking simulations

rms energy spread and bunch length with **average** LSC model

* Smaller discrepancy in energy spread => Average model is better. Remaining discrepancy might be due to 3D effects

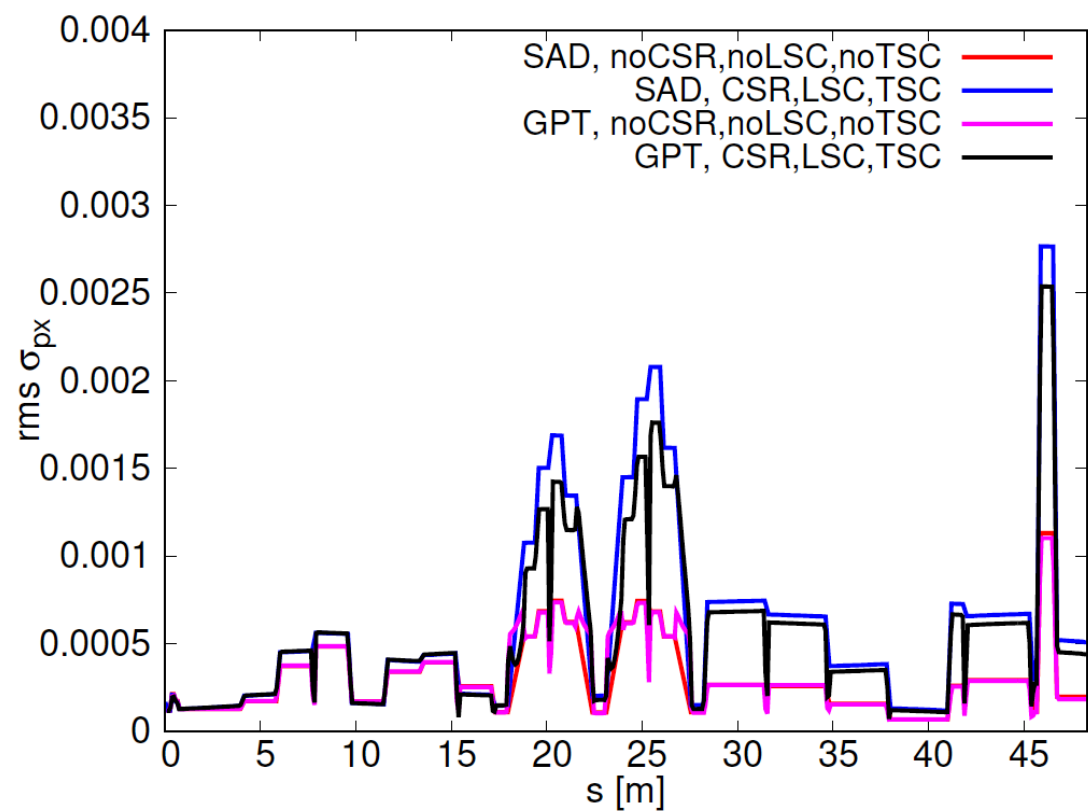
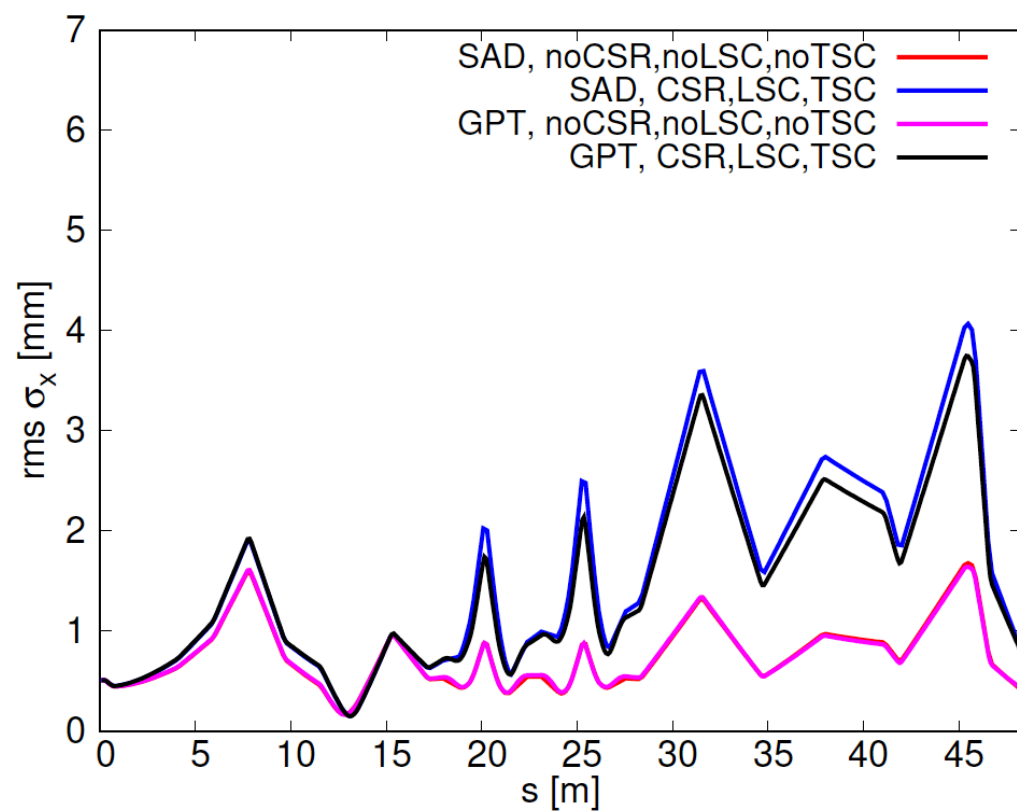
* Discrepancy in bunch length => Related to energy spread and CSR (CSR models are different between SAD and GPT)?



2. Tracking simulations

rms horizontal beam size and momentum spread with **on-axis**
LSC model

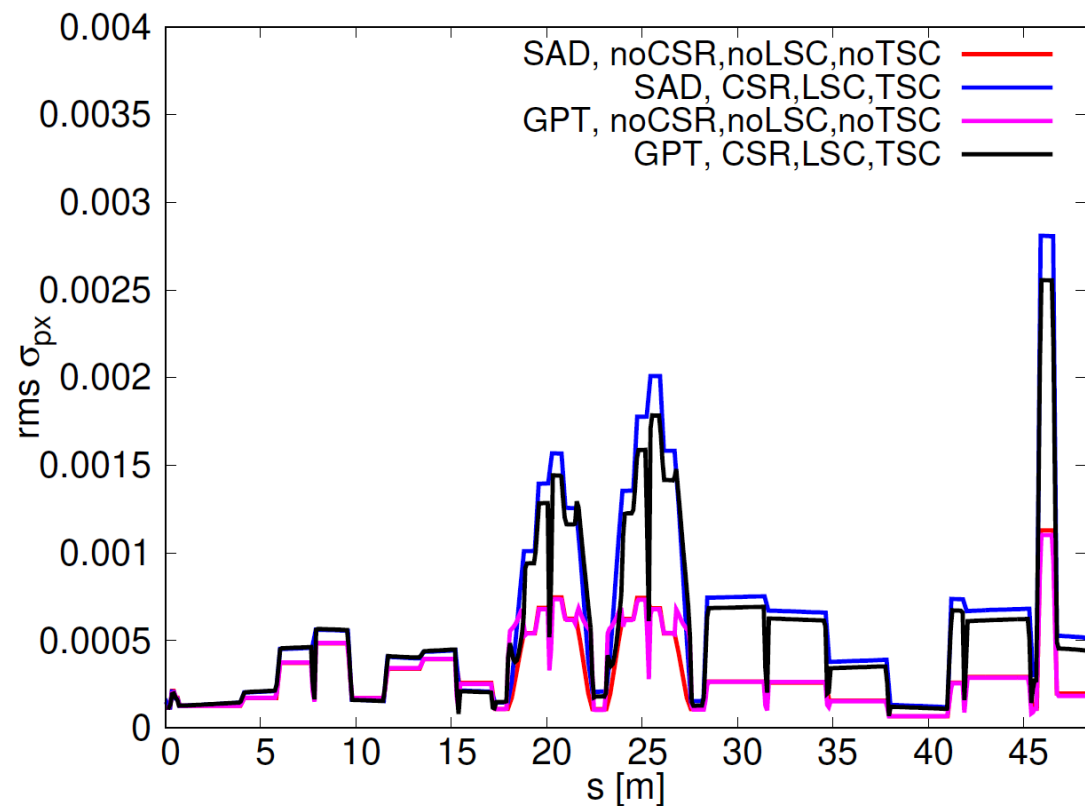
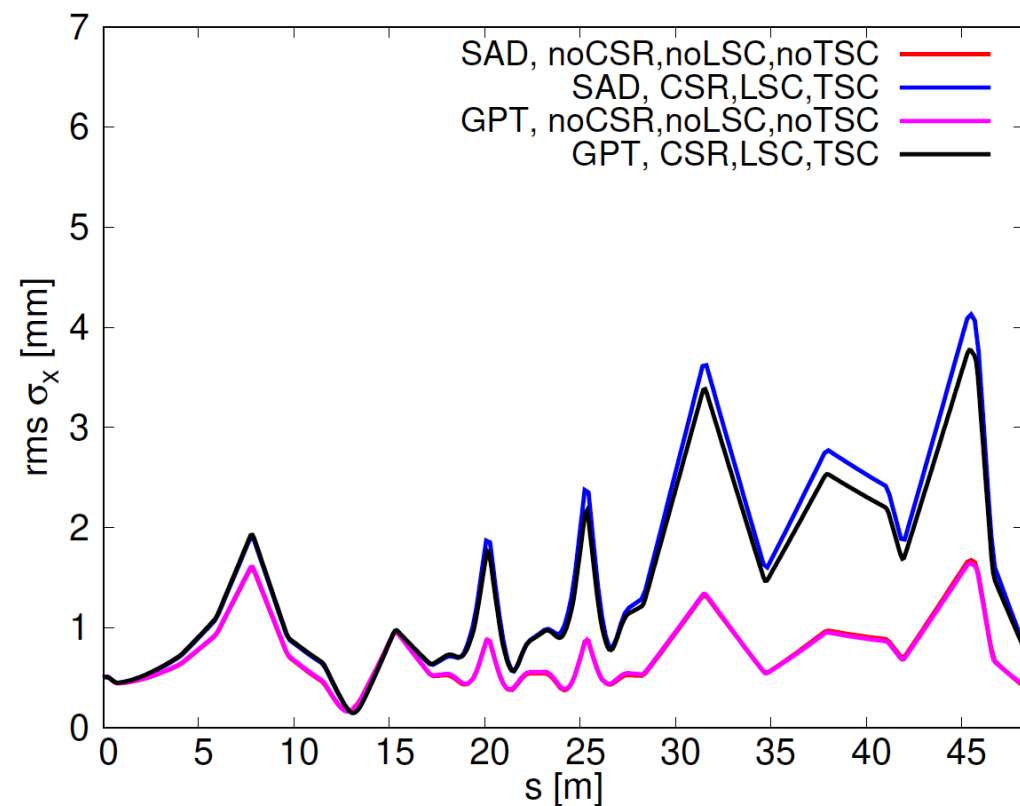
- * Good agreement in general
- * After the arc, difference appear to be remarkable. Can be correlated with the difference in energy spread



2. Tracking simulations

rms horizontal beam size and momentum spread with **average** LSC model

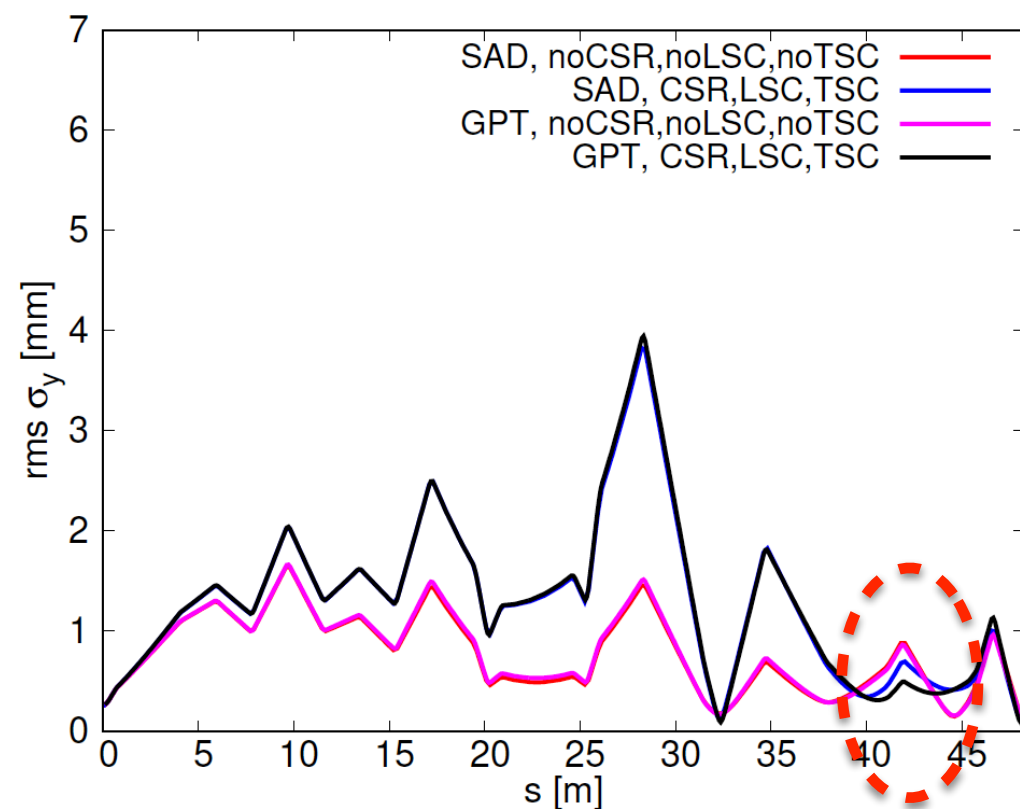
- * Good agreement in general
- * After the arc, difference appear to be remarkable. Can be correlated with the difference in energy spread



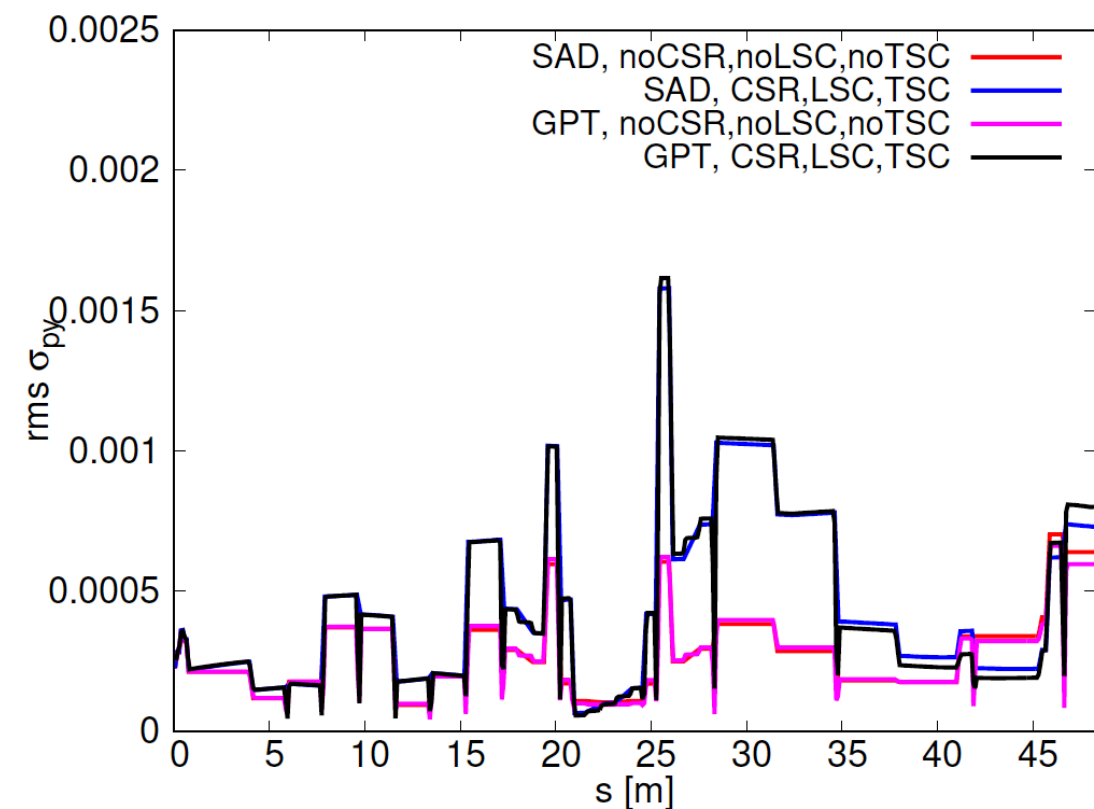
2. Tracking simulations

rms vertical beam size and momentum spread with **on-axis LSC** model

- * Good agreement in general
- * The arc does no make difference since it does not create coupling in Y and Z directions.



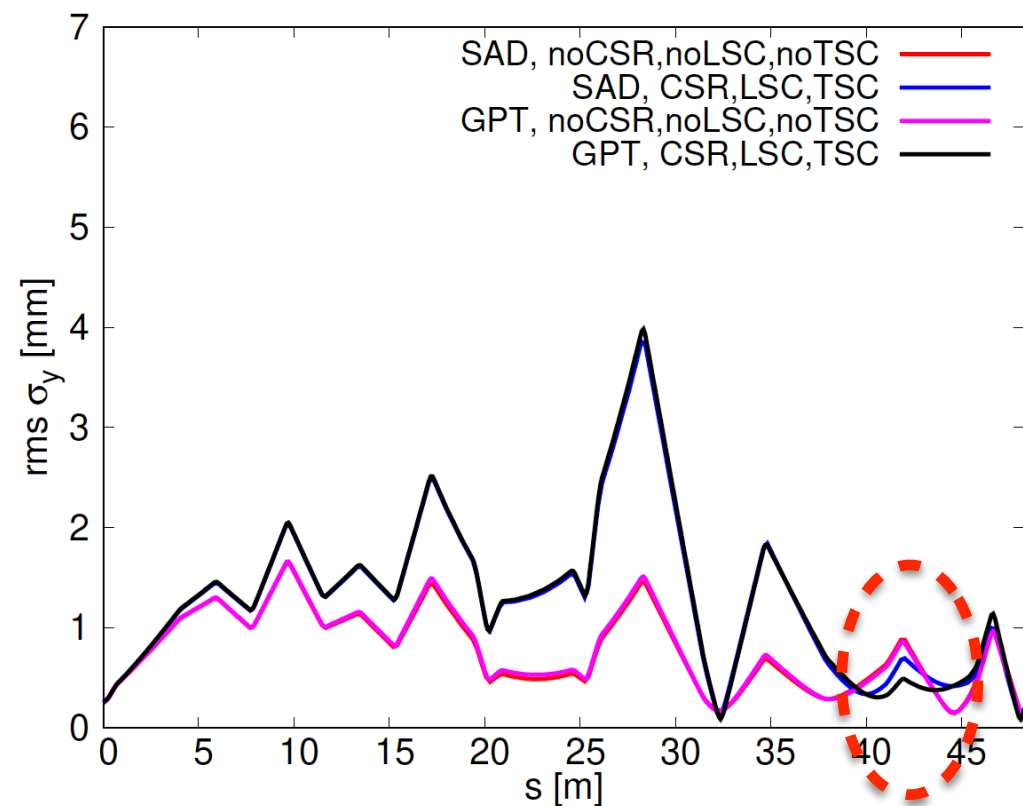
A little strange!



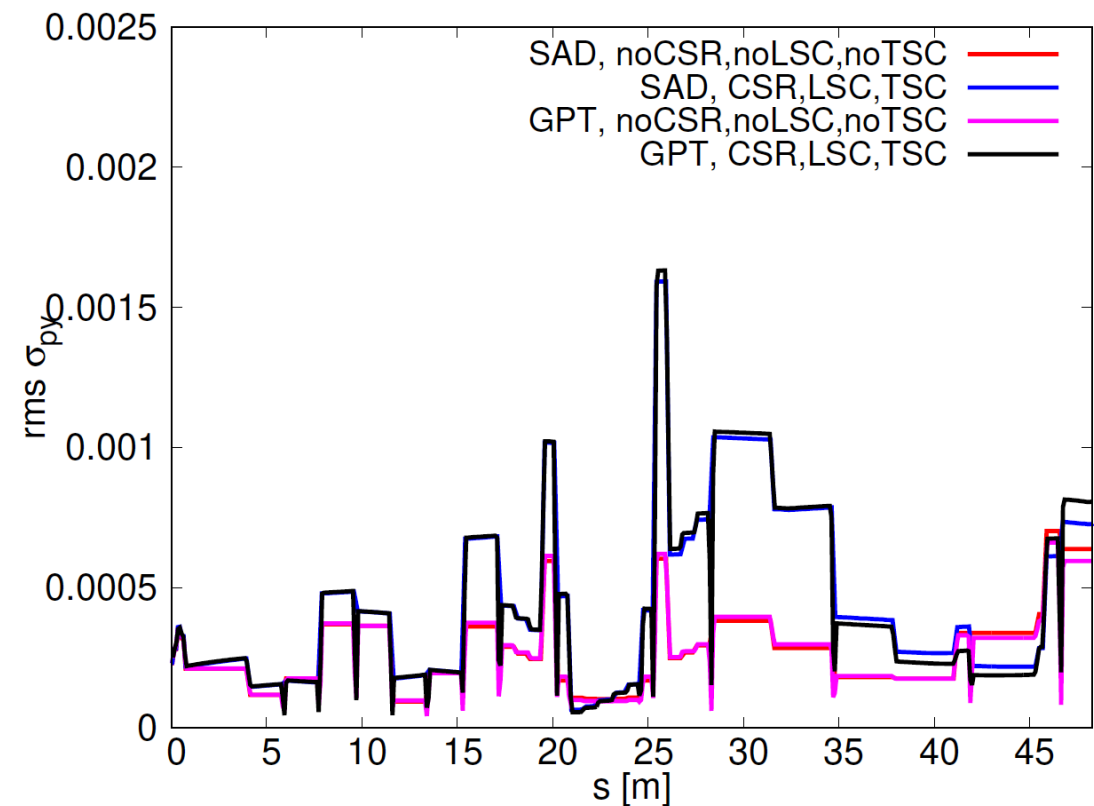
2. Tracking simulations

rms vertical beam size and momentum spread with **average LSC** model

- * Good agreement in general
- * The arc does no make difference since it does not create coupling in Y and Z directions.



A little strange!



3. Theory of SC effects

Scaling law for energy spread increase due to LSC [Refer to my paper under preparation]

* Energy spread due to LSC with path length L along the beam line:

$$\sigma_p^{\text{LSC}} = \frac{3^{-\frac{3}{4}} N r_e \bar{g} L}{\sqrt{2\pi} \gamma^3 \sigma_z^2} \quad \bar{g} = 3 \ln 2 - 2 - 2 \ln (\alpha_x + \alpha_y)$$
$$\alpha_x = \sigma_x / (\gamma \sigma_z) \quad \alpha_y = \sigma_y / (\gamma \sigma_z)$$

This formulation is only valid for $\alpha_x \ll 1$ and $\alpha_y \ll 1$.

* More general equation:

$$\sigma_p'(s) = \frac{\partial \sigma_p(s)}{\partial s} = \frac{3^{-\frac{3}{4}} N r_e \bar{g}(s)}{\sqrt{2\pi} \gamma^3(s) \sigma_z^2(s)}$$
$$\sigma_p(s) = \sqrt{\sigma_{p0}^2 + \left[\int_0^s \sigma_p'(s') ds' \right]^2}$$

* Predictions of LSC driven energy spread:

- ** Linearly proportional to bunch population
- ** Roughly linear proportional to length of beam line
- ** Inverse cubic power law for beam energy
- ** Inverse-square law for bunch length

3. Theory of SC effects

Scaling law for energy spread increase due to LSC [Refer to my paper under preparation]

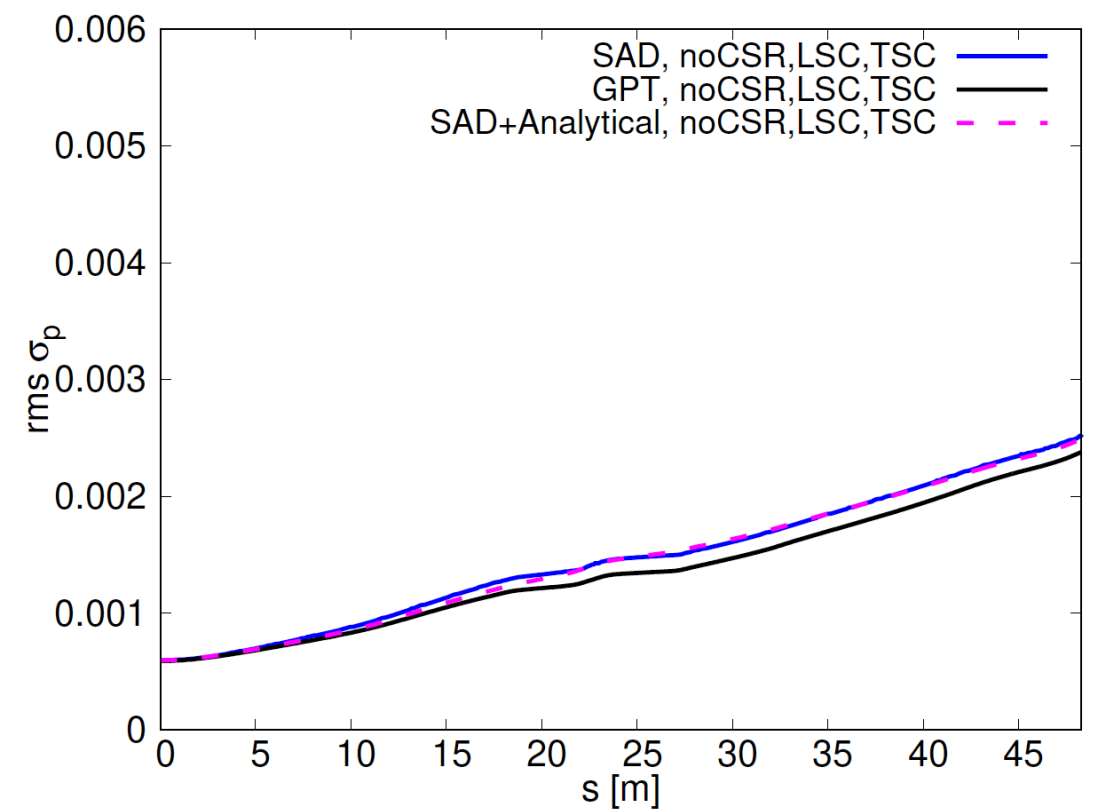
*** Apply to cERL-FEL (magenta line in the figure):**

$$\sigma'_p(s) = \frac{\partial \sigma_p(s)}{\partial s} = \frac{3^{-\frac{3}{4}} N r_e \bar{g}(s)}{\sqrt{2\pi} \gamma^3(s) \sigma_z^2(s)}$$

$$\sigma_p(s) = \sqrt{\sigma_{p0}^2 + \left[\int_0^s \sigma'_p(s') ds' \right]^2}$$

*** Very good agreement:**

**** In the equation, beam sizes taken from SAD simulations**



3. Theory of SC effects

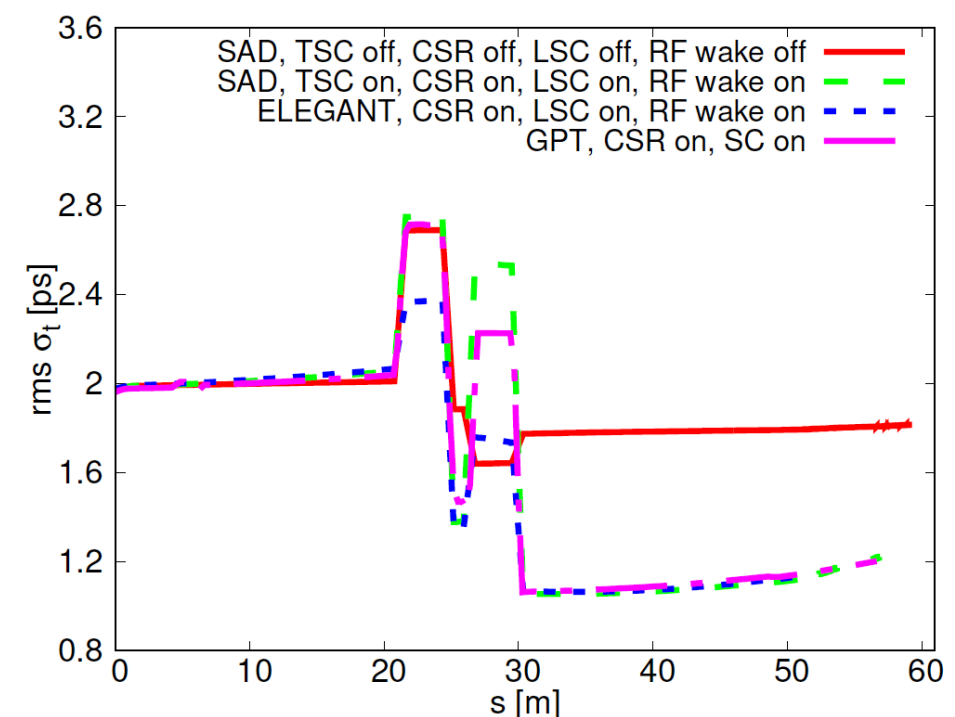
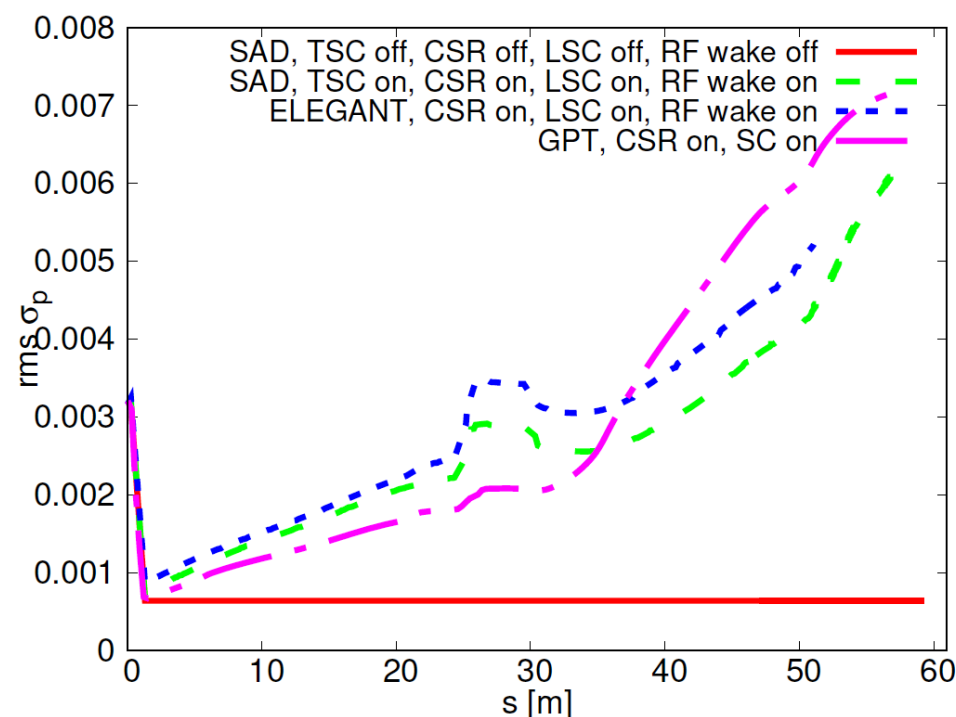
Scaling law for energy spread increase due to LSC [Refer to my paper under preparation]

* Apply to cERL-FEL (magenta line in the figure):

$$\sigma'_p(s) = \frac{\partial \sigma_p(s)}{\partial s} = \frac{3^{-\frac{3}{4}} N r_e \bar{g}(s)}{\sqrt{2\pi} \gamma^3(s) \sigma_z^2(s)} \quad \sigma_p(s) = \sqrt{\sigma_{p0}^2 + \left[\int_0^s \sigma'_p(s') ds' \right]^2}$$

* The theory explains the extra increase of energy spread after the arc with $R_{56}=0.217$ configuration:

** Shorter bunch length causes faster increase of energy spread



3. Theory of SC effects

Scaling law for momentum spread increase due to TSC [Refer to my paper under preparation]

*** Momentum spread due to TSC in a drift:**

$$\sigma'_{p_x}(s) = \frac{\partial \sigma_{p_x}(s)}{\partial s} = \frac{3^{-\frac{1}{4}} \sqrt{\ln \frac{4}{3}} N r_e \bar{g}(s)}{\sqrt{2\pi} \gamma^3(s) \sigma_z(s) \sigma(s)} \quad \sigma = (\sigma_x + \sigma_y)/2$$

*** In a quadrupole, there is another term to be counted.**

*** To be benchmarked with simulations.**

4. Theory of beam envelope equation

So far the simulations and theory I used are based on the assumption of 3D Gaussian beams (The longitudinal distribution is not necessarily to be Gaussian). Initial Gaussian distribution used is also a reason for good agreement between SAD and GPT benchmark simulations

Question: How good is this assumption?

If this is a good assumption, we can derive beam envelope equation using the space-charge forces of Gaussian beams.

4. Theory of beam envelope equation

A revisit of beam envelope equation

II. BEAM ENVELOPE EQUATION WITH SPACE CHARGE

The beam envelope equations with space charge are [5]

$$\sigma_x'' + K_x(s)\sigma_x = \frac{\epsilon_x^2}{\sigma_x^3} + \frac{\langle x f_x \rangle}{\sigma_x}, \quad (9a)$$

$$\sigma_y'' + K_y(s)\sigma_y = \frac{\epsilon_y^2}{\sigma_y^3} + \frac{\langle y f_y \rangle}{\sigma_y}, \quad (9b)$$

$\langle \rangle$ means average over the beam density

$$\langle x f_x \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx dy dz \rho(z) \psi(x, y) x f_x(x, y, z). \quad (10)$$

With nonzero dispersion, the above equations should be extended to include dispersion functions [6, 7]. The space charge force is defined as

$$f_x = \frac{1}{\gamma m_0 \beta^2 c^2} F_x. \quad (11)$$

4. Theory of beam envelope equation

A revisit of beam envelope equation

For beams with transverse Gaussian distribution, there is[8, 9]

$$F_x(x, y) = \frac{Ne^2 Z_0 c}{2\pi\gamma^2} \rho(z) \int_0^\infty \frac{x}{(t + 2\sigma_x^2)^{3/2} (t + 2\sigma_y^2)^{1/2}} e^{-\frac{x^2}{t+2\sigma_x^2} - \frac{y^2}{t+2\sigma_y^2}} dt. \quad (12)$$

This space charge force is valid for $\alpha_{x,y} = \sigma_{x,y}/(\gamma\sigma_z) \ll 1$. The ELEGANT code used the on-axis force to model the linear TSC effects:

$$F_x(0, 0) = \frac{Ne^2 Z_0 c}{2\pi\gamma^2} \frac{x\rho(z)}{\sigma_x(\sigma_x + \sigma_y)}. \quad (13)$$

Linear TSC model used in ELEGANT
=> Overestimate TSC effects

With the above definitions, we can find:

$$\langle x f_x \rangle = \frac{Nr_e \sigma_x}{2\sqrt{\pi}\beta^2\gamma^3\sigma_z(\sigma_x + \sigma_y)}. \quad (14)$$

Apply Eq. (14) to Eq. (9a), we can obtain:

$$\sigma_x'' + K_x(s)\sigma_x = \frac{\epsilon_x^2}{\sigma_x^3} + \frac{I}{I_0\beta^3\gamma^3(\sigma_x + \sigma_y)}, \quad (15)$$

with $I_0 = ec/r_e$, and $I = Ne/(2\sqrt{\pi}\sigma_z/(\beta c)) = Ne/(2\sqrt{\pi}\sigma_t)$. Here I is a bunched beam with Gaussian distribution and also the average force of $\langle x f_x \rangle$ is considered. The scaling factor $2\sqrt{\pi}$ is slightly different from that in the literature.

For a quick check of TSC effects at cERL-FEL, we use some typical beam parameters:

$\beta\gamma\epsilon_x = 2 \mu\text{m}$, $\bar{\sigma}_x = 1.4 \text{ mm}$, $\bar{\sigma}_y = 1.3 \text{ mm}$, $\bar{\sigma}_z = 0.65 \text{ mm}$, $\gamma = 34.2$. Then we can obtain $\frac{\epsilon_x^2}{\bar{\sigma}_x^3} \approx 1.3 \times 10^{-6} \text{ m}^{-1}$, and $\frac{\langle x f_x \rangle}{\bar{\sigma}_x} \approx 4.2 \times 10^{-6} \text{ m}^{-1}$. Since $\frac{\langle x f_x \rangle}{\bar{\sigma}_x} > \frac{\epsilon_x^2}{\bar{\sigma}_x^3}$, we can conclude that TSC is important in cERL-FEL.

4. Theory of beam envelope equation

A revisit of beam envelope equation (references)

- [5] T. P. Wangler, Multiparticle dynamics with space charge, in *RF Linear Accelerators* (John Wiley & Sons, Ltd, 2008) Chap. 9, pp. 282–340.
- [6] H. Okamoto and S. Machida, Particle beam resonances driven by dispersion and space charge, Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment **482**, 65 (2002).
- [7] A. Khan, O. Boine-Frankenheim, F. Hug, and C. Stoll, Beam matching with space charge in energy recovery linacs, Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment **948**, 162822 (2019).
- [8] D. Zhou and Y. Nie, Wake fields and impedance of space charge in cartesian coordinate system (2021), arXiv:2101.04369 [physics.acc-ph].
- [9] Taufik and K. Takayama, Beam-core evolution equation and space-charge limit, Physics Letters A **383**, 125855 (2019).

4. Theory of beam envelope equation

From ELEGANT manual:

This element simulates transverse space charge (SC) kicks using K.Y. Ng's formula [24].

The linear SC force is given by:

$$\Delta x' = \frac{K_{sc} L e^{-z^2/(2\sigma_z^2)}}{\sqrt{2\pi}\sigma_z} \frac{x}{\sigma_x(\sigma_x + \sigma_y)}$$
$$\Delta y' = \frac{K_{sc} L e^{-z^2/(2\sigma_z^2)}}{\sqrt{2\pi}\sigma_z} \frac{y}{\sigma_y(\sigma_x + \sigma_y)}$$

where $K_{sc} = \frac{2Nr_e}{\gamma^3\beta^2}$, L is the integrating length, $\sigma_{x,y,z}$ are rms beam size.

The non-linear SC force is given by:

$$\Delta x' = \frac{K_{sc} L e^{-z^2/(2\sigma_z^2)}}{2\sigma_z \sqrt{\sigma_x^2 - \sigma_y^2}} \text{Im} \left[w \left(\frac{x + iy}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}} \right) - e^{-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2}} w \left(\frac{x \frac{\sigma_y}{\sigma_x} + iy \frac{\sigma_x}{\sigma_y}}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}} \right) \right]$$
$$\Delta y' = \frac{K_{sc} L e^{-z^2/(2\sigma_z^2)}}{2\sigma_z \sqrt{\sigma_x^2 - \sigma_y^2}} \text{Re} \left[w \left(\frac{x + iy}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}} \right) - e^{-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2}} w \left(\frac{x \frac{\sigma_y}{\sigma_x} + iy \frac{\sigma_x}{\sigma_y}}{\sqrt{2(\sigma_x^2 - \sigma_y^2)}} \right) \right]$$

where $w(z)$ is the complex error function

$$w(z) = e^{-z^2} \left[1 + \frac{2i}{\sqrt{\pi}} \int_0^z e^{\zeta^2} d\zeta \right]$$

5. Summary

- * Bench mark of SAD and GPT

- ** Using averaged LSC 1D model shows better agreement with GPT simulations

- * Theory of SC effects

- ** Simple scaling laws were found (Any similar theory in the literature?)

- ** Reproduce the simulations well (Condition: bunches close to Gaussian distribution)

- ** Sigma-matrix equation can be used for simulation and optics matching with space charge)?