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Laser  
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Source



# In-Gas-Jet Resonance Ionization Spectroscopy

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1. What spectral resolution is achievable in the gas jet RIS ? And How?
2. Requirements for laser radiation, pumping system ...



# Types of nozzles that will be discussed

## Convergent- Divergent Nozzle

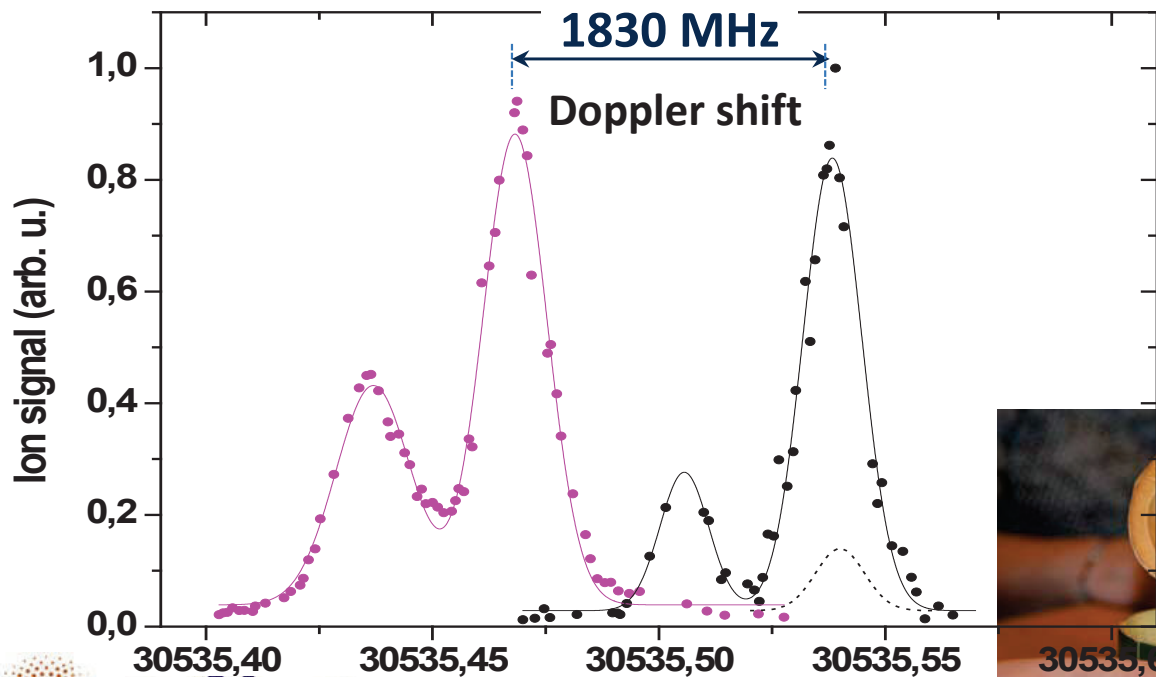


GAS CELL

$T_0 P_0$

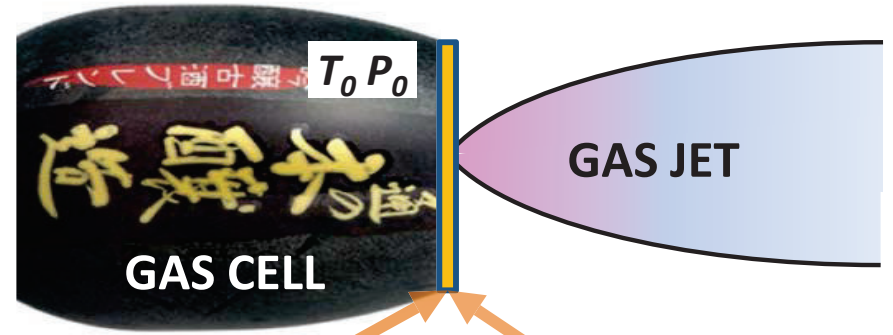
GAS JET

$T_0 = 355 \pm 3K$



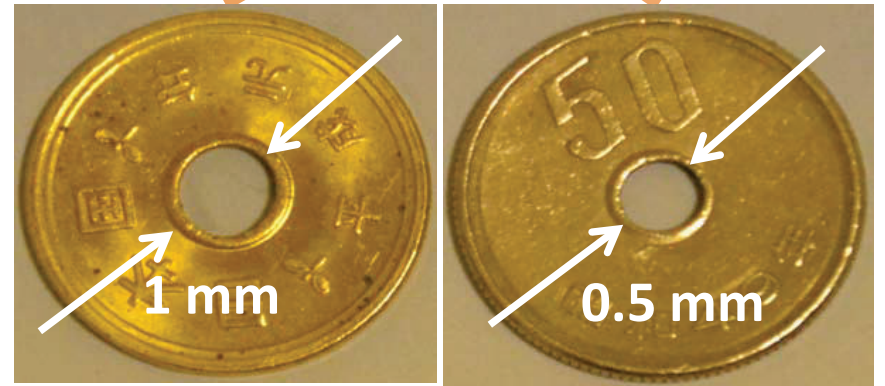
## Free-Jet Nozzle

TOP VIEW



GAS CELL

SIDE VIEW



Exit orifice diameter

Pumping speed to obtain the desired Mach number



# Properties of supersonic beams (I)

$$C_p T_0 = C_p T + \frac{u^2}{2} \quad u_{\max} = \sqrt{2C_p T_0}$$

The sum of specific enthalpy and kinetic energy remains constant during expansion

Mach number -  $M = u / a$

$u$  - gas stream velocity

$a$  - local speed of sound

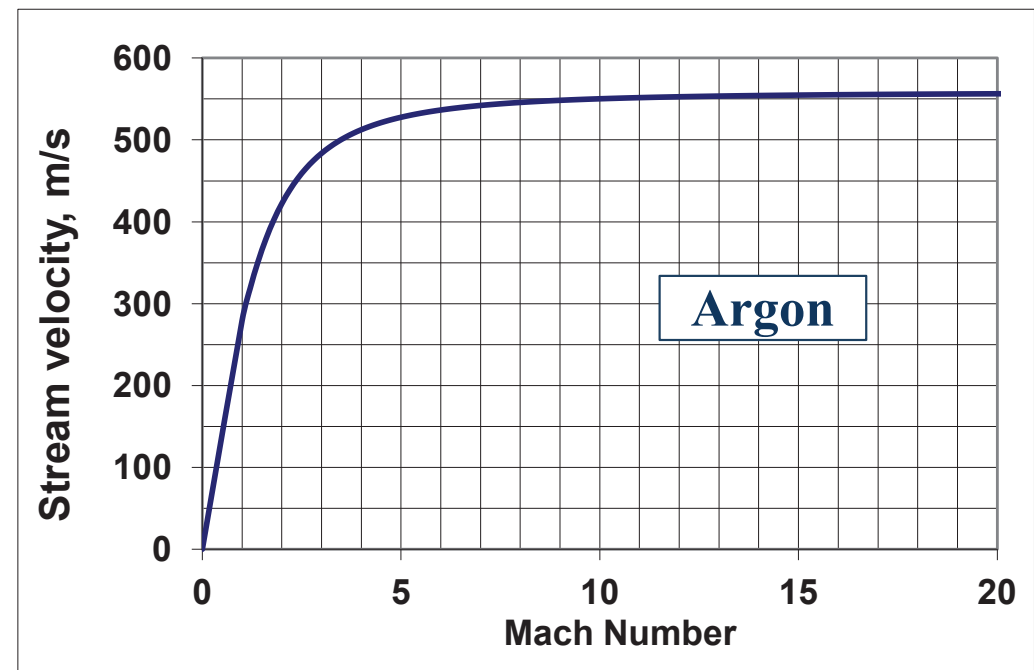
$T$  - gas temperature

$m_{ng}$  - mass of the noble gas

$$a = \sqrt{\frac{\gamma k T}{m_{ng}}}$$

$$u = \sqrt{\frac{k T_0 \gamma M^2}{m_{ng} \left\{ 1 + \left[ \frac{\gamma - 1}{2} \right] M^2 \right\}}}$$

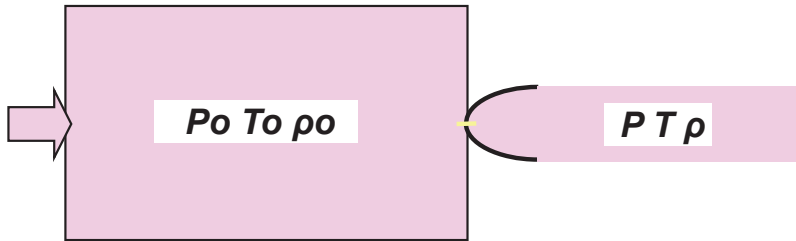
$\gamma = C_p / C_v$  - ratio of specific heat capacities = 5/3



Stream velocity Ar-550 m/s

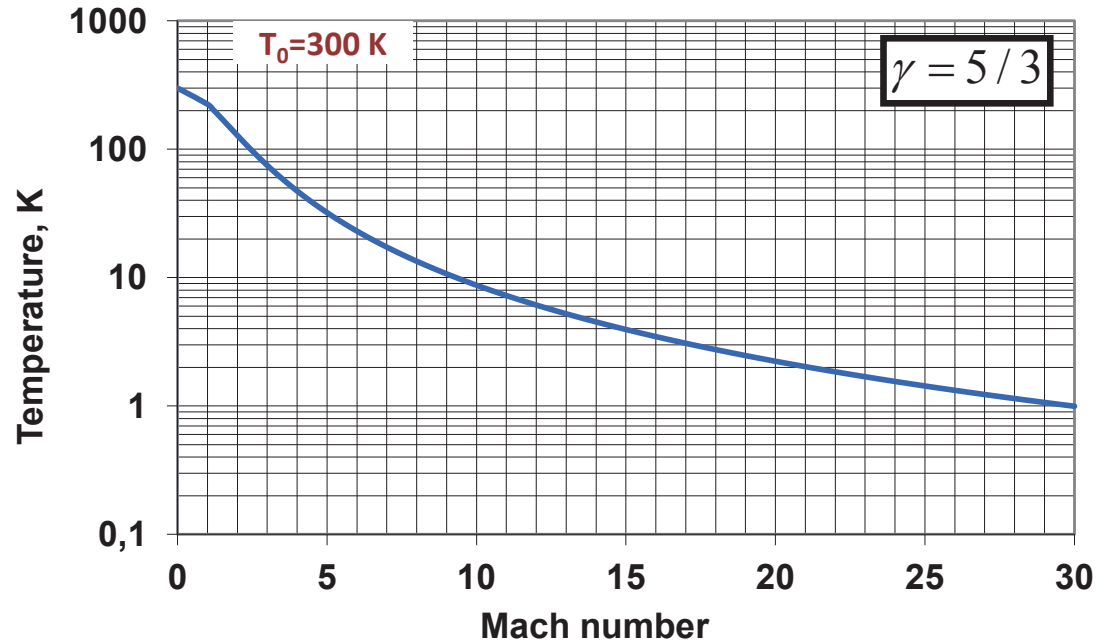
He  $\sim \sqrt{m_{ng}}^{-1}$

# Properties of supersonic beams (II)

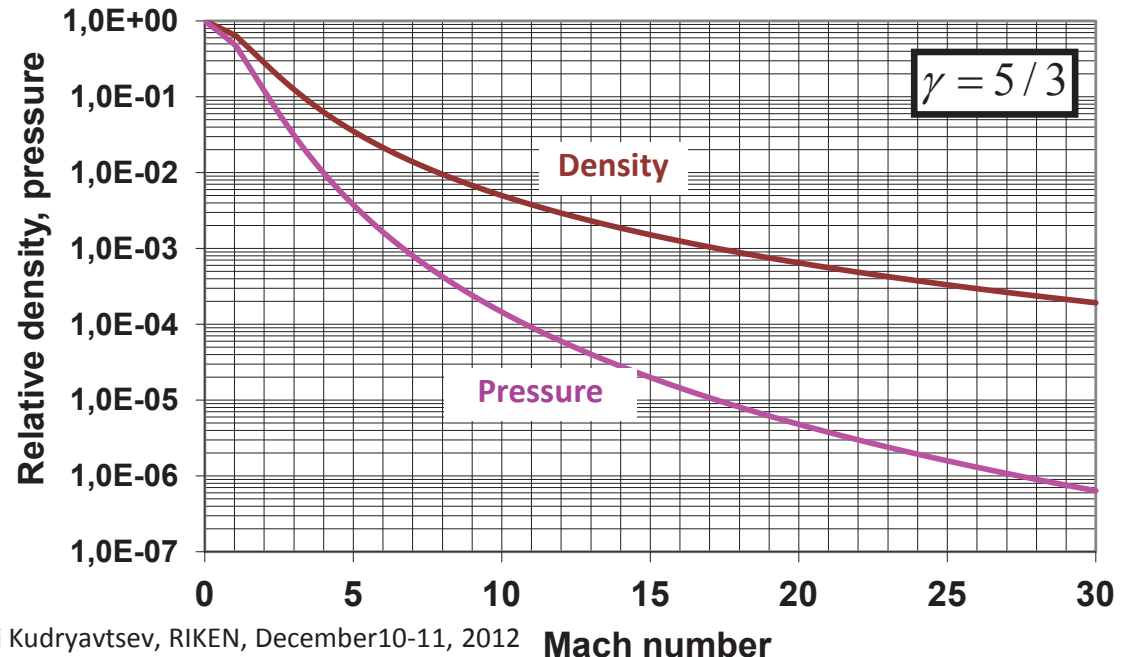


Isentropic flow - gas entropy is constant

$$\frac{T}{T_0} = \left[ 1 + \left( \frac{\gamma - 1}{2} \right) M^2 \right]^{-1}$$

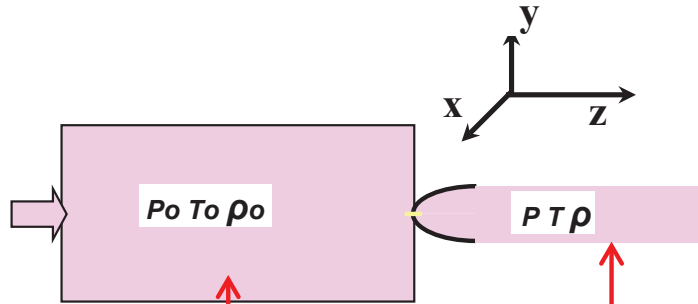


$$\frac{\rho}{\rho_0} = \left[ 1 + \left( \frac{\gamma - 1}{2} \right) M^2 \right]^{-\frac{1}{\gamma - 1}}$$



$$\frac{P}{P_0} = \left[ 1 + \left( \frac{\gamma - 1}{2} \right) M^2 \right]^{-\frac{\gamma}{\gamma - 1}}$$

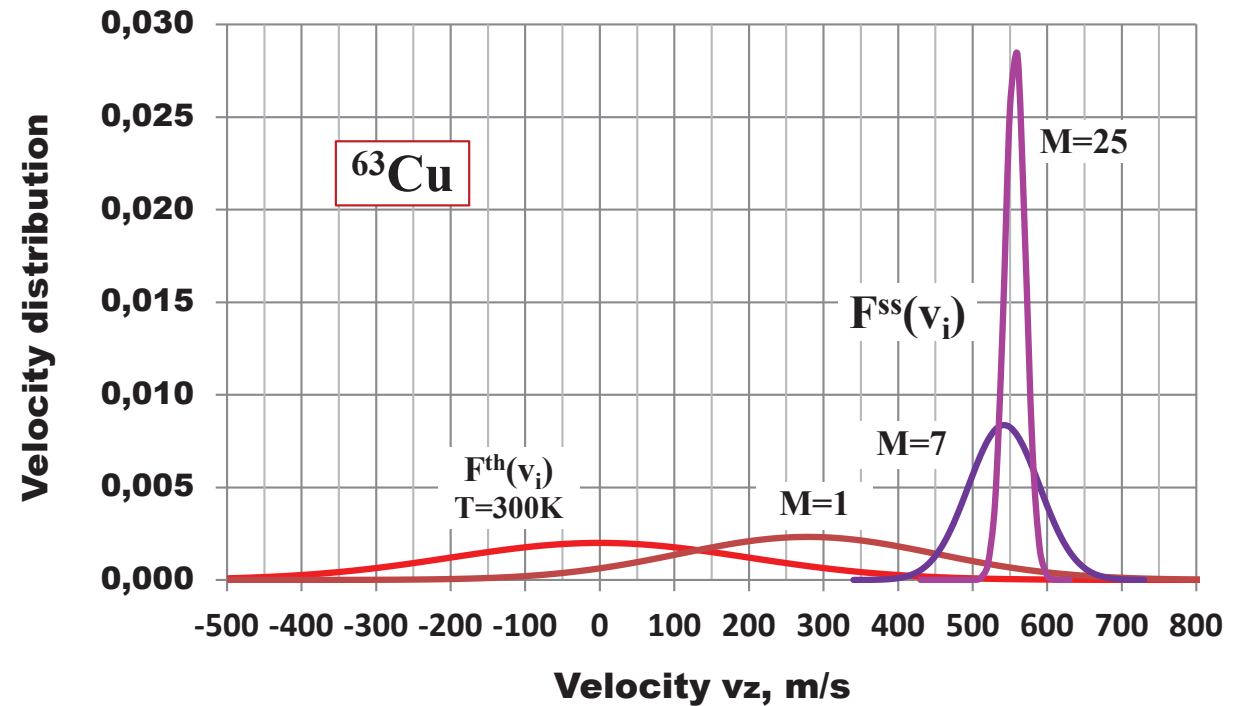
# Properties of supersonic beams (III)



$$F^{th}(v_i) = \sqrt{\frac{m}{2\pi kT_0}} \exp\left(\frac{-mv_i^2}{2kT_0}\right)$$

$$F^{ss}(v_z) = \sqrt{\frac{m}{2\pi kT}} \exp\left(\frac{-m(v_z-u)^2}{2kT}\right)$$

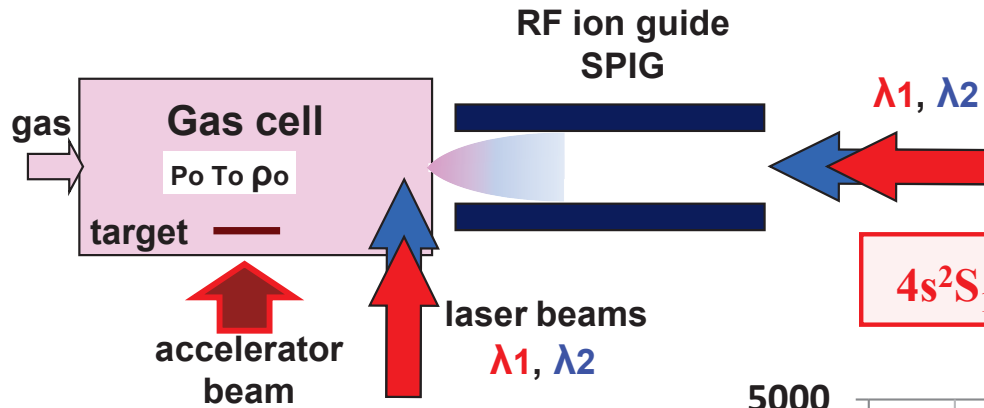
One dimensional Maxwell-Boltzmann velocity distribution



Radioactive atoms are in thermal equilibrium with buffer gas atoms

$$\delta(F) = 2\sqrt{\ln 2} \sqrt{\frac{2kT}{m}} \text{ - Full Width at Half Maximum (FWHM)}$$

# Doppler and Collision Contributions to the Spectral Line Width



$4s^2S_{1/2} - 4p^2P_{1/2}$ , 327.4 nm  $^{63}\text{Cu}$  transition,  $\nu_0 = 30535.3 \text{ cm}^{-1}$

## Collision/pressure contribution

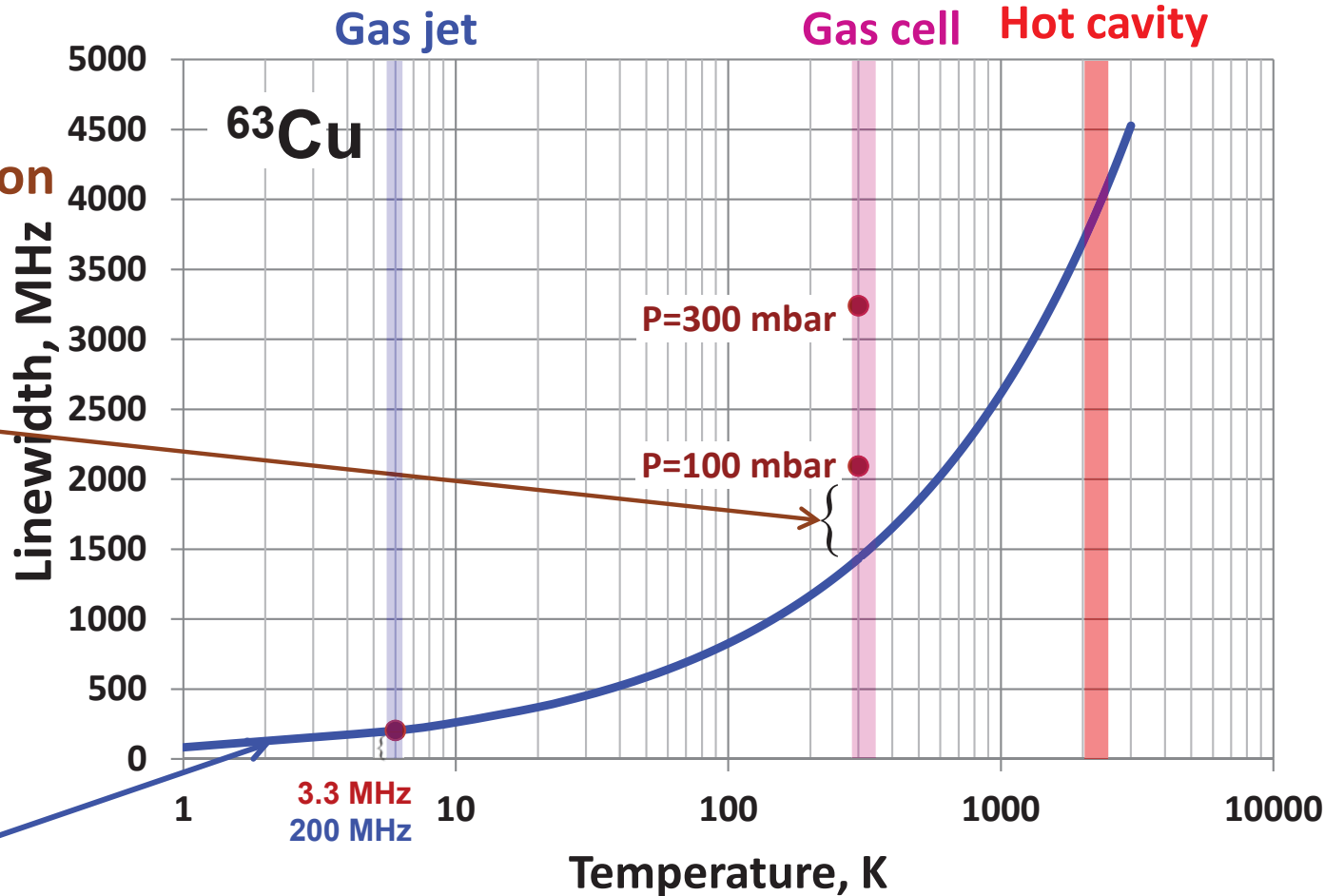
$$\Delta \nu_{coll} = \gamma_{coll} \times \rho$$

$\gamma_{coll}$  - collision broadening coefficient,  $1.5 \cdot 10^{-20} \text{ cm}^{-1}/\text{cm}^{-3}$  (8 MHz/mbar)

$\rho$  - gas density ( $\text{atom}/\text{cm}^3$ )

## Doppler contribution

$$\Delta \nu_{Doppler} = 2\sqrt{\ln 2} \frac{\nu_0}{c} \sqrt{\frac{2kT}{m}}$$



# Doppler Gaussian and collision- and natural Lorentzian contributions to the spectral line shape

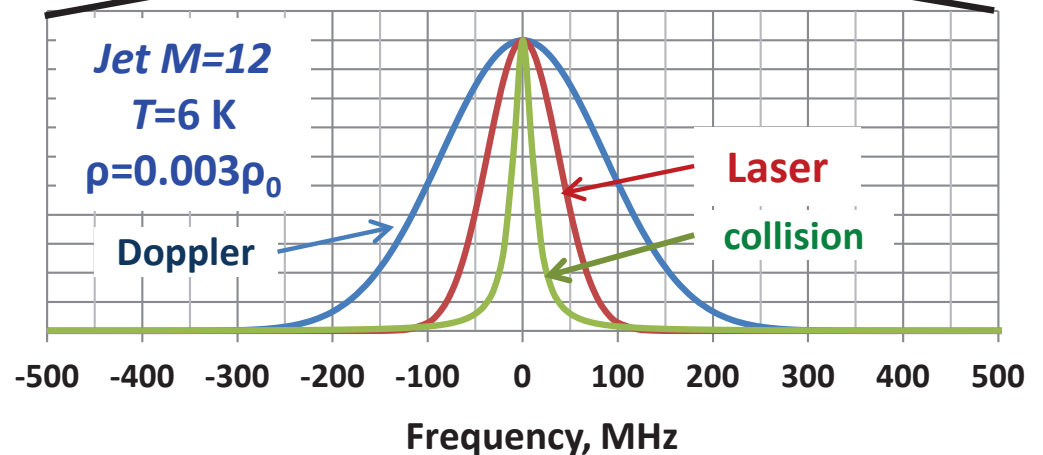
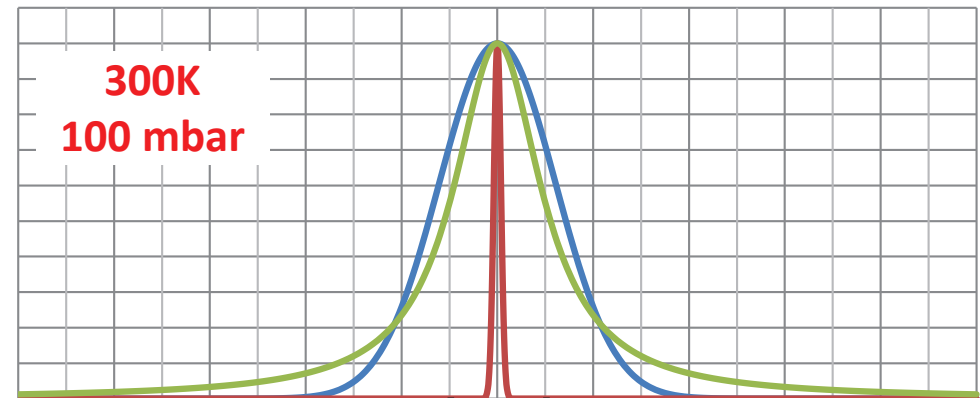
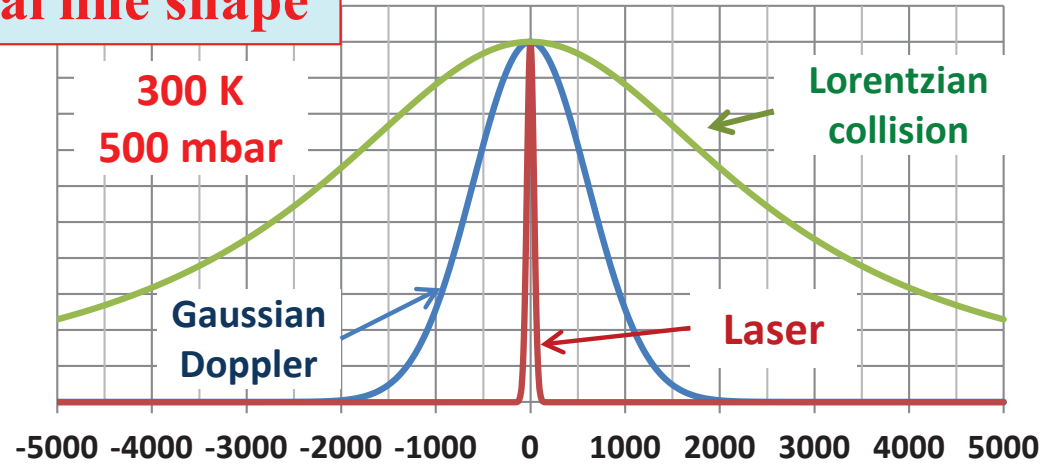
$$G(\nu) = G_0 \exp \left[ -\frac{c^2 (\nu - \nu_0)^2}{\nu_0^2 \frac{2kT}{m}} \right]$$

$$L(\nu - \nu_0) = \frac{1}{2\pi} \frac{\Gamma}{(\nu - \nu_0 + \Gamma sh)^2 + (\Gamma/2)^2}$$

Laser bandwidth –  $\delta_{laser}$  Gaussian if laser time profile is Gaussian

$$\delta_{laser} = 441 / \tau_{pulse}$$

$$\tau_{pulse} = 5 \text{ ns} \quad \delta_{laser} = 88 \text{ MHz}$$



# Laser Resonance Ionization in Supersonic Beams

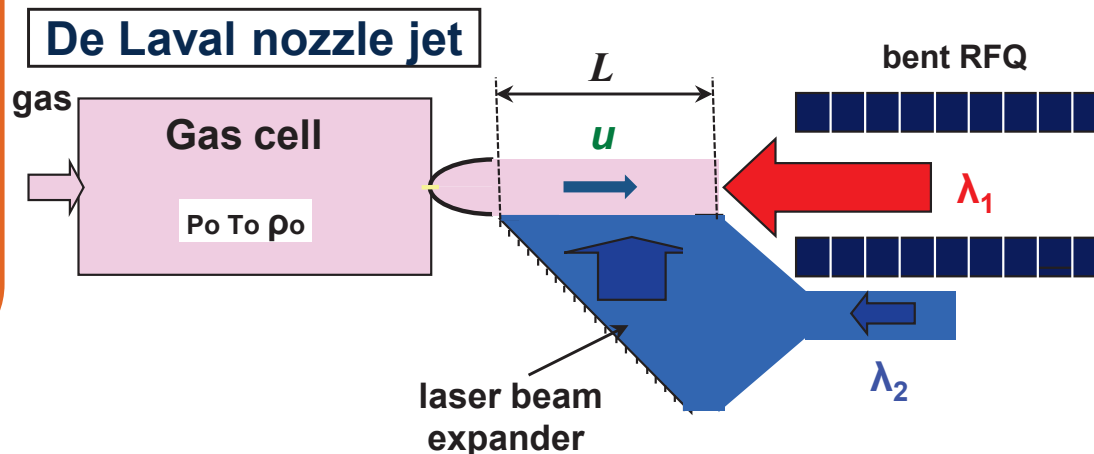
**NO laser ionization inside the cell !**  
**Laser ionization- only in the cold jet**

$\lambda_2$  !

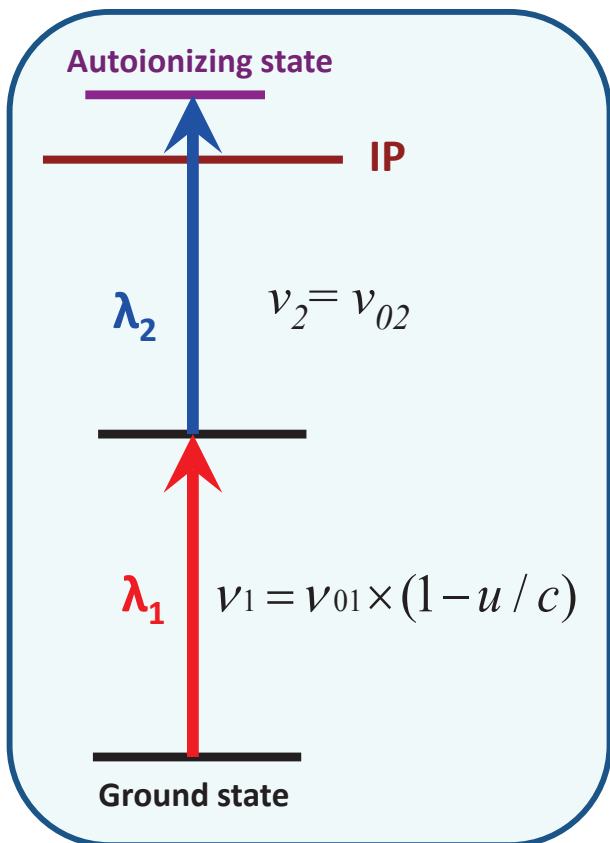
$$f_{laser} \geq 1 / (L/u) \geq 10 \text{ kHz, argon jet - } L = 5.5 \text{ cm}$$

$u$  – stream velocity, 550m/s

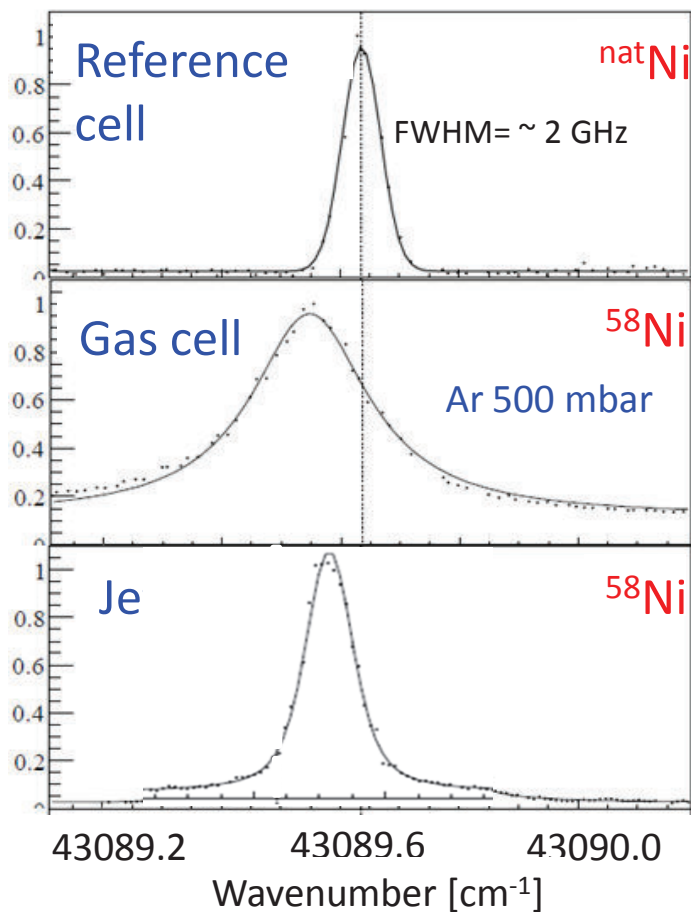
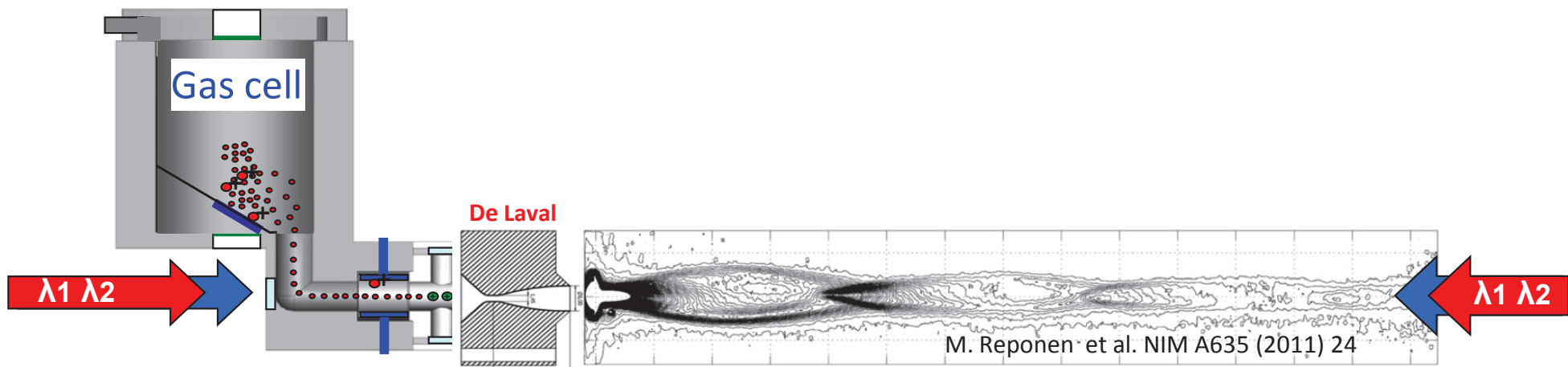
## Crossed laser beams with supersonic jet



**The parallel beam from de Laval nozzle !**  
**No broadening due to the beam divergence**  
**Very careful design of the nozzle is required**





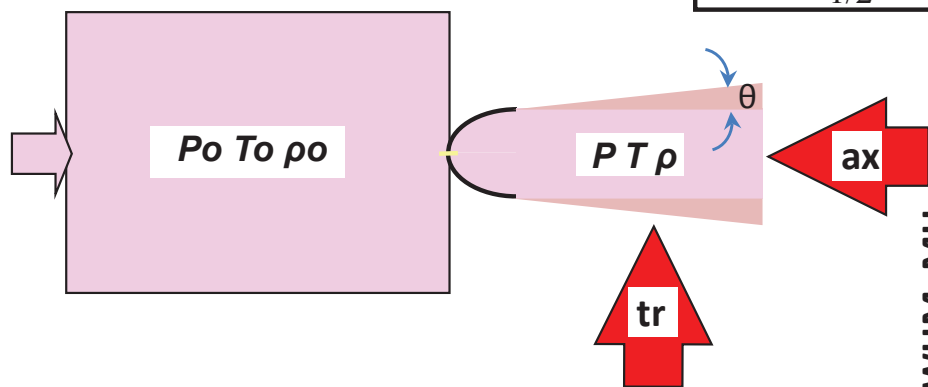


T. Sonoda et al. NIM B267 (2009) 2908

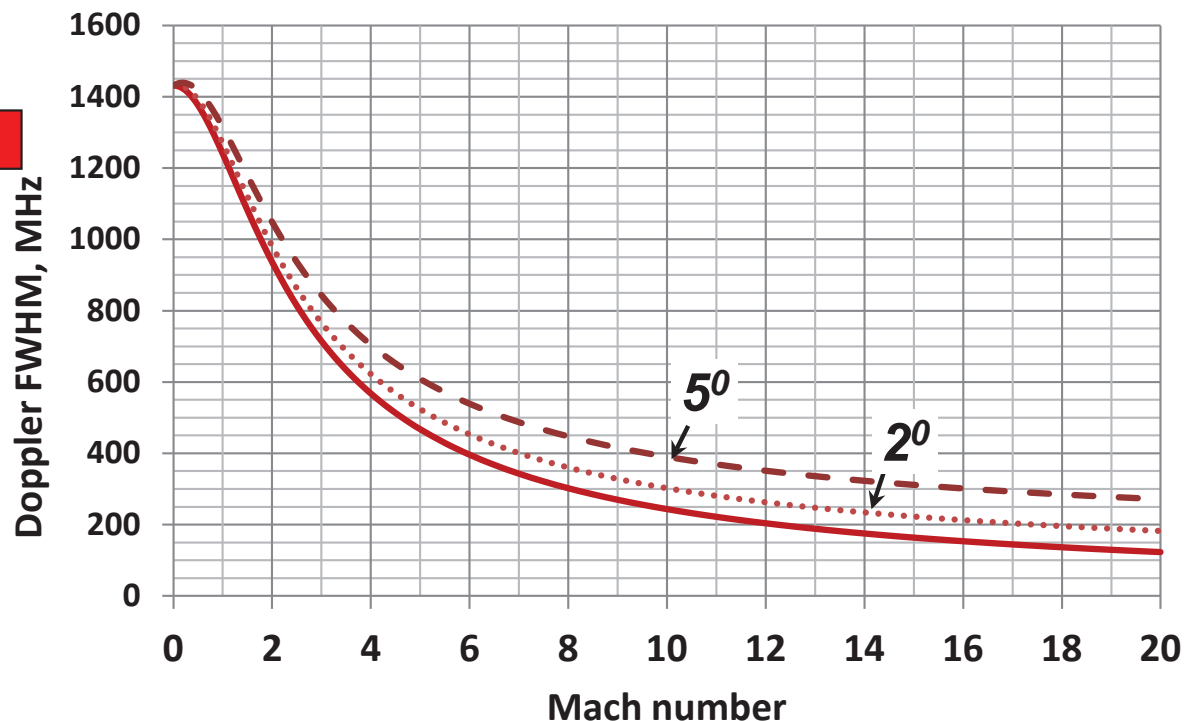
1. **Contra-propagating direction of laser- and supersonic beams makes conditions even worth in comparison co-propagating one**
2. **Applying blocking potential to the RF ion guide will lead to collection of laser-produced ions**
3. **No ionization in shock wave zones**

# Doppler broadening in supersonic beam

$4s^2S_{1/2} - 4p^2P_{1/2}$ , 327.4 nm  $^{63}\text{Cu}$  transition,  $V_0 = 30535.3 \text{ cm}^{-1}$



$$\Delta \nu_{\text{Doppler}} = 2\sqrt{\ln 2} \frac{v_0}{c} \sqrt{\frac{2kT}{m}} \quad \text{- FWHM}$$



$$\frac{A}{A^*} = \frac{1}{M} \left[ \left( \frac{2}{\gamma+1} \right) \left( 1 + \frac{\gamma-1}{2} M^2 \right) \right]^{(\gamma+1)/2(\gamma-1)}$$

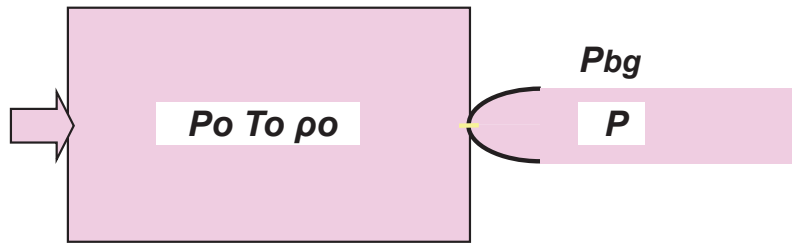
$$A^* = \pi d^2 / 4$$

$$d = 1 \text{ mm}$$

$$M=7 \text{ Diameter} - 4.88 \text{ mm}$$

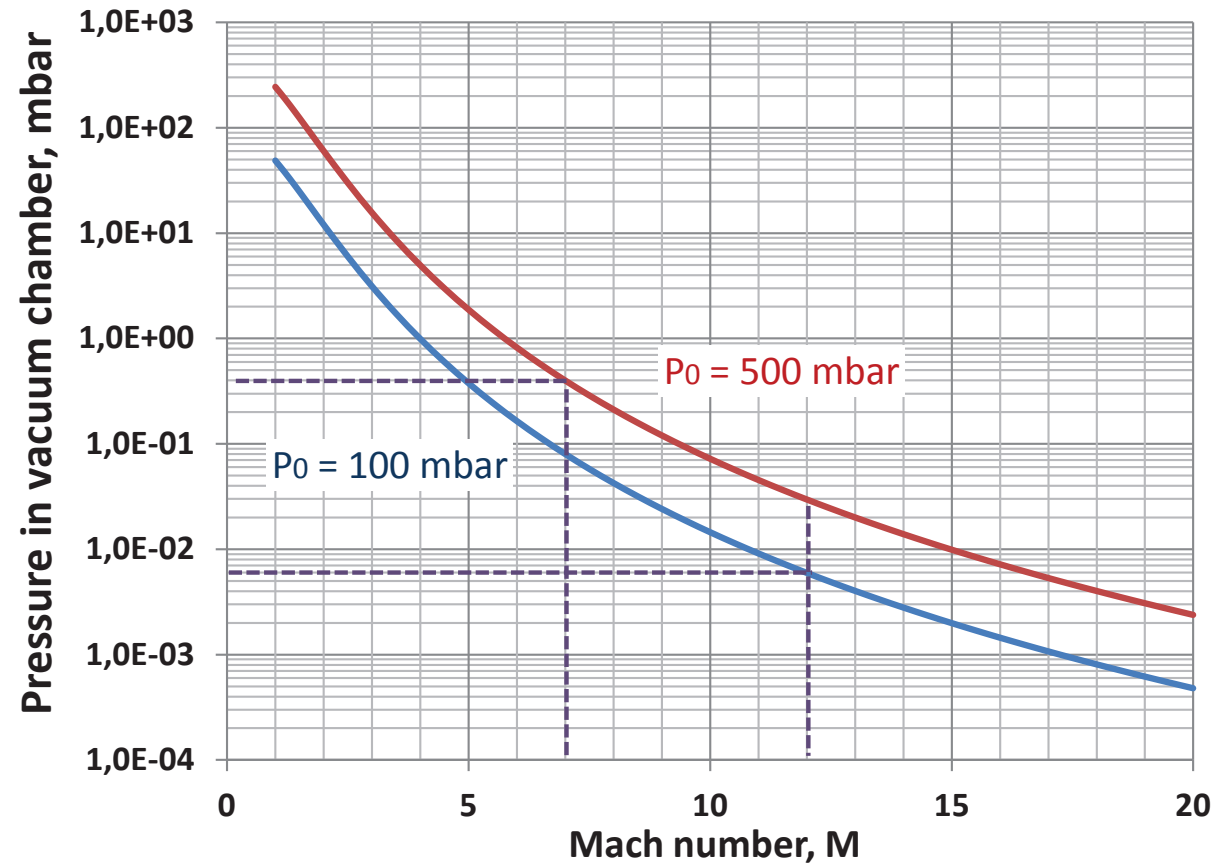
$$M=12 \text{ Diameter} - 10.51 \text{ mm}$$

# Required pumping capacity for the vacuum system (I)



$$P_{bg} = P \quad !$$

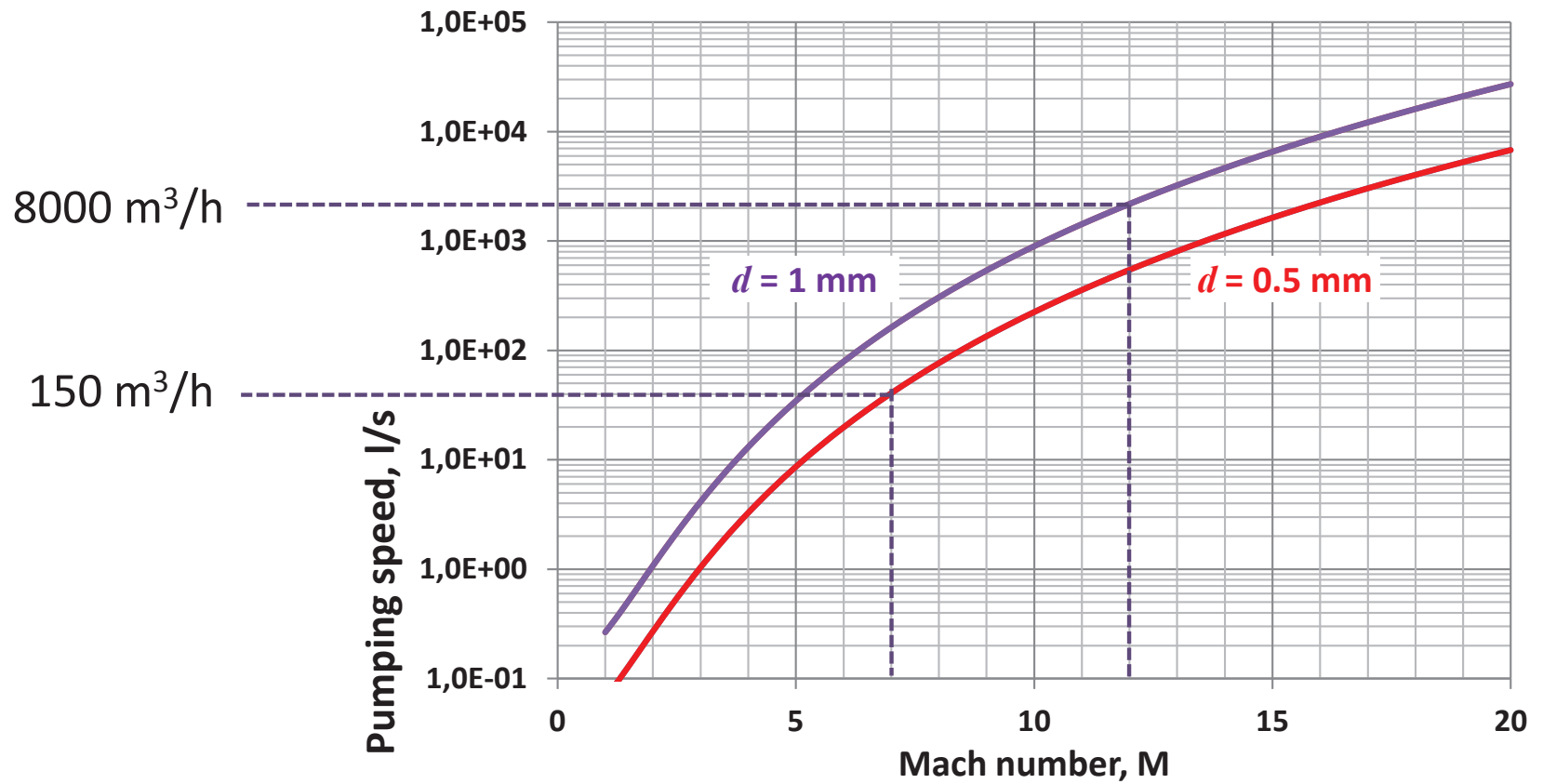
$$\frac{P}{P_0} = \left[ 1 + \left( \frac{\gamma - 1}{2} \right) M^2 \right]^{\frac{-\gamma}{\gamma - 1}}$$



# Required pumping capacity for the vacuum system (II)

$Q$  - the volume flow rate of the buffer gas through the throat diameter  $d$  (in l/s)  $Q = 0.052d^2 \sqrt{\frac{T_0}{A}}$   
 $A$  - in atomic mass units

$W$  - pumping speed of the vacuum system  $W = \frac{QP_0}{P_{bg}}$   $P_{bg} = P$



# Formation of clusters and chemical reactions

Hagena parameter  $G^*$

$$G^* = \eta \frac{(d)^{0.85}}{T_0^{2.29}} P_0$$

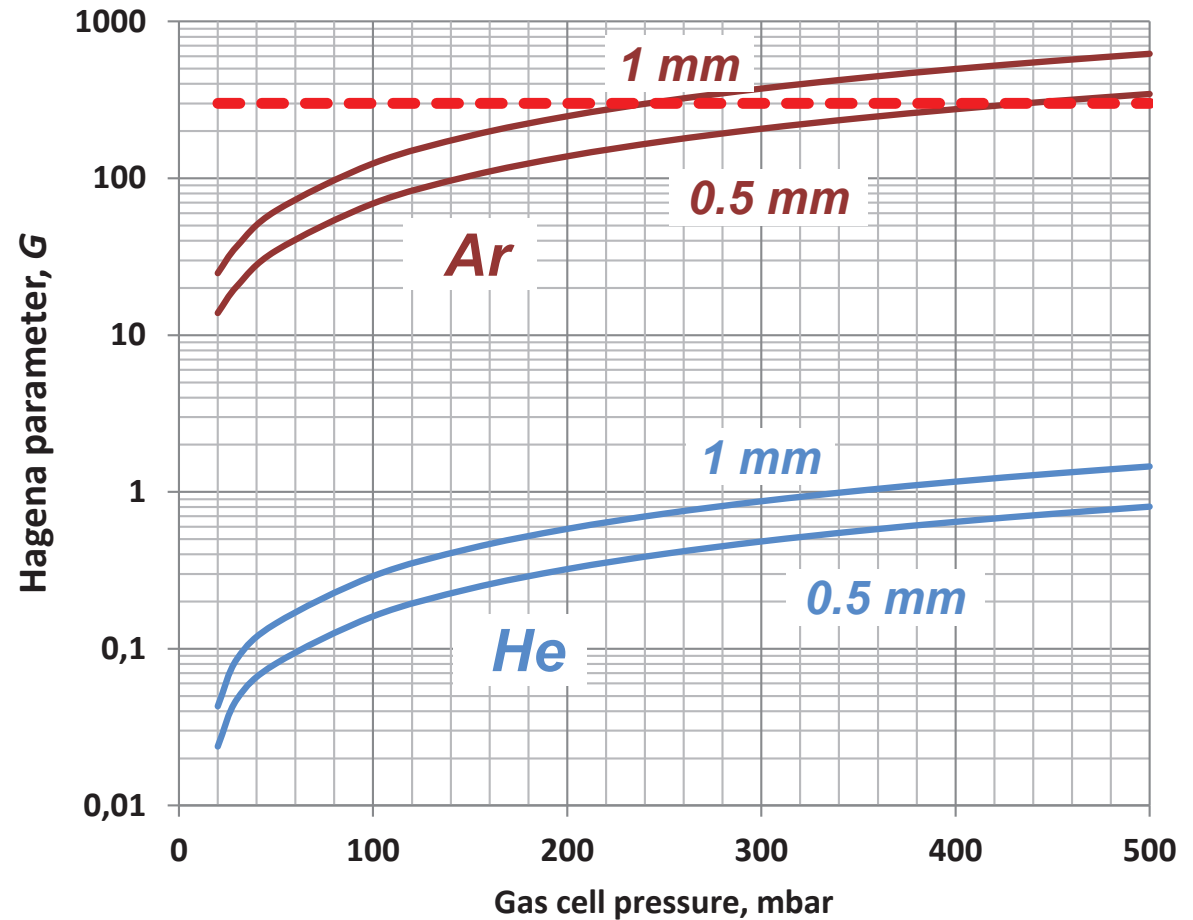
$\eta$  is the condensation parameter related to the bond formation.

number of atoms per clusters

$$N_c \sim G^{*2.0-2.5}$$

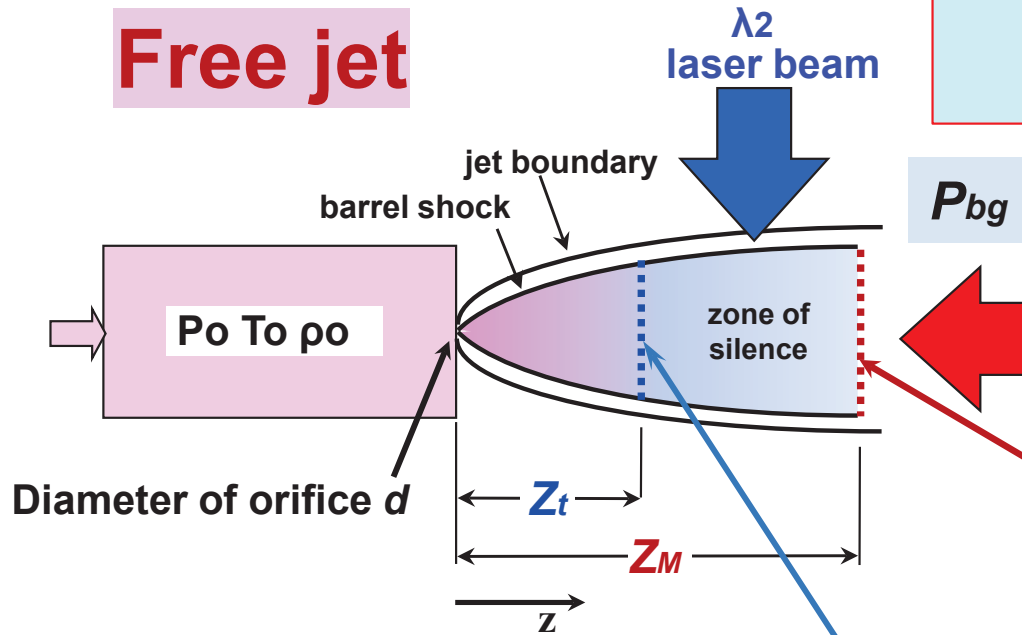
$$N_c \sim T_0^{-5}$$

Clustering starts if the Hagena parameter  $G^* > 300$



# Two-Step Laser Ionization in a Free Jet

## Free jet



1951 free jet – A. Kantrowitz, J. Grey

Mach disk,  $T, \rho \uparrow$   
 $Z_M$  – position of the Mach disk
 
$$\frac{Z_M}{d} = 0.67 \sqrt{\frac{P_0}{P_{bg}}}$$

$M_t$  - terminal Mach number

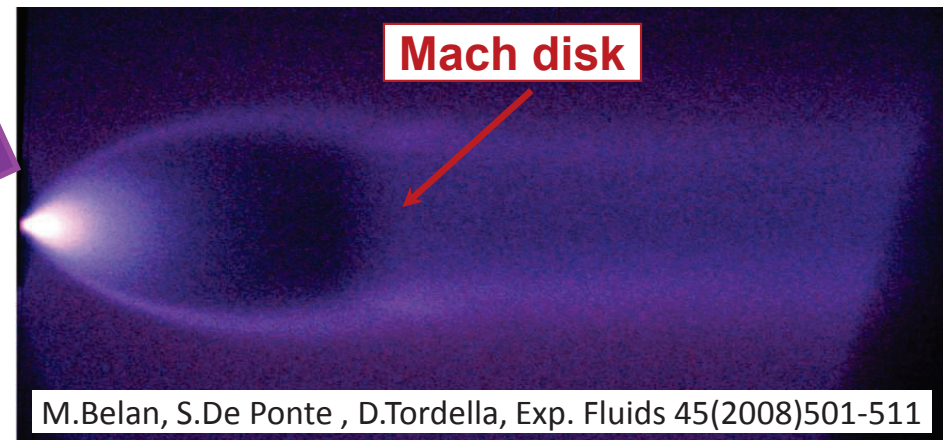
$Z_t$  – position of terminal Mach number

$$M_t = 3.32 (P_0 d)^{0.4}$$

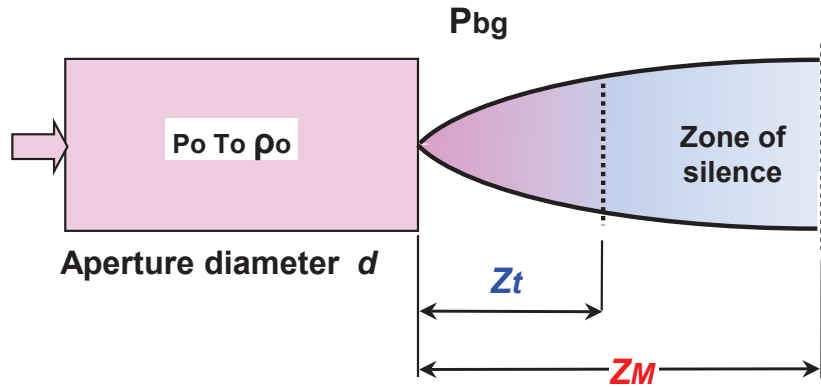
(mbar, mm)

$$\frac{Z_t}{d} = \left( \frac{M_t}{3.26} \right)^{1.5}$$

## Visualization of free jet

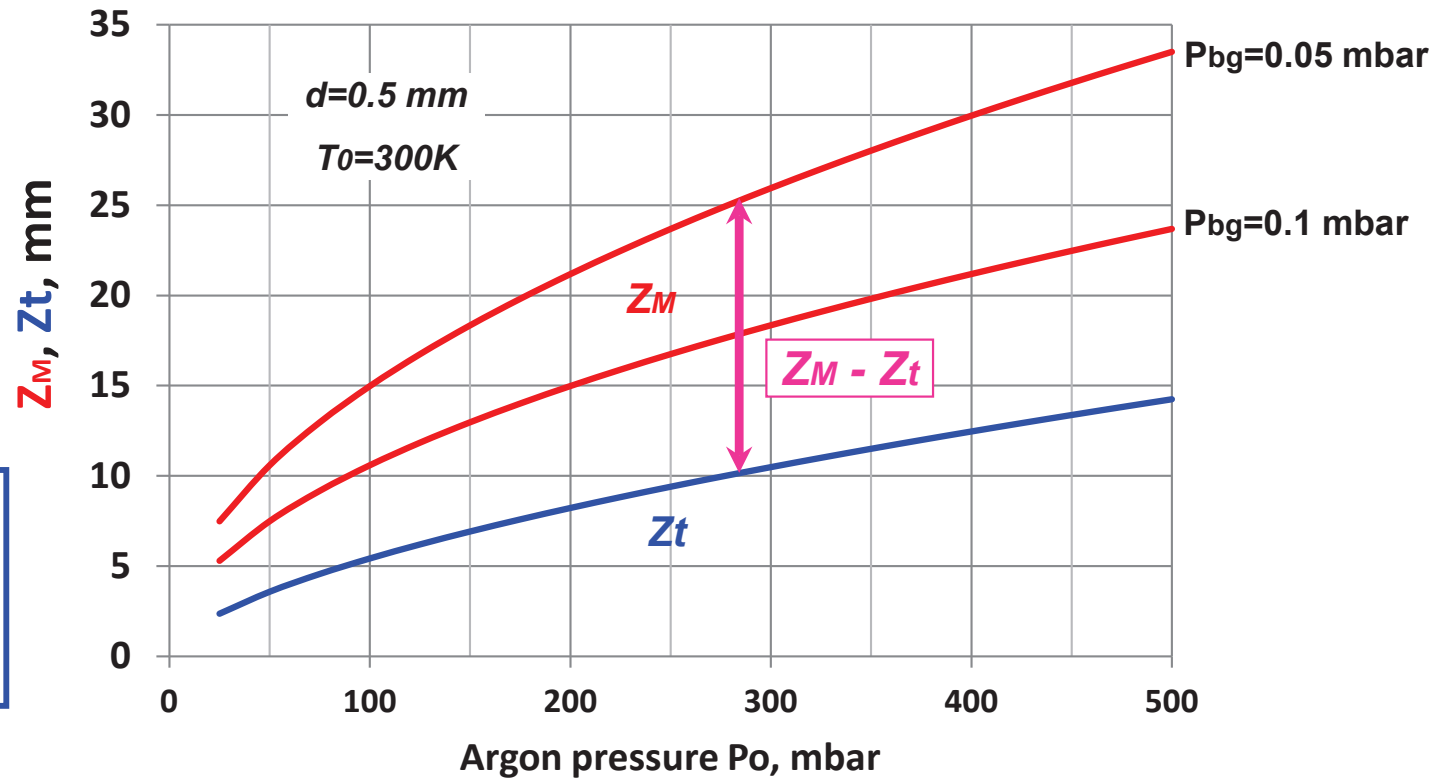


# Properties of free jet (I)



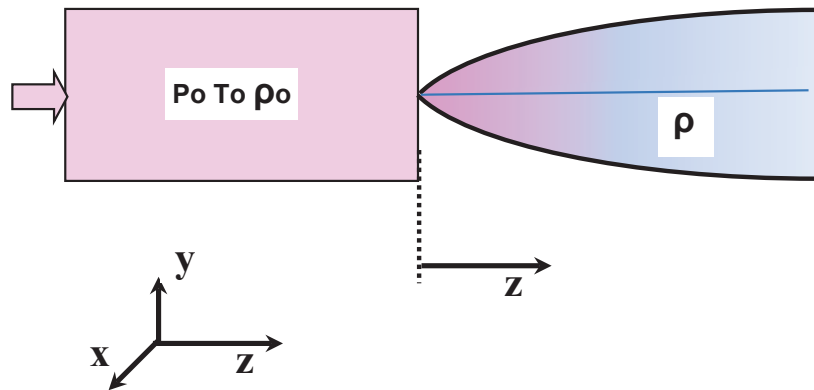
$$\frac{Z_M}{d} = 0.67 \sqrt{\frac{P_0}{P_{bg}}}$$

$$\frac{Z_t}{d} = \left[ \frac{133(P_0 d)^{0.4}}{3.26} \right]^{1.5}$$



Diameter of the Mach disk –  $0.5 \cdot Z_M$

# Properties of Free Jet (II)

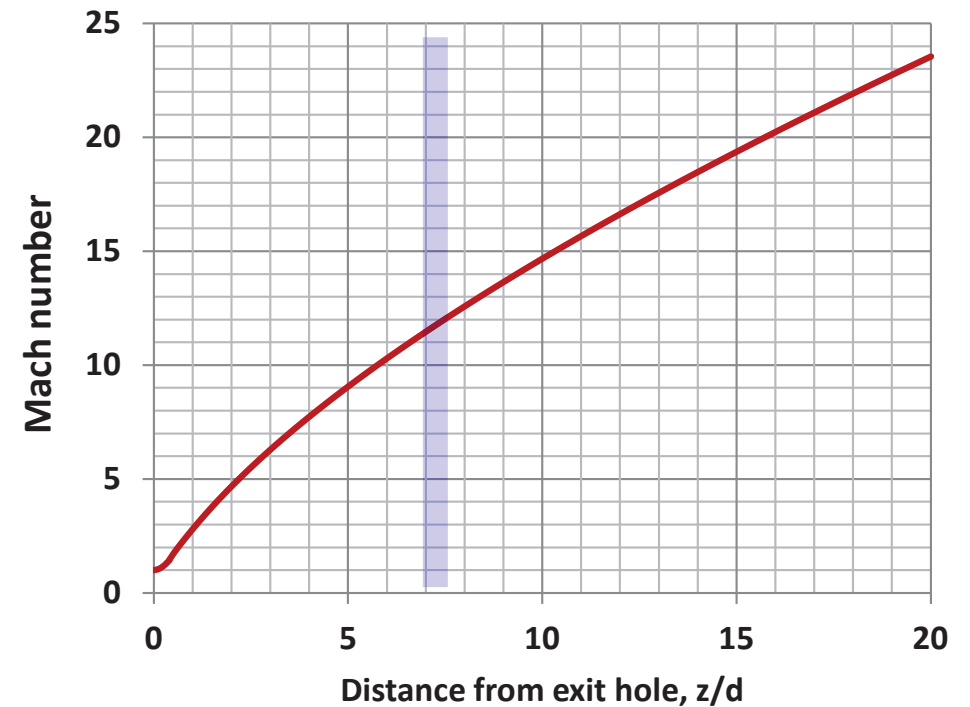
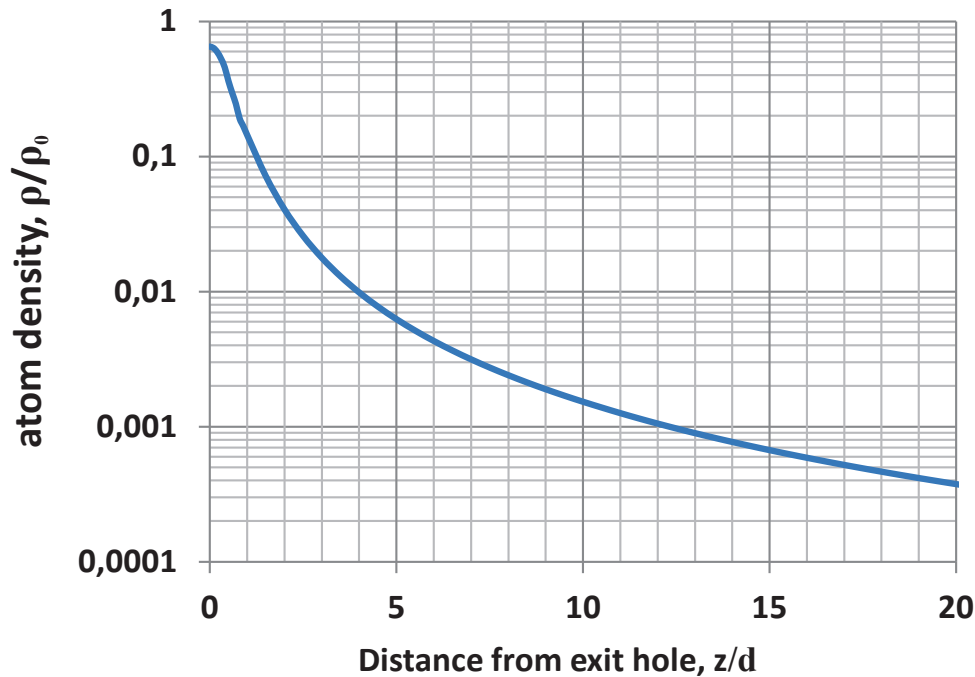


Centerline Mach number correlations

$A$	$B$	$C_1$	$C_2$	$C_3$	$C_4$
3.337	-1.541	3.232	-0.7563	0.3937	-0.0729

$$0 < \frac{Z}{d} < 1.0 \quad M = 1.0 + A \left( \frac{Z}{d} \right)^{-2} + B \left( \frac{Z}{d} \right)^3$$

$$\frac{Z}{d} > 0.5 \quad M = \left( \frac{Z}{d} \right)^{(\gamma-1)} \left[ C_1 + \frac{C_2}{\left( \frac{Z}{d} \right)} + \frac{C_3}{\left( \frac{Z}{d} \right)^2} + \frac{C_4}{\left( \frac{Z}{d} \right)^3} \right]$$



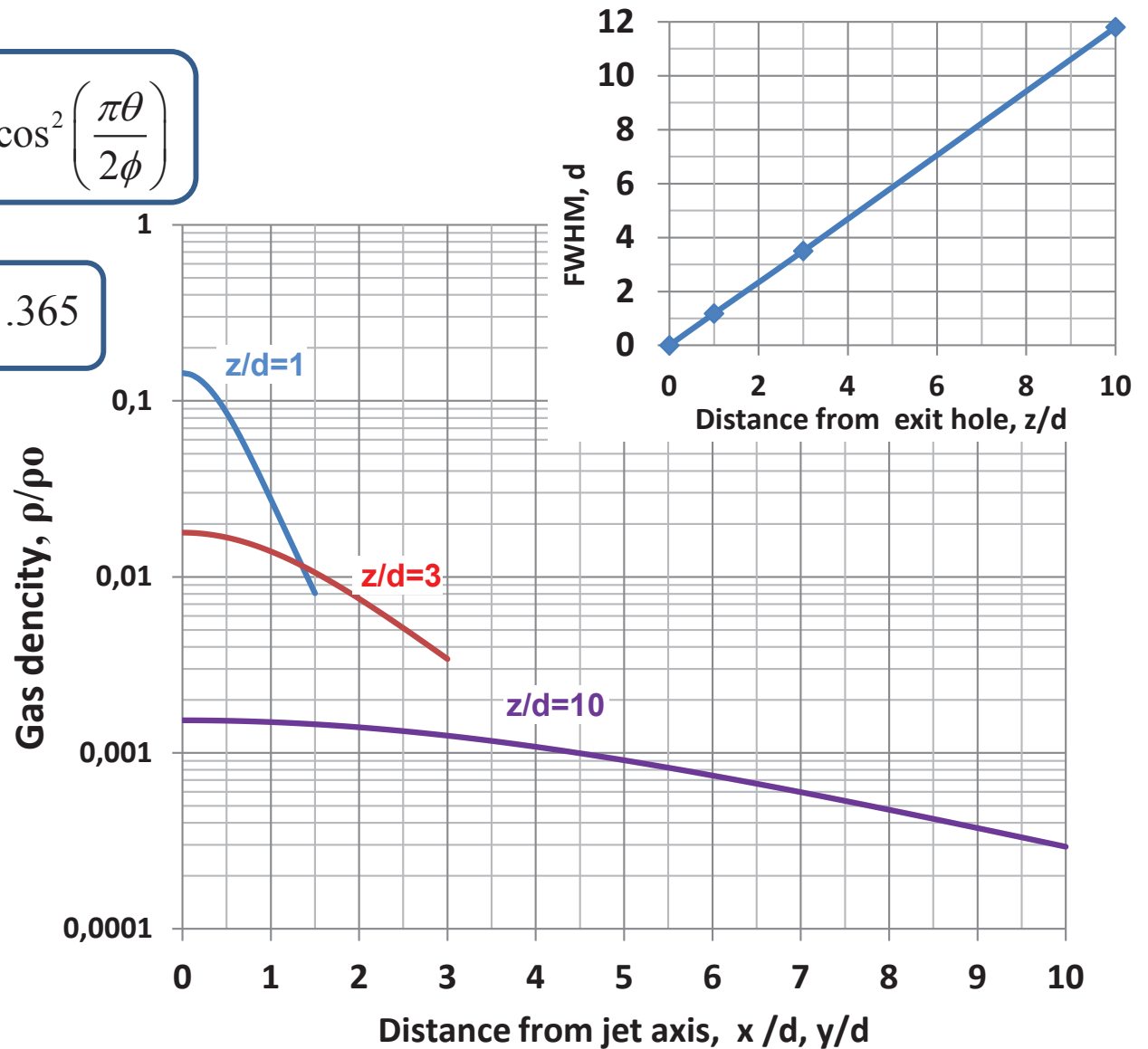
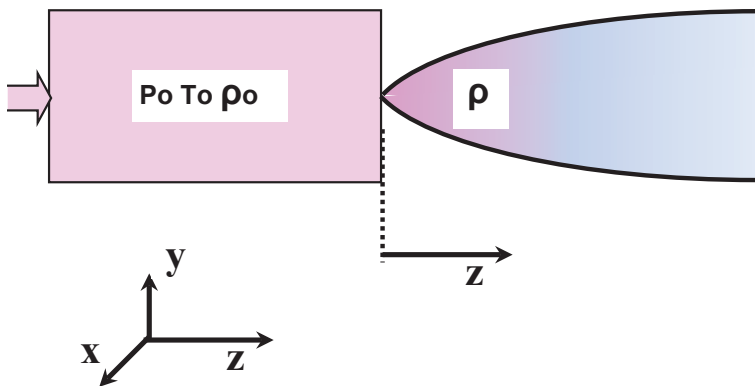


# Properties of Free jet (III)

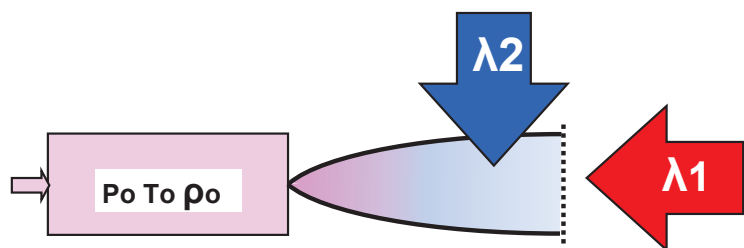
## Off-axis density distribution for axisymmetric free jet flow

$$\frac{\rho(y,z)}{\rho(0,z)} = \cos^2 \theta \cos^2 \left( \frac{\pi\theta}{2\phi} \right) \quad \frac{\rho(R,\theta)}{\rho(R,0)} = \cos^2 \left( \frac{\pi\theta}{2\phi} \right)$$

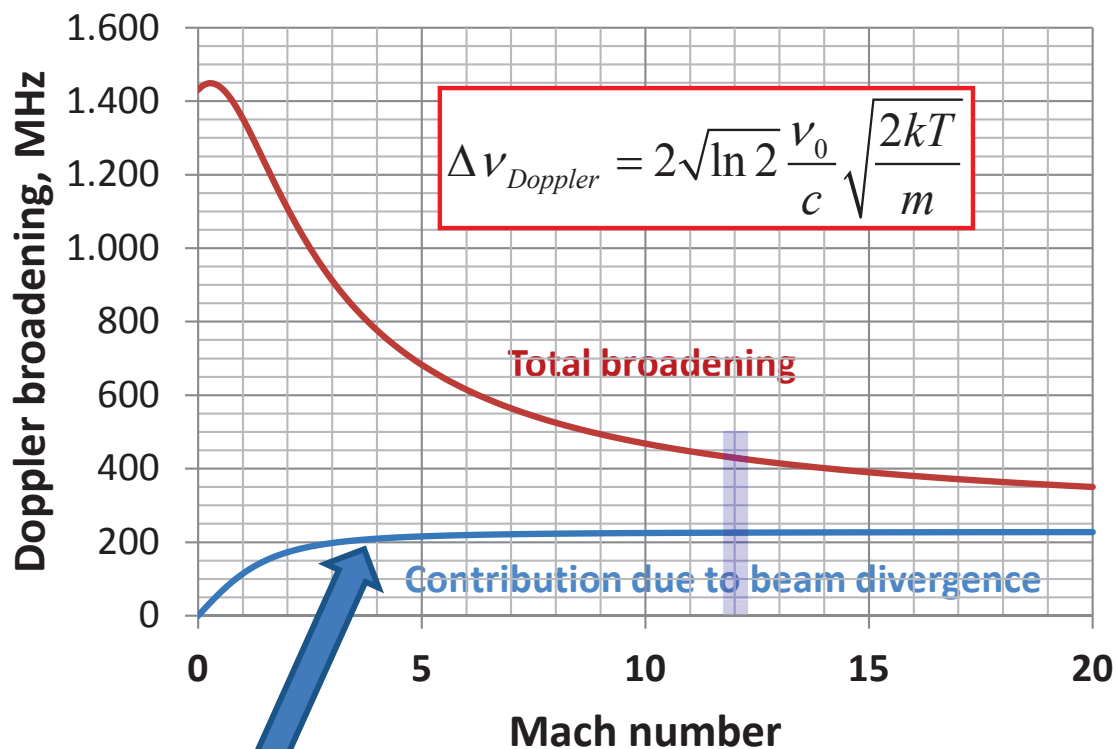
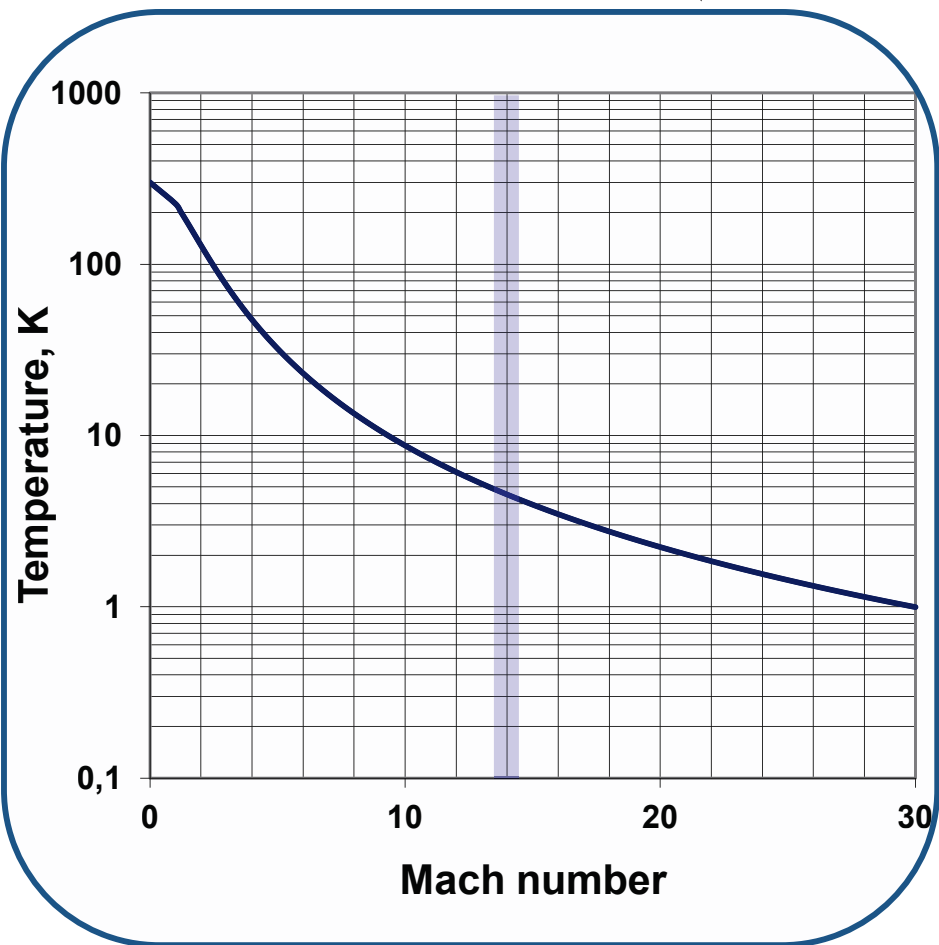
$$\tan \theta = \frac{y}{z} \quad R^2 = z^2 + y^2 \quad \gamma = 5/3 \quad \phi = 1.365$$



# Doppler Broadening in the Free Jet Supersonic Beam



$4s^2S_{1/2} - 4p^2P_{1/2}$ , 327.4 nm  $^{63}\text{Cu}$  transition,  $v_0 = 30535.3 \text{ cm}^{-1}$



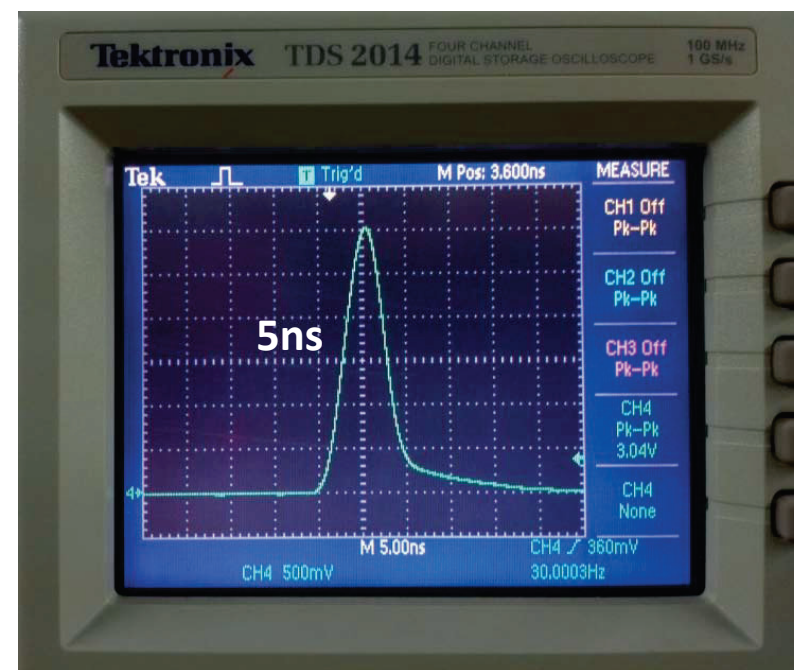
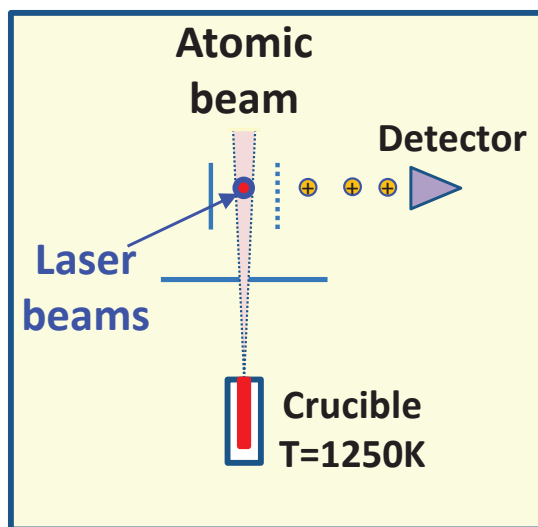
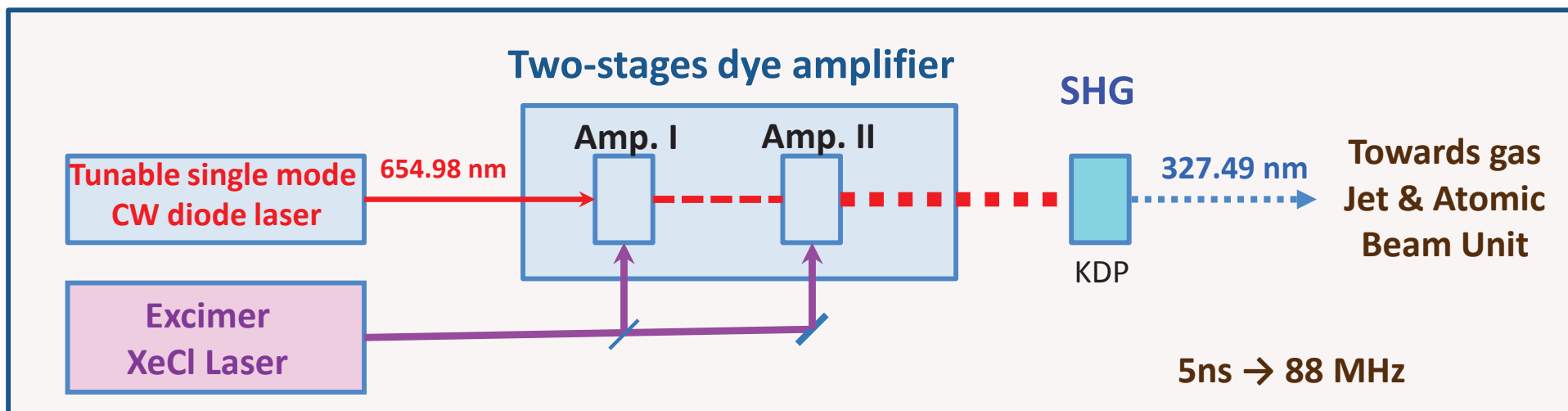
$$\Delta \nu_{Doppler} = 2\sqrt{\ln 2} \frac{v_0}{c} \sqrt{\frac{2kT}{m}}$$

$$\Delta_{Doppler}^{ax} = v_0 \cdot u (1 - \cos \theta) / c - \text{axial laser beam direction}$$

T=4K, Dopp. W.=200 MHz

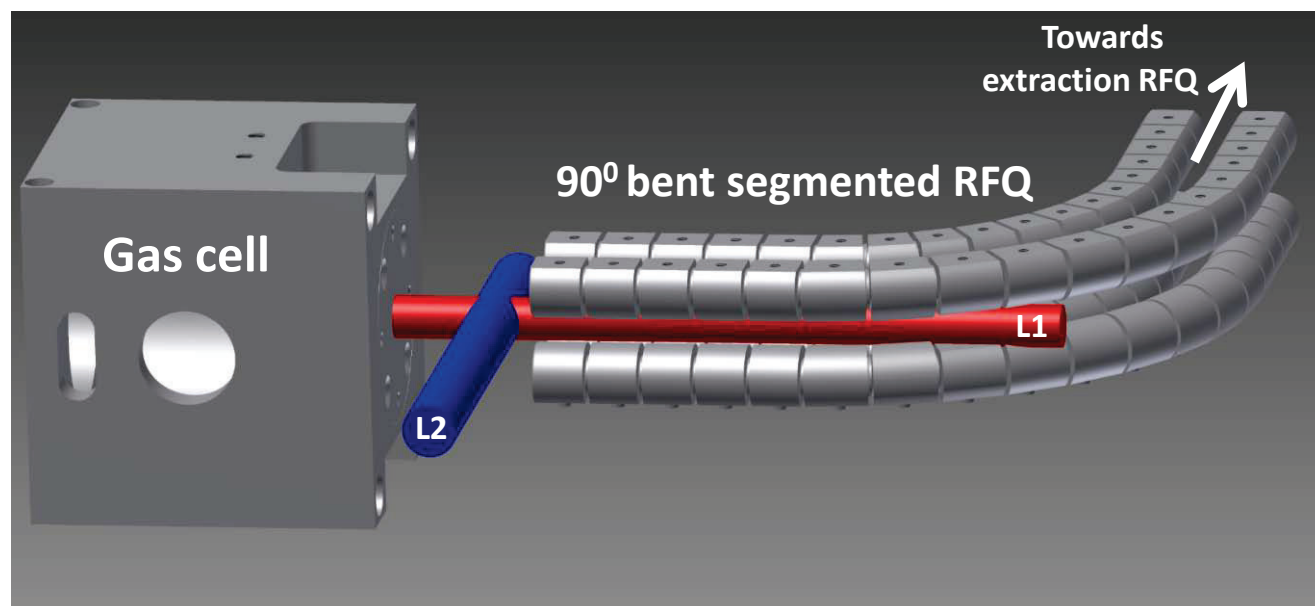
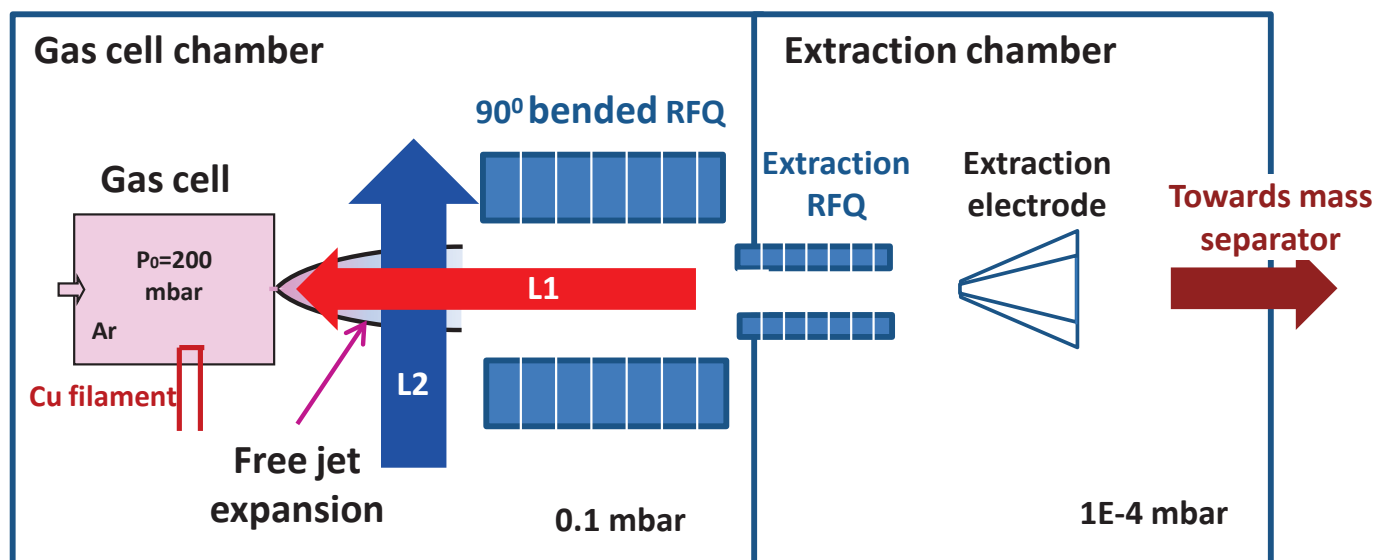
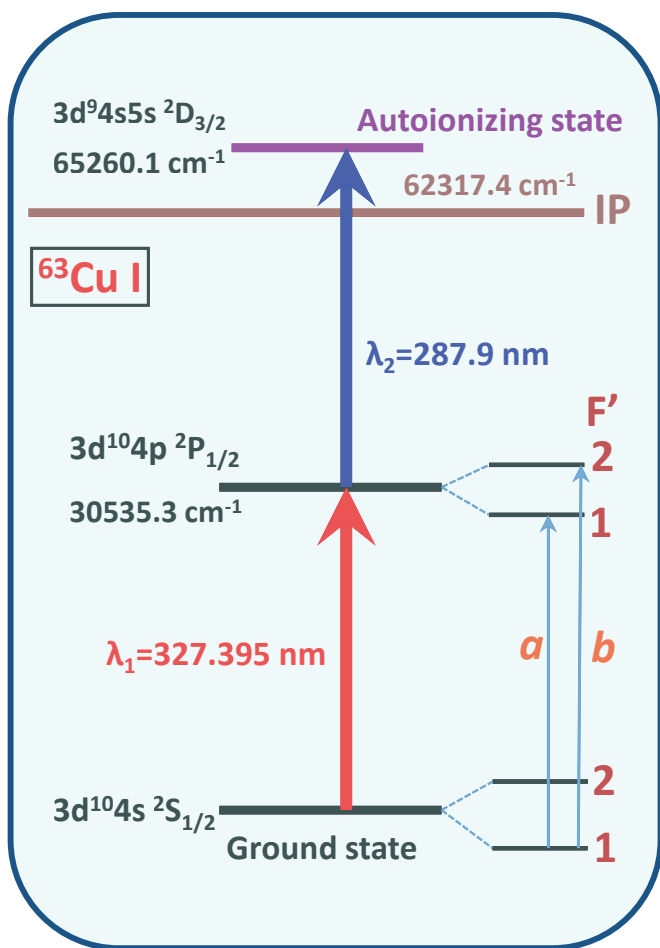
Tot. broad. = 420 MHz

# Amplification of CW Single Mode Diode Laser Radiation in a Pulsed Dye Amplifier



# Resonance Ionization Spectroscopy in a Free Gas Jet

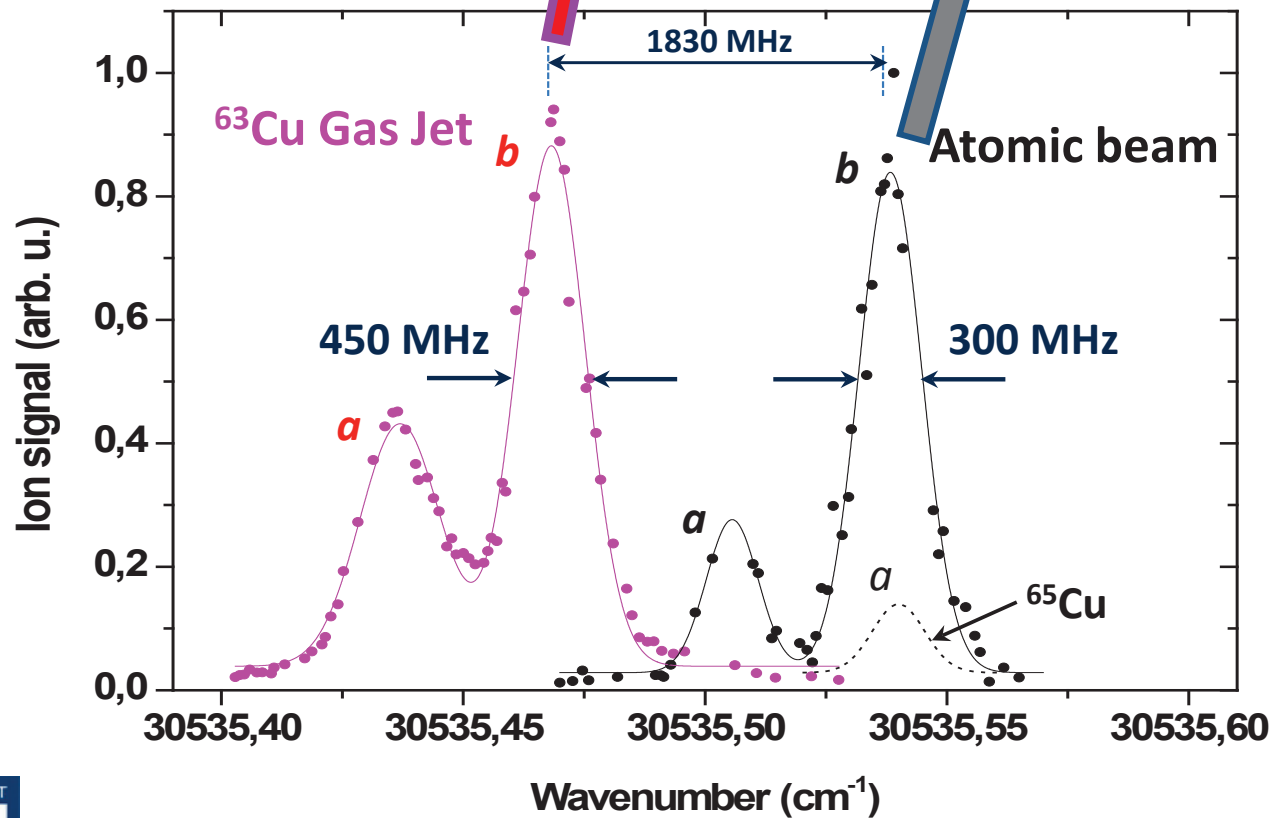
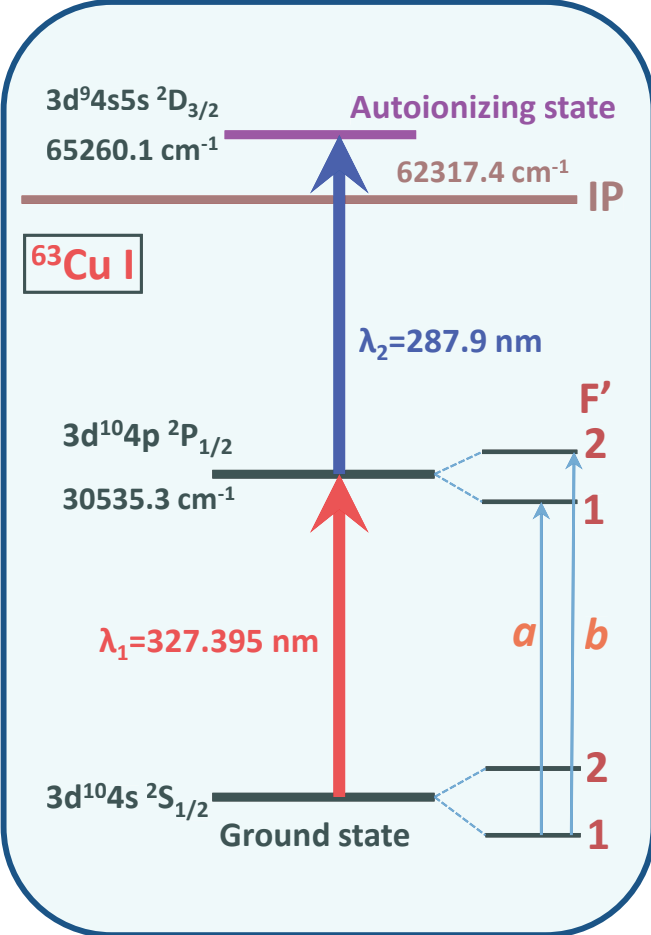
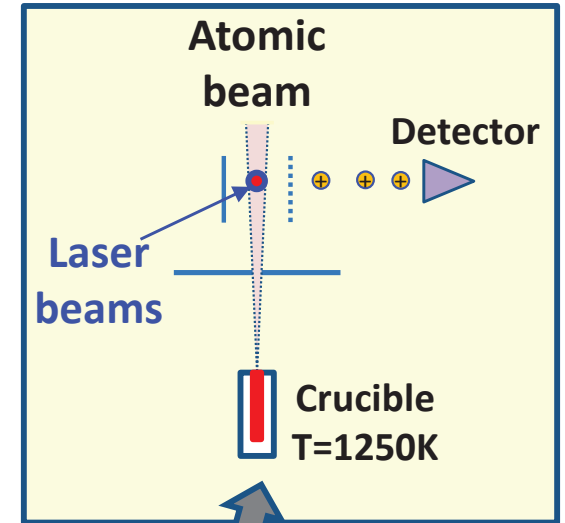
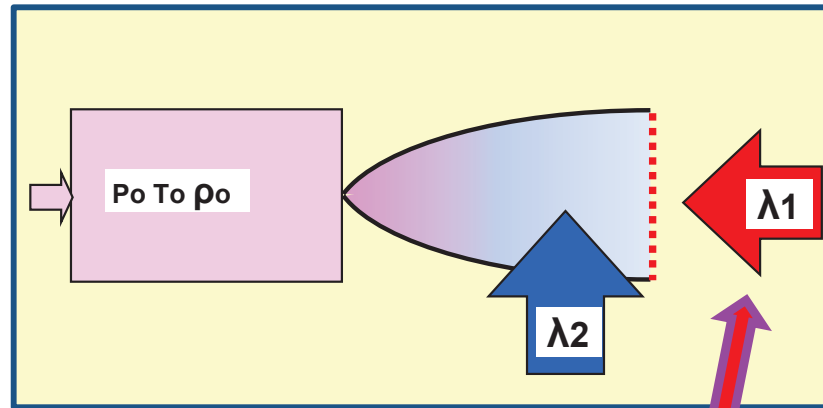
Accepted NIMB, Yu. Kudryavtsev et al, <http://arxiv.org/abs/1211.6649>



# Resonance Ionization Spectroscopy in a Free Gas Jet

$$u = \sqrt{\frac{kT_0 \gamma M^2}{m_{ng} \left\{ 1 + \left[ \frac{\gamma - 1}{2} \right] M^2 \right\}}}$$

1830 MHz  $\rightarrow$   $T_0 = 355 \pm 3$  K



# Summary

1. The crossed laser beams with supersonic jet has been proposed and realized at off-line conditions for two-step photo ionization in a free jet.
2. Using this method, the spectral resolution can be improved by more than one order of magnitude (200 MHz,  $\Delta\nu/\nu = 2.3\text{E-}7$ ) in comparison to the gas cell.

Thank you for your attention