

中性子ハロー核におけるテンソル 相関の役割とその効果

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Collaborators

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Contents

- Tensor correlation (TC) in light nuclei.
- He isotopes
 - ^4He : Tensor-optimized shell model to describe TC.
 - ^5He : $^4\text{He}+n$ phase shifts, LS splitting.
 - ^6He : Coupled $^4\text{He}+n+n$ model, Tensor suppression in 0_2^+ .
- Li isotopes
 - ^9Li : Tensor(**pn**) and Pairing(**nn**) correlations.
 - ^{11}Li : Coupled $^9\text{Li}+n+n$ model.
Tensor suppression for the halo formation.

Motivation and purpose of our studies on tensor correlation (TC)

- The tensor force is an essential component in nuclear force and plays a significant role in nuclear structure.
- Physical effect of the tensor force (V_T) is not easy to understand from the various kinds of ab initio calculations in a transparent manner.
- In the nuclear models such as mean field model and cluster model, V_T has been renormalized into V_C , V_{LS} and V_ρ , based on the **Brueckner-Theory**.
 - ${}^4\text{He} : (0s)^4 \rightarrow \langle V_T \rangle = 0$.
- We would like to understand the physical aspects of the tensor force on the nuclear structure **by describing TC explicitly in the model space**.

○ Recently, various kinds of ab initio calculations are available for light mass nuclei.

($A \simeq 10-12$)

– Green's Function Monte Carlo

R.B.Wiringa, S.C.Pieper, J.Carlson, V.R.Pandharipande, PRC62(2001)

– No Core Shell Model P.Navrátil, J.P.Vary, B.R.Barrett, PRL84(2000)

– Faddeev-Yakubovsky A.Nogga, H.Kamada, W.Glöckle, PRL85(2000)

– Variational method such as...

Stochastic variational method Y.Suzuki, K.Varga, PRC52(1995)

Gaussian expansion method E.Hiyama, Y.Kino, M.Kamimura, PPNP51(2003)

Hyperspherical Harmonics method M.Viviani, A.Kievsky, S.Rosati, FB-S18(1995)

– ATMS Y.Akaishi et al., Int.Review of Nucl.Phys.4(1986)

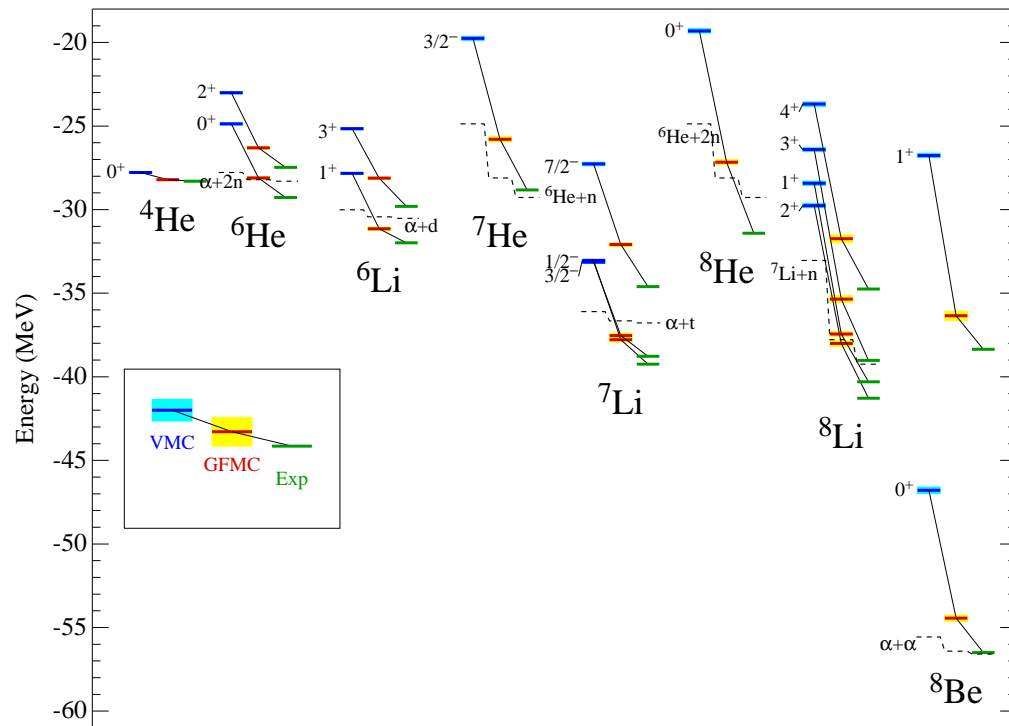
– Unitary operator approach

UMOA Fujii-Okamoto-Suzuki, PRC69(2004), UCOM Neff-Feldmeier, NPA713(2003).

– In ${}^4\text{He}$, $\langle V_{\text{tensor}} \rangle = -68.4$, $\langle V_{\text{central}} \rangle = -55.3$, $\langle T \rangle = 102.4$, unit in MeV

$P(D) = 13.9\%$. Kamada et al., PRC64('01)044001.

Variational calculation in real model space: Illinois-Argonne Group



- C. Pieper, R.B. Wiringa, Annu. Rev. Nucl. Part. Sci. 51 (2001)
- R.B. Wiringa, S.C. Pieper, J. Carlson, V.R. Pandaripande, PRC 62 (2000) 014001.

○ $H = T + V_{NN}$, $V_{NN} = V_{\pi} + V_{ST}^R$

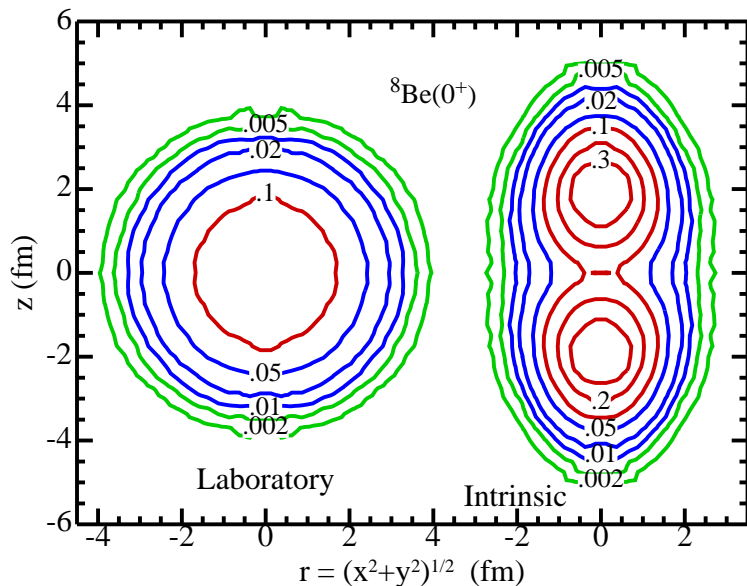
- AV18 Potential with three-body force.

- Variational Monte-Carlo and Green Function Monte-Carlo.

- Reproduction of binding properties and spectra.

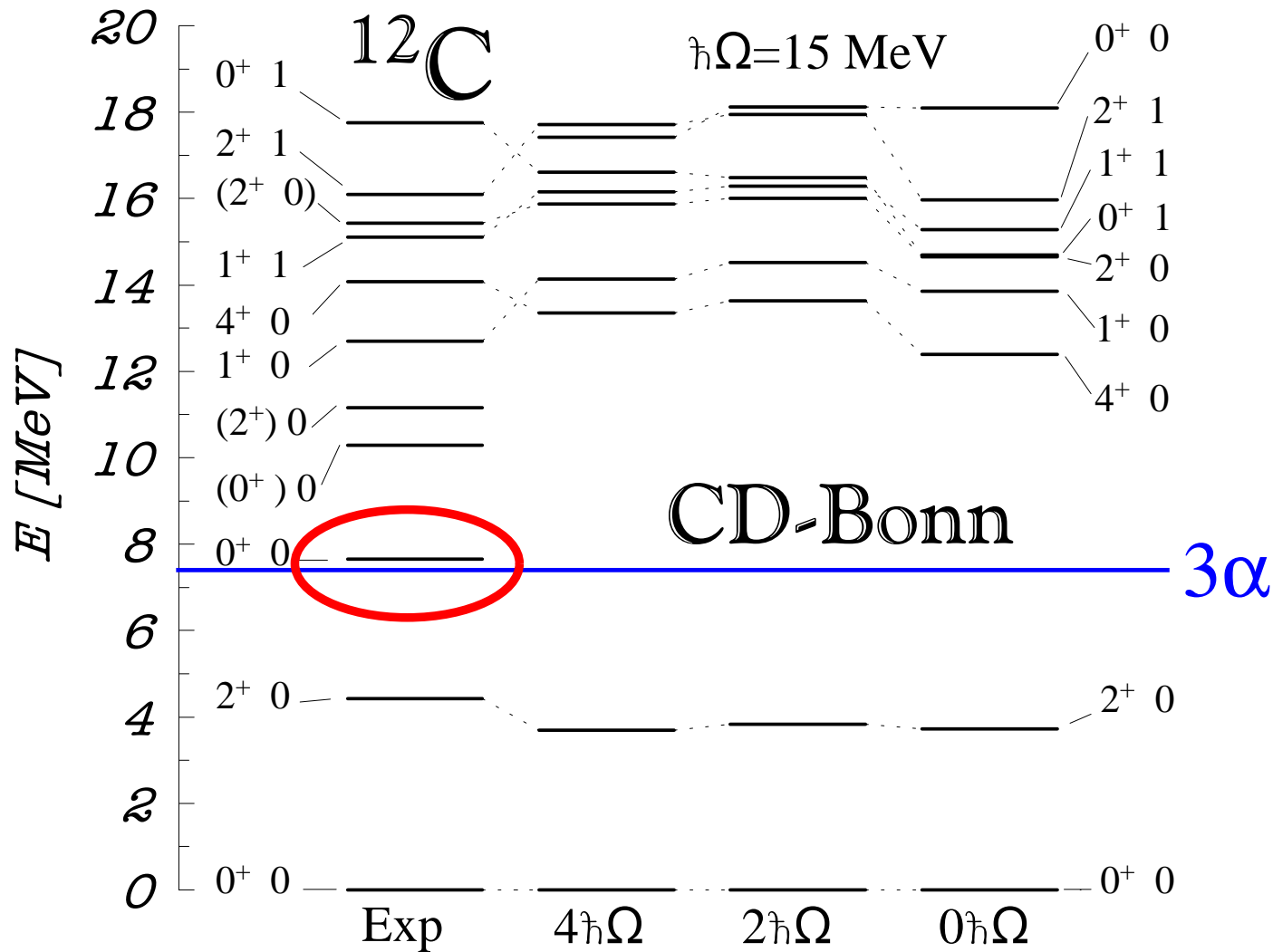
○ V_{π} is 70 – 80 % of the attraction.

- Cluster structure of ^8Be is reproduced.



No core-shell model calculation

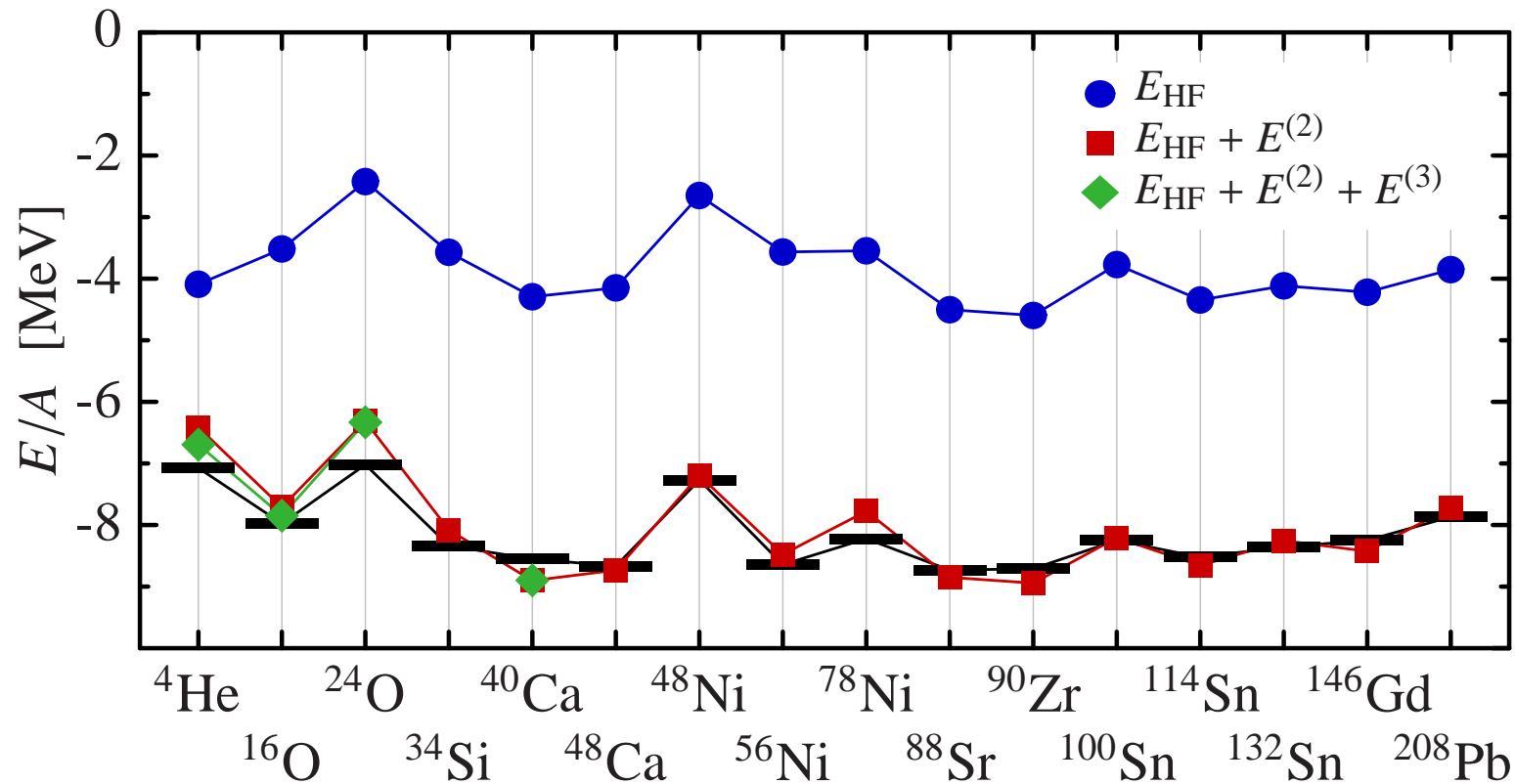
P.Navrátil, J.P.Vary, B.R.Barrett, PRL84(2000)



Unitary Correlation Operator Method (UCOM)

– $\Phi_{\text{corr.}} = C_{\omega} \cdot C_r \cdot \Phi_{\text{uncorr.}}$ ← HF, FMD, ...
 tensor and short-range correlator

– $C_r = \exp[-i \sum_{(i<j)} g_{ij}]$, $g = \frac{1}{2} [p_r s(r) + s(r) p_r]$



AV18

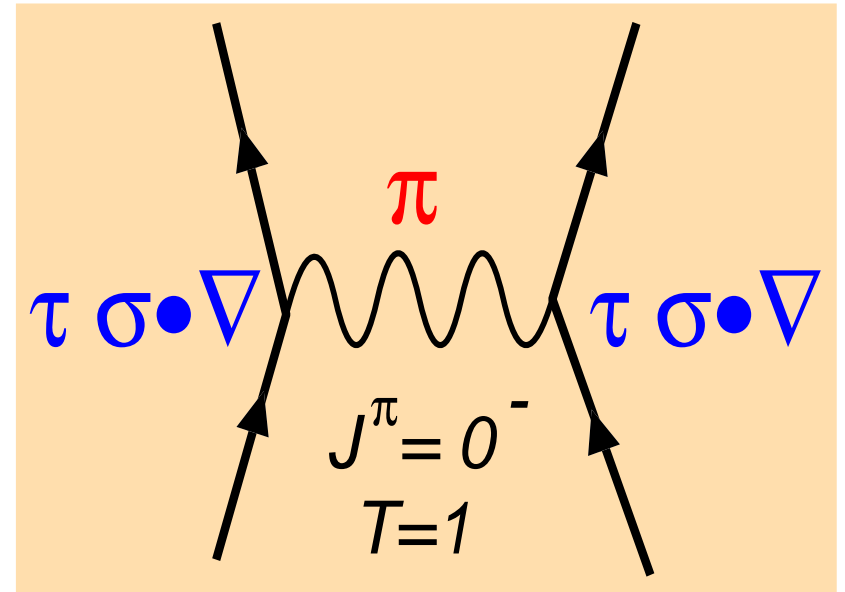
R. Roth, P. Papakonstantinou,

N. Paar, H. Hergert, T. Neff,

H. Feldmeier, PRC73(2006)

Mean Field approach

Tensor correlation in the mean field theory
by Sugimoto-Toki-Ikeda(NPA740('04))
and Akaishi(NPA738('04)).



- **CPPHF** : Charge- and Parity-Projected Hartree-Fock method
- Single particle orbit breaks the parity and charge.

$$- \phi_i = \phi_i^p + \phi_i^n \quad (p: \text{proton}, n: \text{neutron}),$$

$$- \phi_{j=1/2}^p = \phi_{0s1/2}^p + (\boldsymbol{\sigma} \cdot \mathbf{r}) \phi_{0s1/2}^p = \phi_{0s1/2}^p + \phi_{0p1/2}^p.$$

cf. Ogawa et al., CPP-RMF, PRC73(2006), Doté et al., CPP-AMD, PTP115(2006),
Isshiki-Naito-Ohnishi, AMD with coherent pion, PTP114(2005)

Our approach based on the extended shell model

- T.Myo, K.Katō, K. Ikeda, PTP113(2005) –
- We use the **extended shell model approach** to take into account TC explicitly, “**Tensor-optimized shell model**”, referring to the results of CPPHF.
 - **1p1h + 2p2h excitations**: Configuration mixing beyond the major shell.

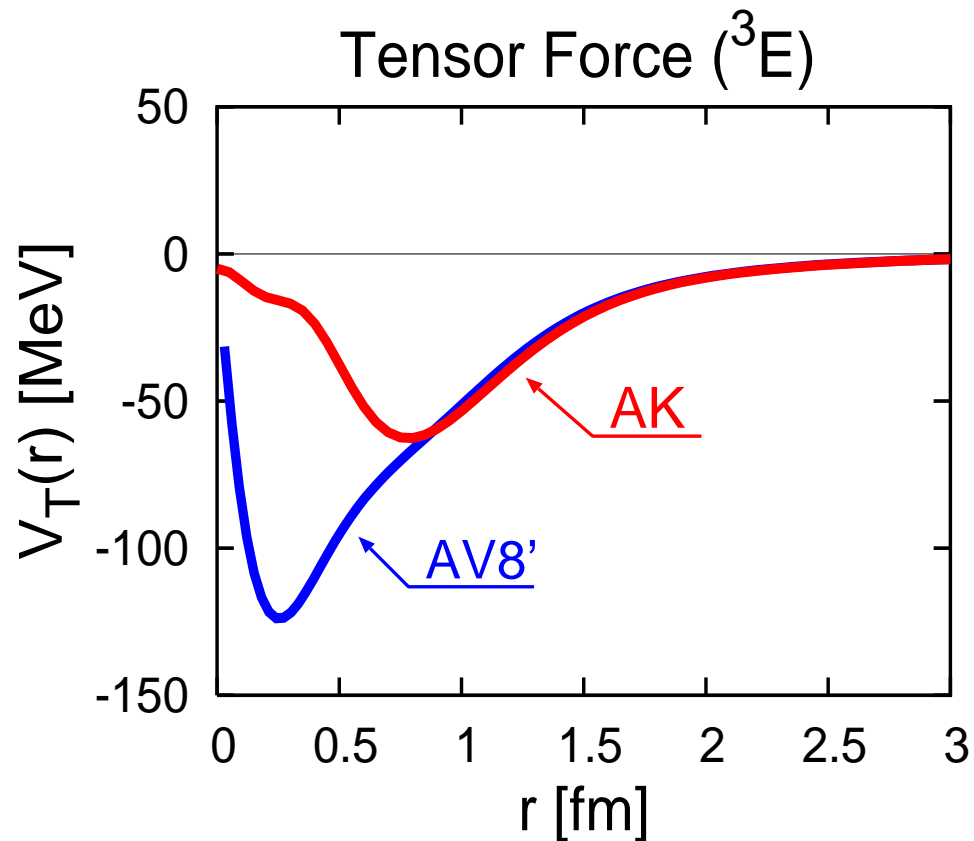
$$\Phi(^4\text{He}) = \sum_i C_i \psi_i(\{b_\alpha\}) = C_1 (0s)^4 + C_2 (0s)^2(\overline{0p}_{1/2})^2 + \dots$$

- HO basis with **independent length parameters** $\{b_\alpha\}$, such as $b_{0s} \neq b_{\overline{0p}}$, ... to represent the high momentum components. (cf. works by Sugimoto et al./ Akaishi)

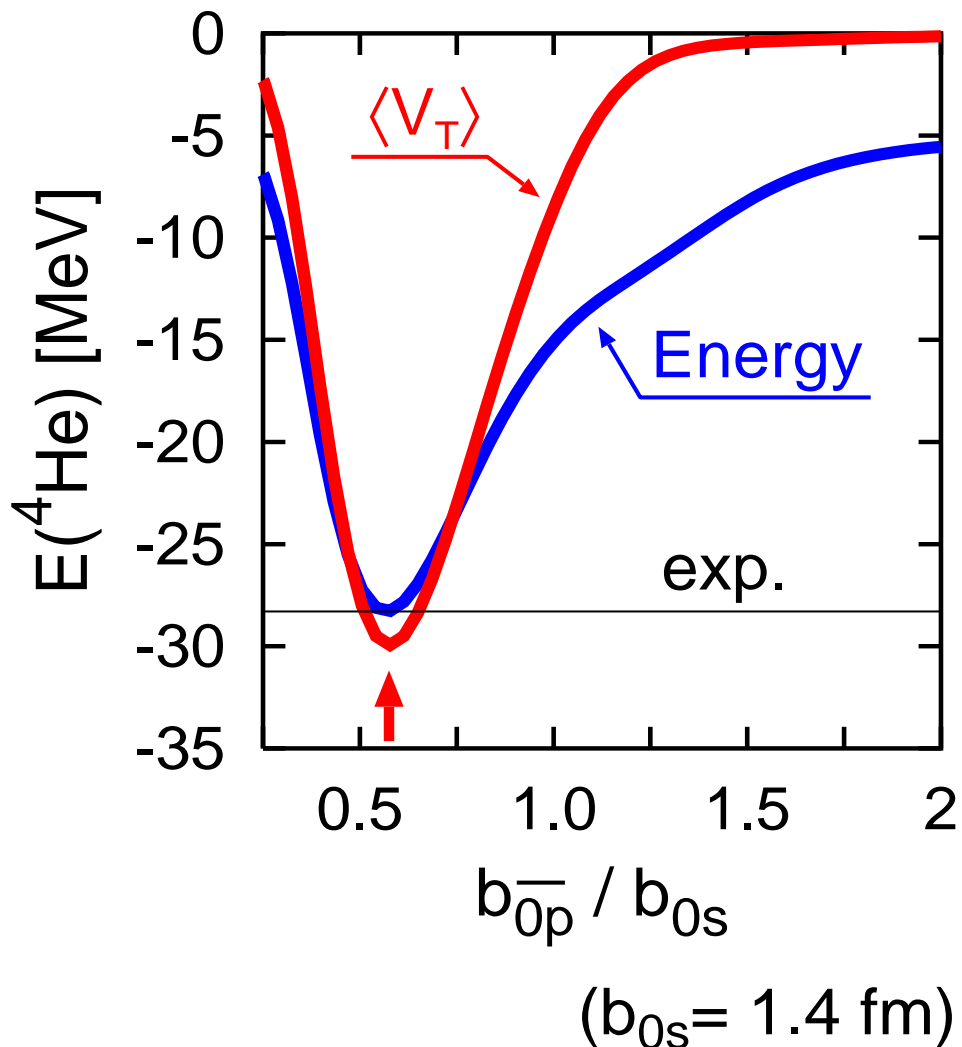
$$H = \sum_{i=1}^A t_i - T_G + \sum_{i < j}^A v_{ij}, \quad v_{ij} = v_{ij}^C + v_{ij}^T + v_{ij}^{LS} + v_{ij}^{Cmb}$$

$$\delta \frac{\langle \Phi | H | \Phi \rangle}{\langle \Phi | \Phi \rangle} = 0 \quad \Rightarrow \quad \frac{\partial \langle H - E \rangle}{\partial b_\alpha} = 0, \quad \frac{\partial \langle H - E \rangle}{\partial C_i} = 0.$$

- Interaction : Akaishi force (AK) (NPA738)
 - G-matrix using AV8' with $k_Q = 2.8 \text{ fm}^{-1}$ ($> k_F = 1.4 \text{ fm}^{-1}$)
 - \Rightarrow Long and intermidiate ranges of the tensor force survive.
 - Central part : We adjust the intermidiate range to fit B.E. and R_m of ${}^4\text{He}$.



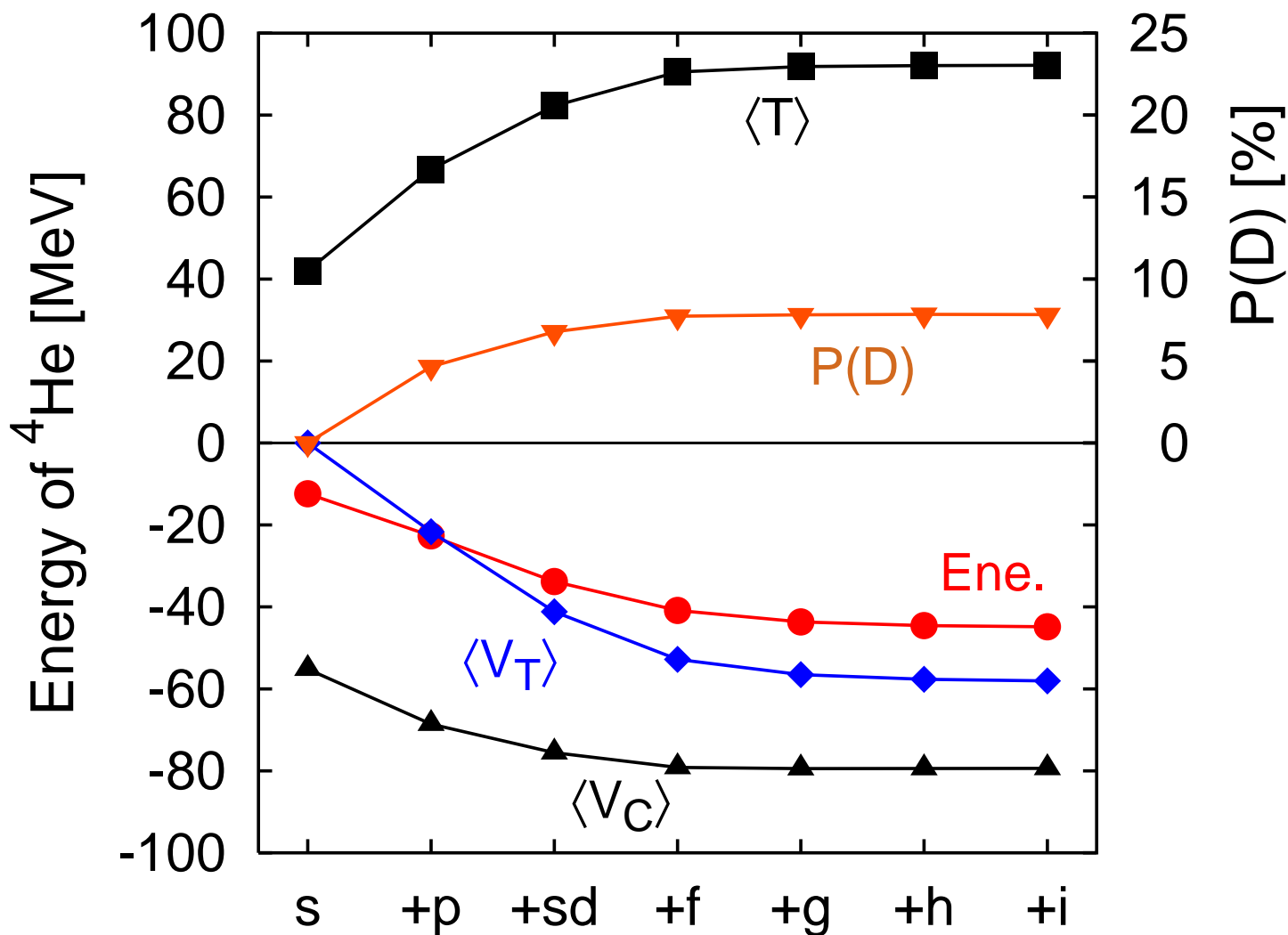
Previous results : ${}^4\text{He}$ with $0S+0P$ of HO bases, $V_T \times 1.5$



$(0s_{1/2})^4$	87.5 %
$(0s_{1/2})_{JT}^2 (\overline{0p}_{1/2})_{JT}^2$ (JT)=(10)	7.7 %
(JT)=(01)	1.3 %
$(0s_{1/2})^2 (\overline{0p}_{3/2})^2$	3.4 %
$(0s_{1/2})^2 (\overline{0p}_{1/2}) (\overline{0p}_{3/2})$	0.06 %
P[D]	6.9 %
R_m	1.48 fm

- narrow p-orbit \Rightarrow higher shell effect
- 0^- coupling of $0s_{1/2} - 0p_{1/2}$
 \Rightarrow pion nature
- (J,T)=(1,0) \Rightarrow deuteron correlation

^4He with HO basis by adding high angular momentum states.



○ Length parameters

orbit	\bar{b}_{nlj} / b_{0s}
$\bar{0}p_{1/2}$	0.65
$\bar{0}p_{3/2}$	0.58
$\bar{1}s_{1/2}$	0.63
$\bar{0}d_{3/2}$	0.58
$\bar{0}d_{5/2}$	0.53
$\bar{0}f_{5/2}$	0.66
$\bar{0}f_{7/2}$	0.55

○ The solutions show a good convergence.

\Rightarrow higher shell effect $\sim 16 \hbar\omega$

^4He with $0S+0P+SD+F+G$, 4 Gaussians

Energy [MeV]	-28.0
$\langle V_T \rangle$ [MeV]	-51.0
$(0s_{1/2})^4$	85.0 %
$(0s_{1/2})^2_{JT}(\overline{0p}_{1/2})^2_{JT} \quad (JT)=(10)$	5.0 %
$(0s_{1/2})^2_{JT}(\overline{0p}_{1/2})^2_{JT} \quad (JT)=(01)$	0.3 %
$(0s_{1/2})^2[(\overline{1s}_{1/2})(\overline{0d}_{3/2})]_{10}$	2.4 %
$(0s_{1/2})^2[(\overline{0p}_{3/2})(\overline{0f}_{5/2})]_{10}$	2.0 %
$(0s_{1/2})^2[(\overline{0p}_{1/2})(\overline{0p}_{3/2})]_{10}$	0.9 %
P[D]	9.6 %

$$\langle T \rangle = 71.2 \text{ [MeV]}$$

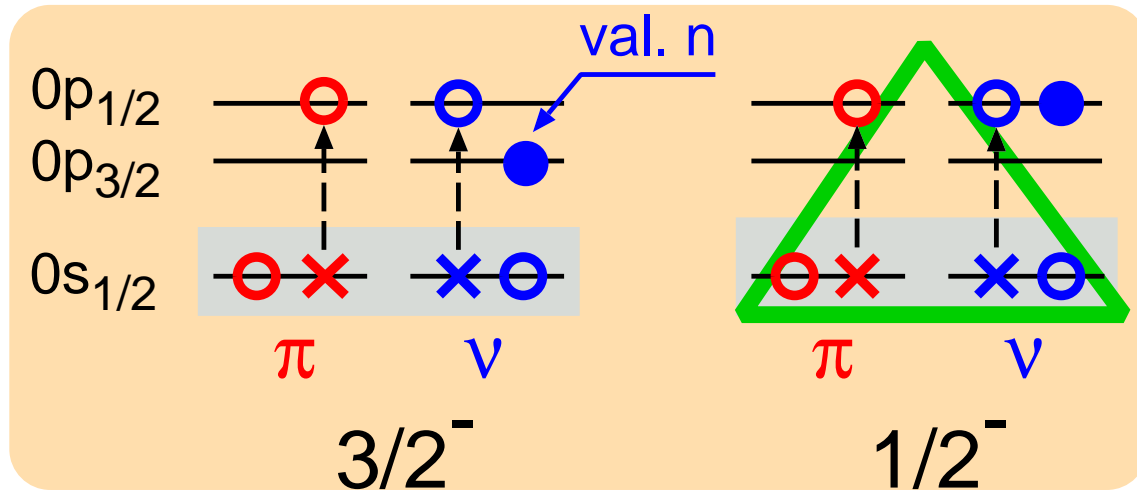
$$\langle V_C \rangle = -48.6 \text{ [MeV]}$$

c.m. excitation ~ 0.6 MeV.

\Leftarrow Three cases are selectively mixed.

- one shows 0^- pion nature
- deuteron correlation with $JT=(01)$
- $(\Delta L=\Delta S=2)$ properties of V_T .

LS splitting in ${}^5\text{He}$ in a coupled ${}^4\text{He}+n$ model



Tensor correlation is suppressed in $1/2^-$ due to the Pauli-blocking



An appreciable contribution to the doublet (LS) splitting

[Ref] T. Terasawa, PTP22('59), S. Nagata, T. Sasakawa, T. Sawada and R. Tamagaki, PTP22('59)
K. Ando and H. Bando PTP66('81)

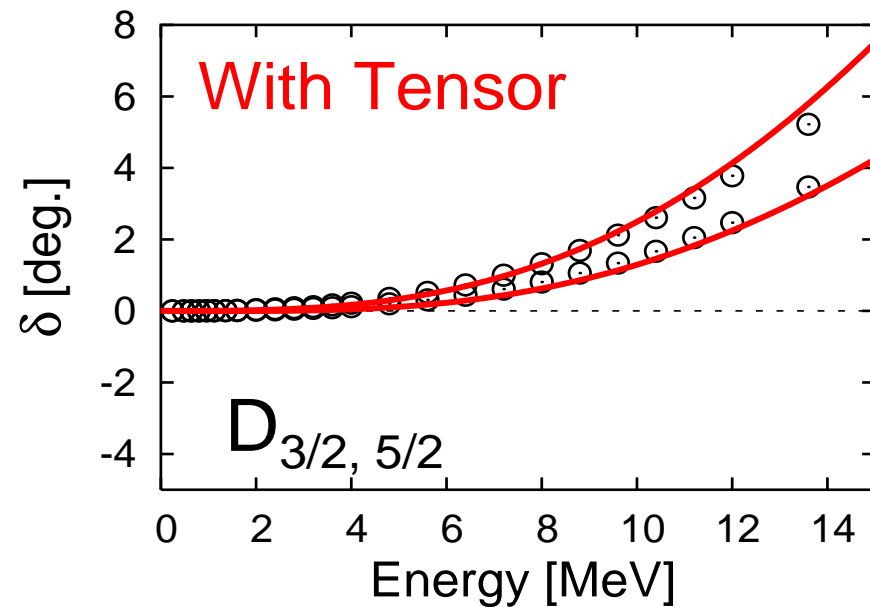
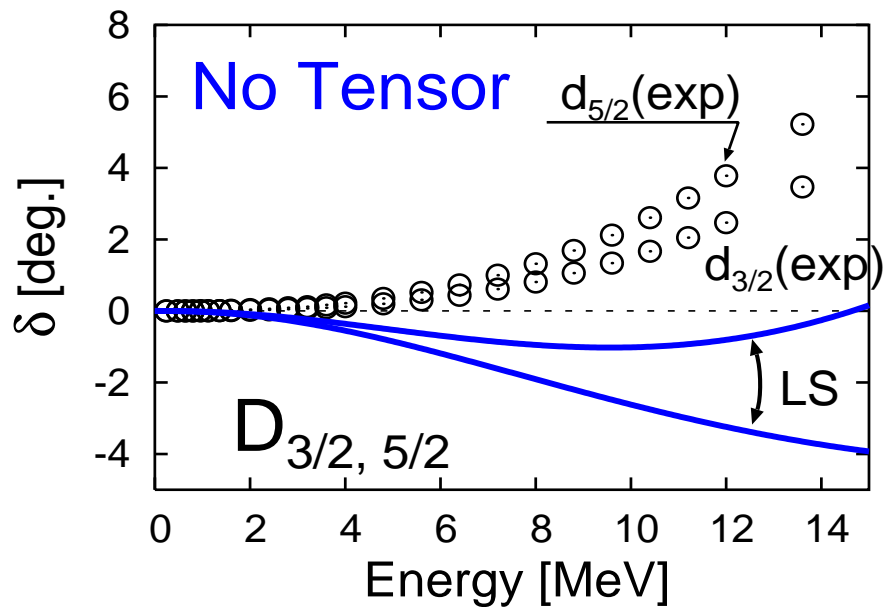
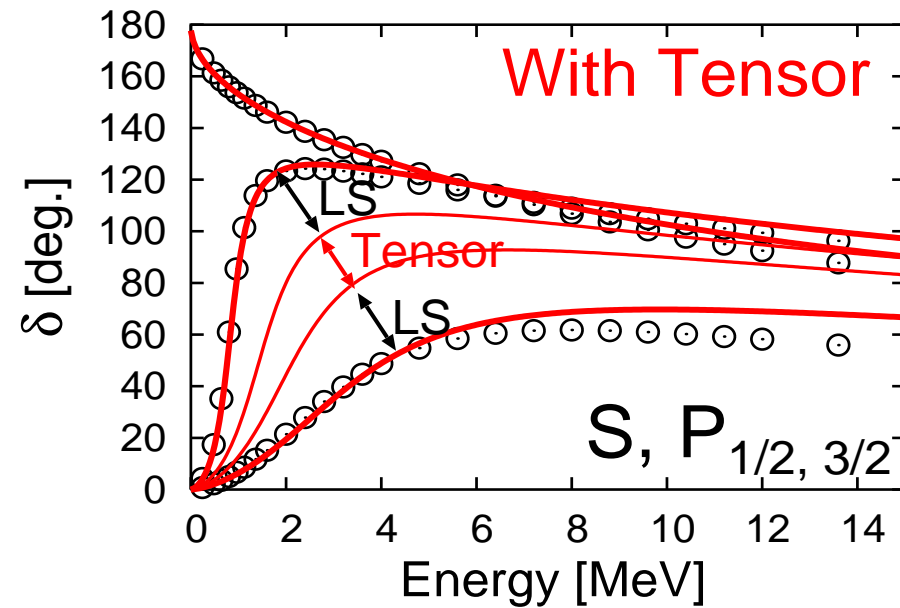
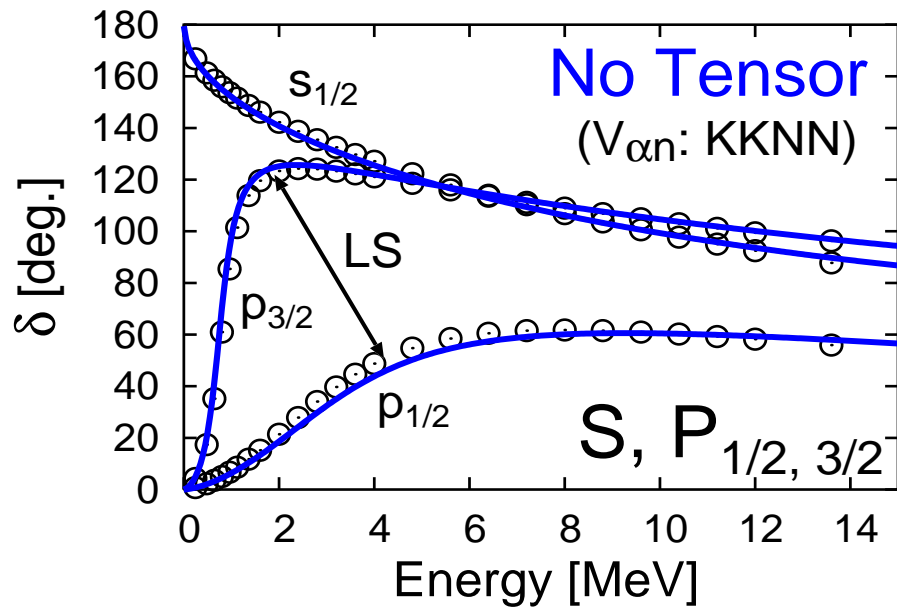
- The system is solved by OCM equation.

$$\circ H({}^5\text{He}) = H({}^4\text{He}) + H_{\text{rel}}, \quad \Phi^J({}^5\text{He}) = \mathcal{A} \left\{ \sum_{i=1}^N \psi_i({}^4\text{He}) \cdot \chi_i^J(r) \right\}$$

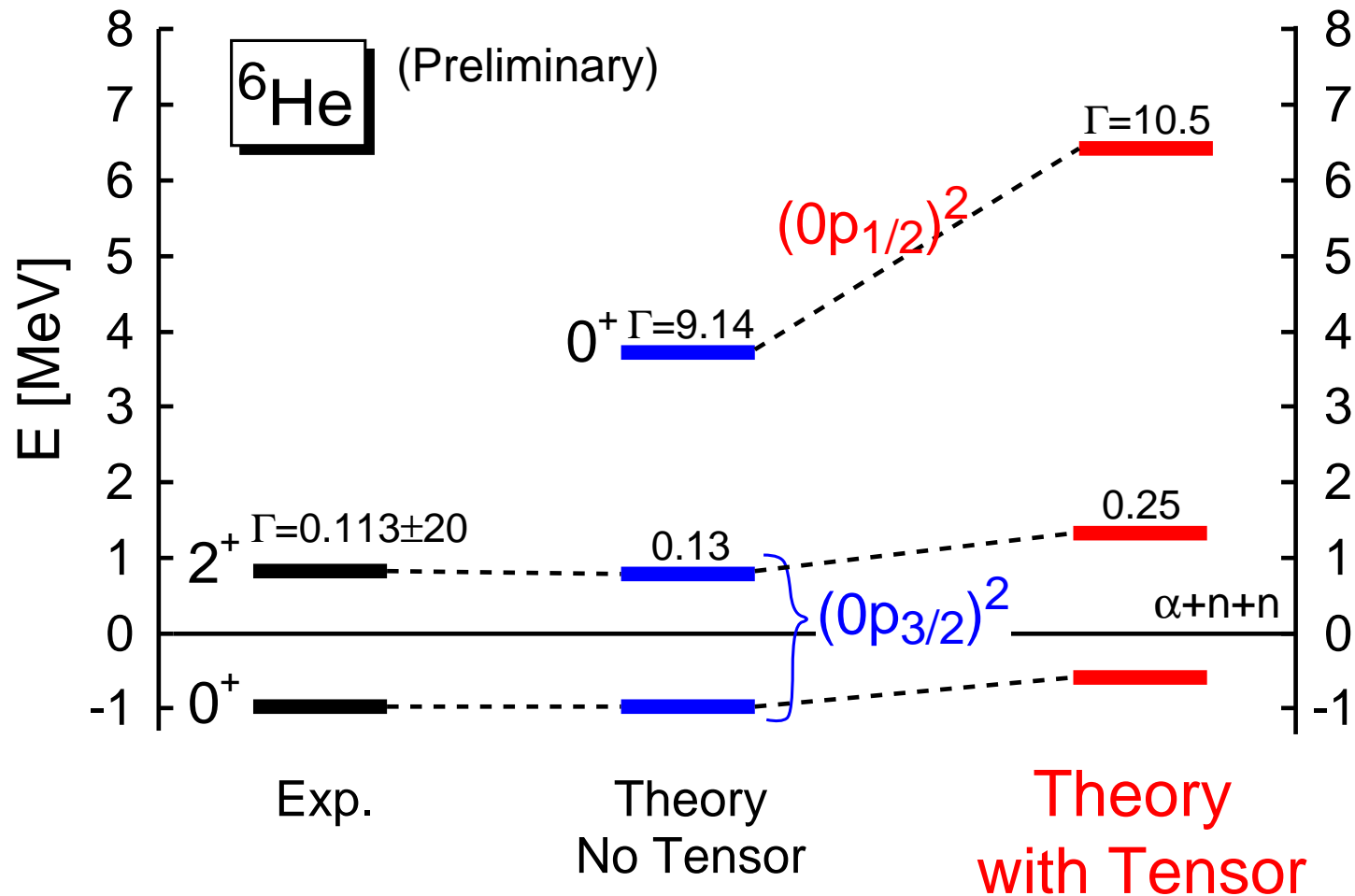
$$\circ \sum_{i=1}^N \left[h_{ij}({}^4\text{He}) + (T_r + V_{\alpha n} + \Lambda_i) \delta_{ij} \right] \chi_j^J(r) = E \chi_i^J(r)$$

$$\Lambda_i = \lambda \cdot \sum_{\alpha \in {}^4\text{He}} |\phi_\alpha\rangle \langle \phi_\alpha|, \quad \text{for } j=1, \dots, N.$$

Phase shifts of ${}^4\text{He-n}$ with TC



${}^6\text{He}$ in a coupled 3-body ${}^4\text{He}+n+n$ model



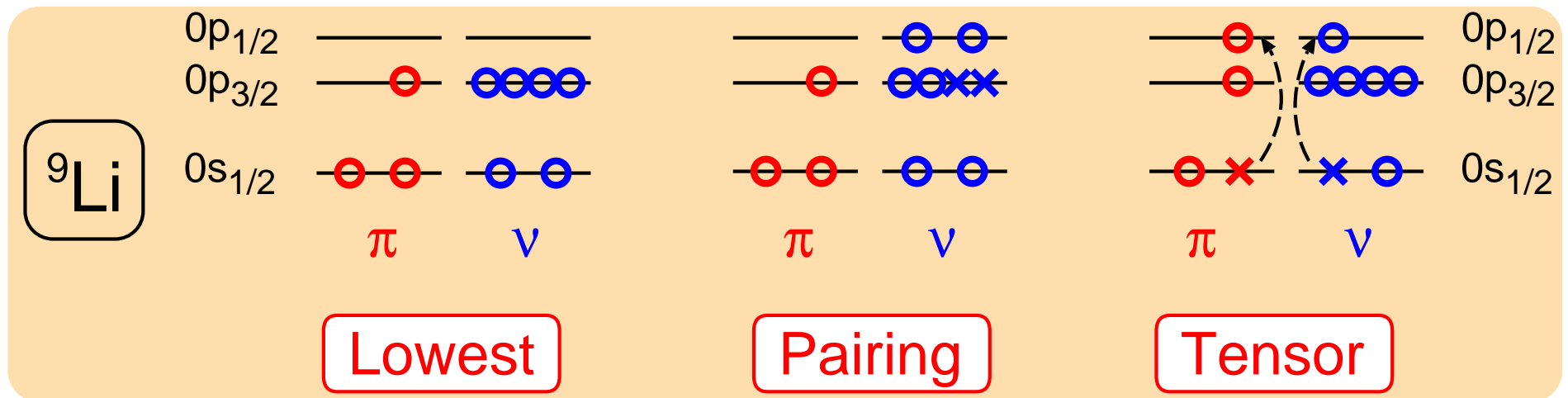
- $(0p_{3/2})^2$ does not change the energy \Rightarrow Naive ${}^4\text{He}+n+n$ model is applicable.
- $(0p_{1/2})^2$ loses the energy \Rightarrow Tensor suppression in 0_2^+
 \Rightarrow Contribute to Level inversion in $N=7,8$ isotones.

Structures of ^{11}Li

- Expt. : $S_{2n}=0.31$ MeV, $R_m=3.12\pm 0.16$ / 3.53 ± 0.06 fm (^9Li : 2.32 fm).
- Breaking of magic number $N=8$, halo structure in G.S.
 - Simon et al.(PRL83, expt) : $(1s)^2\sim 50\%$, Mechanism is unclear.
- Virtual s-state in ^{10}Li :
 - Expt. : Invariant mass spectrum of $^9\text{Li}-n$.
 - Many theories assume a deep s-wave $^9\text{Li}-n$ potential.
 - Thompson-Zhukov(PRC49) / Garrido-Fedorov-Jensen(NPA700).
 - Halo structure is reproduced for G.S., and soft dipole resonances appear.
- Our group performs the extended three-body model analysis with configuration mixing for ^9Li (PTP101, PTP108, PLB576, pairing).

Explicit tensor and pairing correlations in ${}^9\text{Li}$ for analysis of ${}^{11}\text{Li}$

- Configuration mixing with H.O. basis function (TM, K.Katō, K. Ikeda, PTP113)
(S. Sugimoto et al for ${}^{12}\text{C}$, ${}^{16}\text{O}$)
- $0s + \overline{0p} + \overline{1s0d}$ within **2p2h excitations**.
- Length parameters $\{b_\alpha\}$ are determined **independently and variationally**.
This is useful to represent the **high momentum component**.



Superposition of the tensor and pairing correlations in ${}^9\text{Li}$

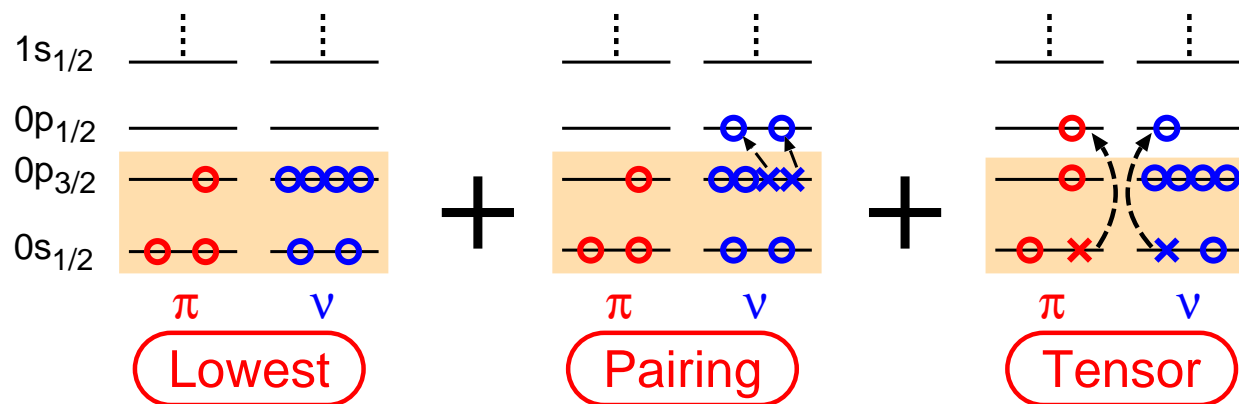
Energy [MeV]	-44.3
$\langle V_T \rangle$ [MeV]	-31.9
R_m [fm]	2.30
0p0h	78.5 %
$(0p_{3/2})_{01}^{-2}(\overline{0p}_{1/2})_{01}^2$	8.8 %
$(0s_{1/2})_{JT}^{-2}(\overline{0p}_{1/2})_{JT}^2$ (JT)=(10)	6.8 %
(JT)=(01)	0.2 %
$(0s_{1/2})_{10}^{-2}[(\overline{1s}_{1/2})(\overline{0d}_{3/2})]_{10}$	1.9 %
$(0s_{1/2})_{10}^{-2}(\overline{0d}_{3/2})_{10}^2$	1.2 %

- Tensor correlation:

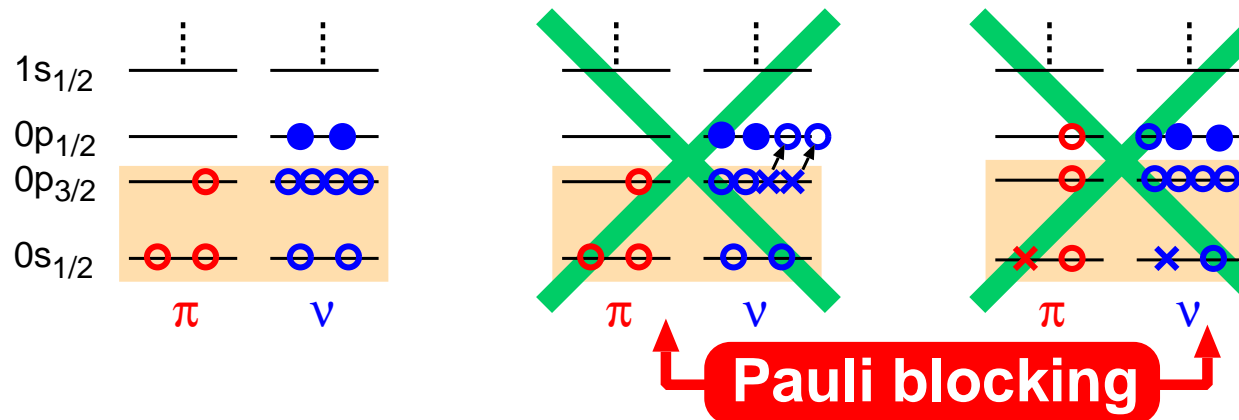
- 0^- coupling of $0s_{1/2} - 0p_{1/2}$
 \Rightarrow pion nature of V_T
- (J,T)=(1,0)
 \Rightarrow deuteron correlation

Effect of pairing and tensor correlation in ^{11}Li

^9Li
(GS)

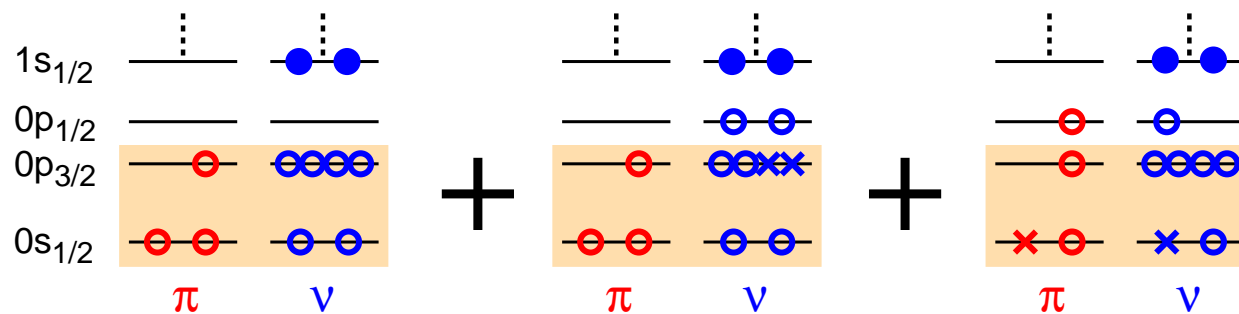


^{11}Li
(p^2)



energy loss

^{11}Li
(s^2)



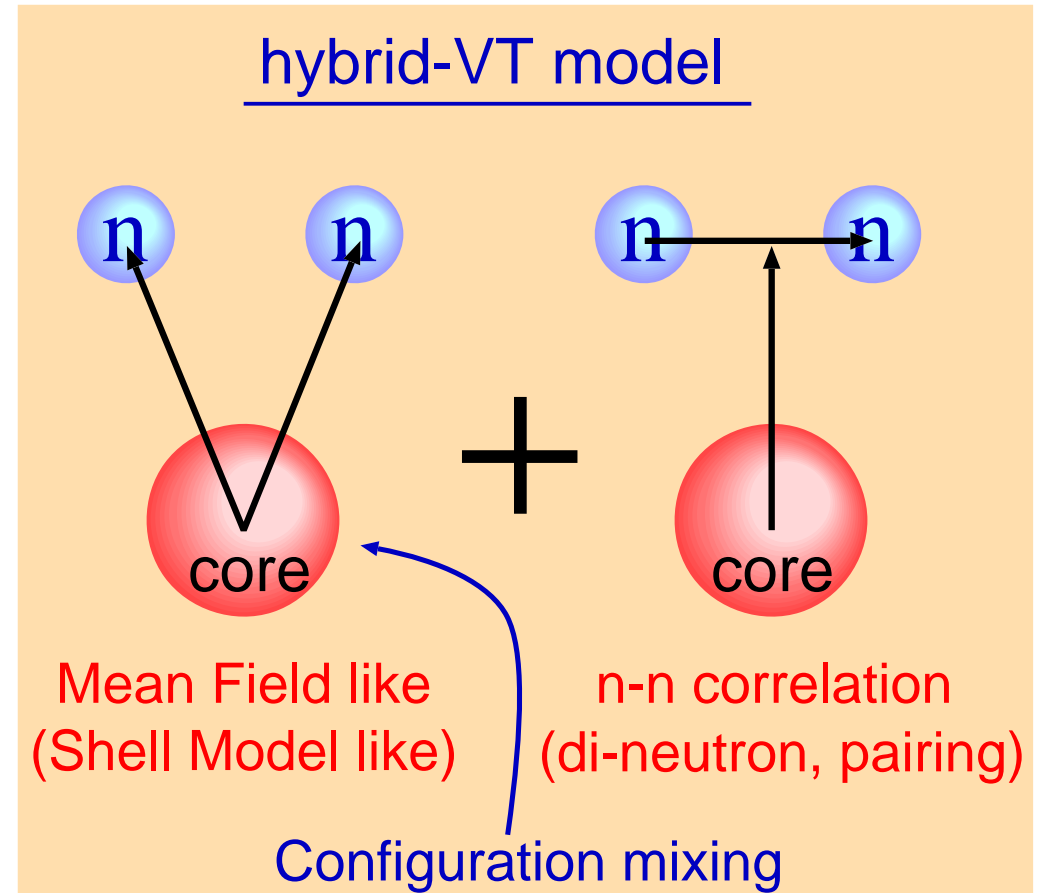
energy gain

increase
 s^2 mixing

Pairing-blocking : K.Katō, T.Yamada, K.Ikeda, PTP101('99)119, H.Masui, S.Aoyama, TM, K.Katō, K.Ikeda, NPA673('00)207.
TM, S.Aoyama, K.Katō, K.Ikeda, PTP108('02)133, H.Sagawa, B.A.Brown, H.Esbensen, PLB309('93)1

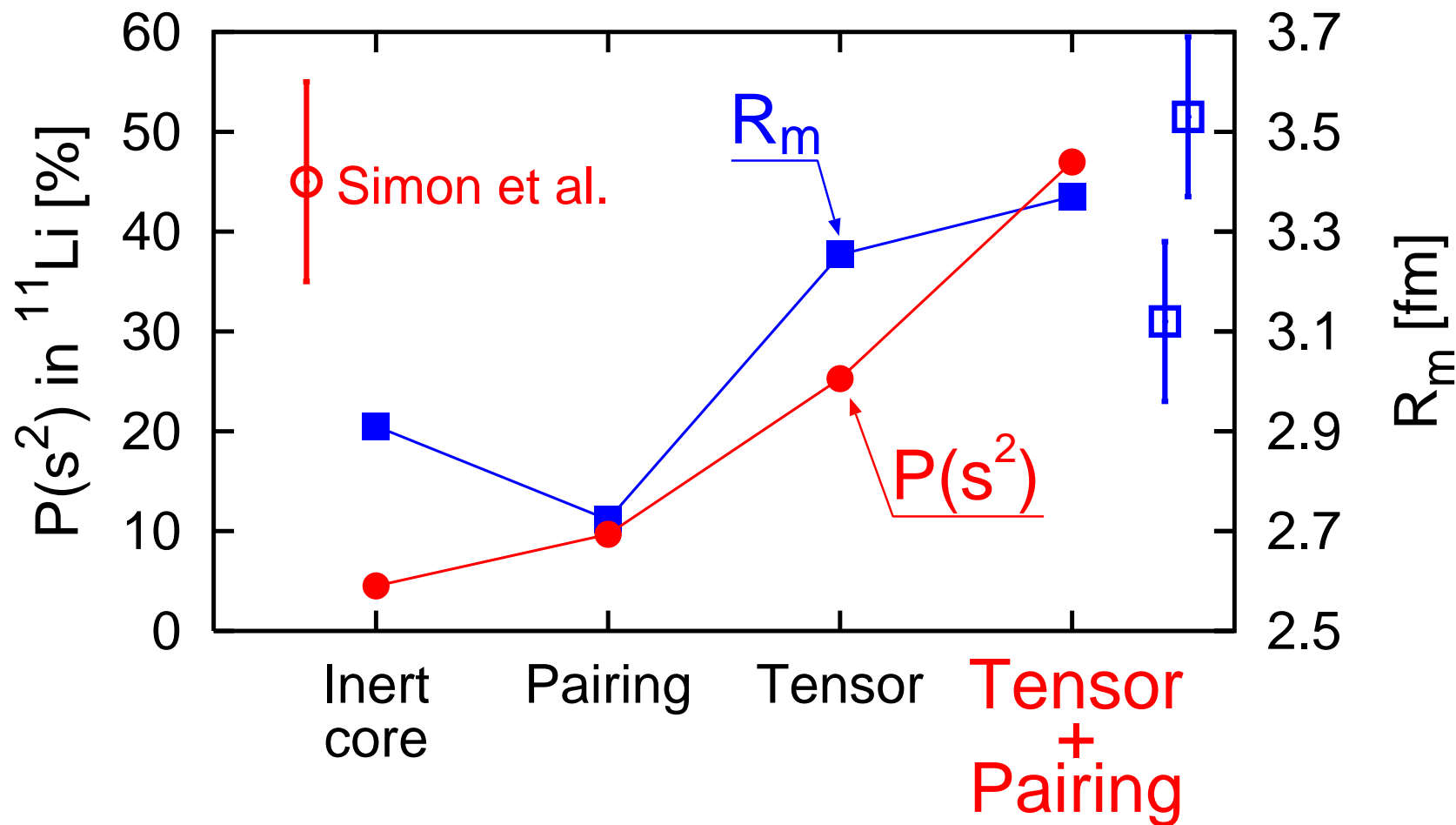
Hamiltonian for ^{11}Li in the orthogonarity condition model

- Folding potential with MHN(G-matrix)
 - Same strength for s- and p-waves.
 - Adjust to reproduce $S_{2n}=0.31$ MeV.
- Argonne potential (AV8') for last 2n.



[Ref] TM, S. Aoyama, K.Kato, K.Ikeda, PTP108(2002)

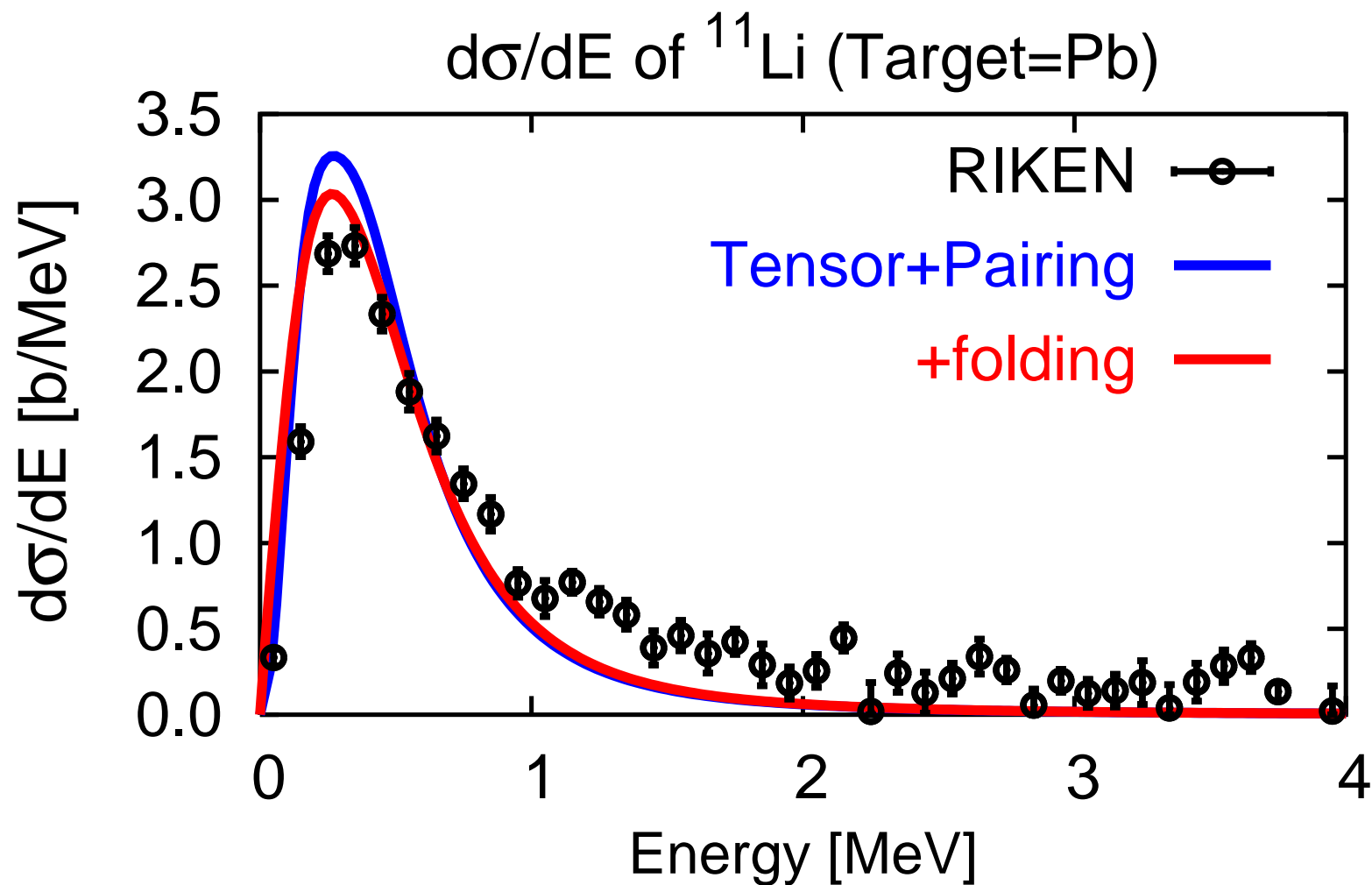
^{11}Li G.S. properties with tensor and pairing correlations



$E(s^2) - E(p^2)$	2.1	1.4	0.5	-0.1	[MeV]
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pairing correlation couples $(0p)^2$ with $(1s)^2$ for last 2n

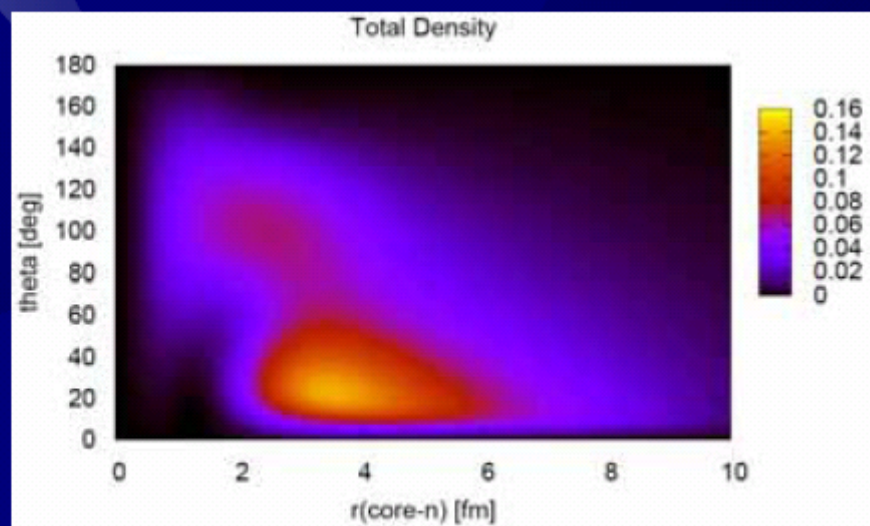
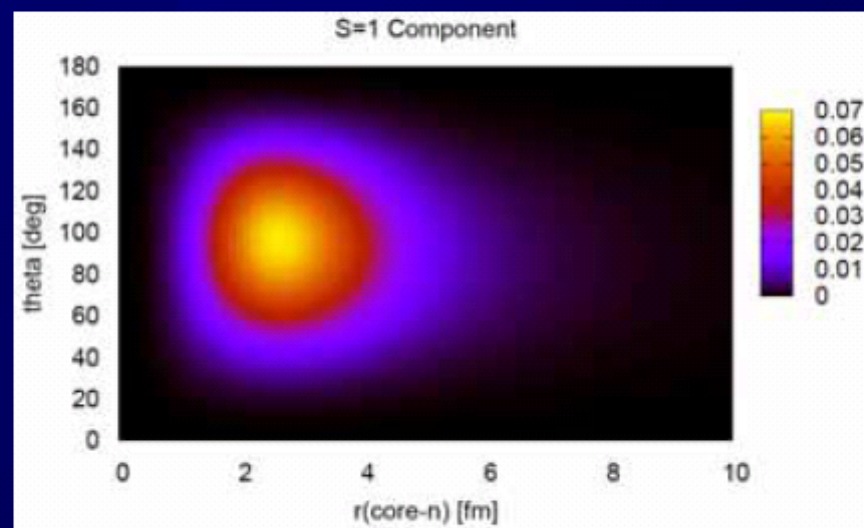
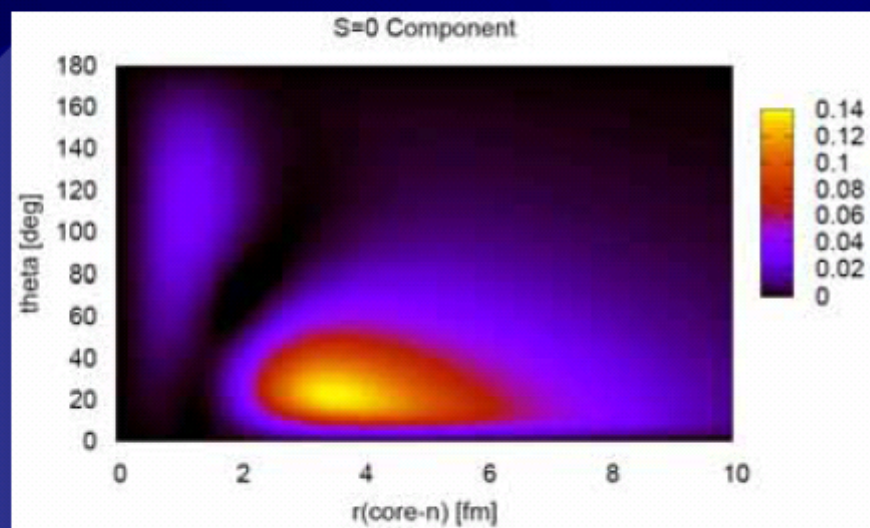
Coulomb breakup strength into ${}^9\text{Li}+n+n$ system



- No three-body resonances.
- Exp. : Nakamura et al., RIKEN Accel.Prog.Rep.38.

コア励起を含む場合の確率密度分布

4. Core(3 config.) & Tensor+Pairing $\Rightarrow s^2=46.9\%$



Summary

1. **Explicit effects of the tensor correlation (TC)** are investigated in light nuclei.
2. ${}^4\text{He}$:
 - i) **Three specific 2p2h configurations** are favored to describe TC.
 - ii) **spatial shrinkage of the particle states** is important to describe TC.
3. ${}^5\text{He}$: Contribute to LS splitting / ${}^6\text{He}$: Tensor suppression in 0_2^+ .
 - $0p_{1/2}$ -orbit neutrons cause the strong **Pauli-blocking**.
4. ${}^{11}\text{Li}$: Tensor suppression leads to the large admixture of $(1s)^2$ in G.S.
 - **“Tensor+Pairing”** naturally explains **50 % of $(1s)^2$** .
5. Future problem : How to treat the short-range correlation.