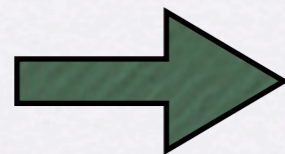


DM and new physics

What is the connection to DM and LHC

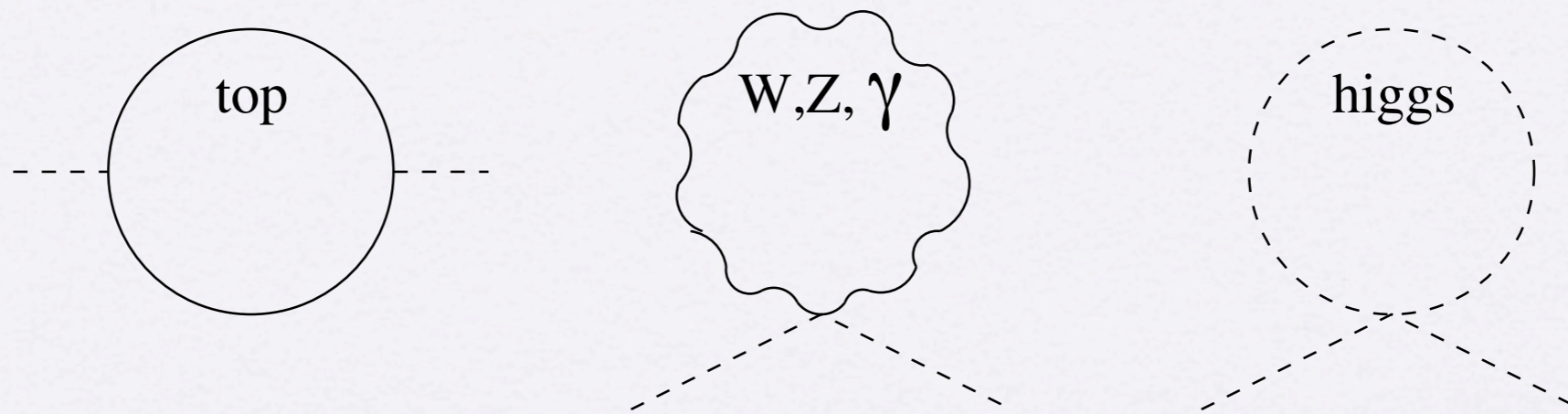
- In the big-bang scenario DM must be produced at high energy collisions. LHC, pp collider at 14 TeV make a collision of particle at 1 TeV
- Standard model has hierarchy problem. To solve this the modification of gauge and top sector required.



Observation of DM production at LHC

New Physics, Clue

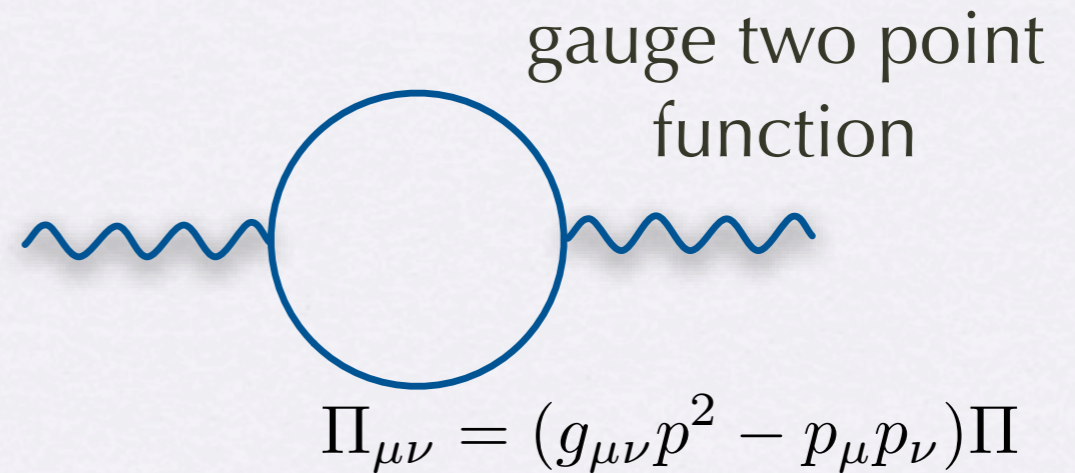
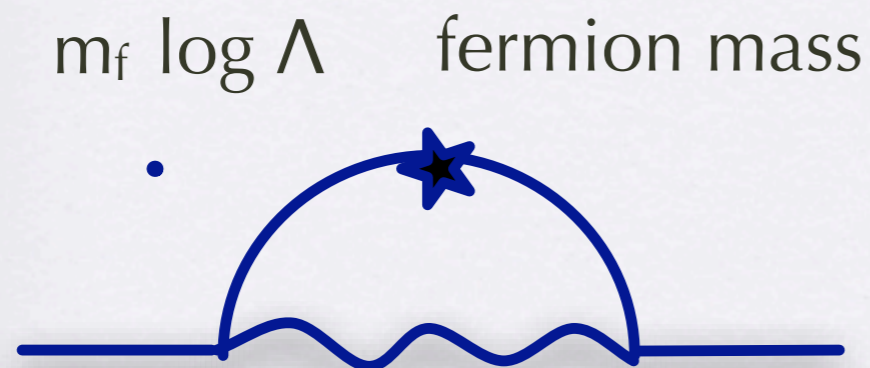
Fine tuning in the Higgs sector



top loop	$-\frac{3}{8\pi^2} \lambda_t^2 \Lambda^2$	$\sim -(2 \text{ TeV})^2$
$SU(2)$ gauge boson loops	$\frac{9}{64\pi^2} g^2 \Lambda^2$	$\sim (700 \text{ GeV})^2$
Higgs loop	$\frac{1}{16\pi^2} \lambda^2 \Lambda^2$	$\sim (500 \text{ GeV})^2$.

Others are reasonable

Why Higgs vev is $O(200) \text{ GeV}??$



New Symmetry \rightarrow New Particle

- Need control on the radiative correction to the Higgs sector

- ideas

- chiral symmetry (extended to boson sector)
- global symmetry (little Higgs model)
- gauge symmetry (gauge higgs unification)

- Or planck scale is low (Extra dimension model)

- On the other hand \ll we see no effect of BSM in radiative correction

$$\delta L = \frac{(h^\dagger D_\mu h)^2}{\Lambda^2} \quad \Lambda > 5\text{TeV}$$



Classic Solution: Supersymmetry

- exchange boson and fermion. $\phi \leftrightarrow \psi$
 - sfermions(0), gaugino(1/2), higgsinos(1/2)
- SUSY change “dimension” (1 for boson 3/2 for fermion), relate mass and couplings
$$\Phi = \frac{1}{g^2} + M\theta^2 \quad \Phi_{WW} = \frac{1}{g^2} F_{\mu\nu} F^{\mu\nu} + M\tilde{g}\tilde{g}$$
- chiral symmetry is extended to boson sector. No new dimensionless coupling and no quadratic divergence
$$\lambda\psi_L\psi_R H \rightarrow \lambda\phi_L\psi_R\tilde{H} + \lambda\psi_L\phi_R\tilde{H}$$
- R parity conservation. New stable particle \rightarrow DM candidate.

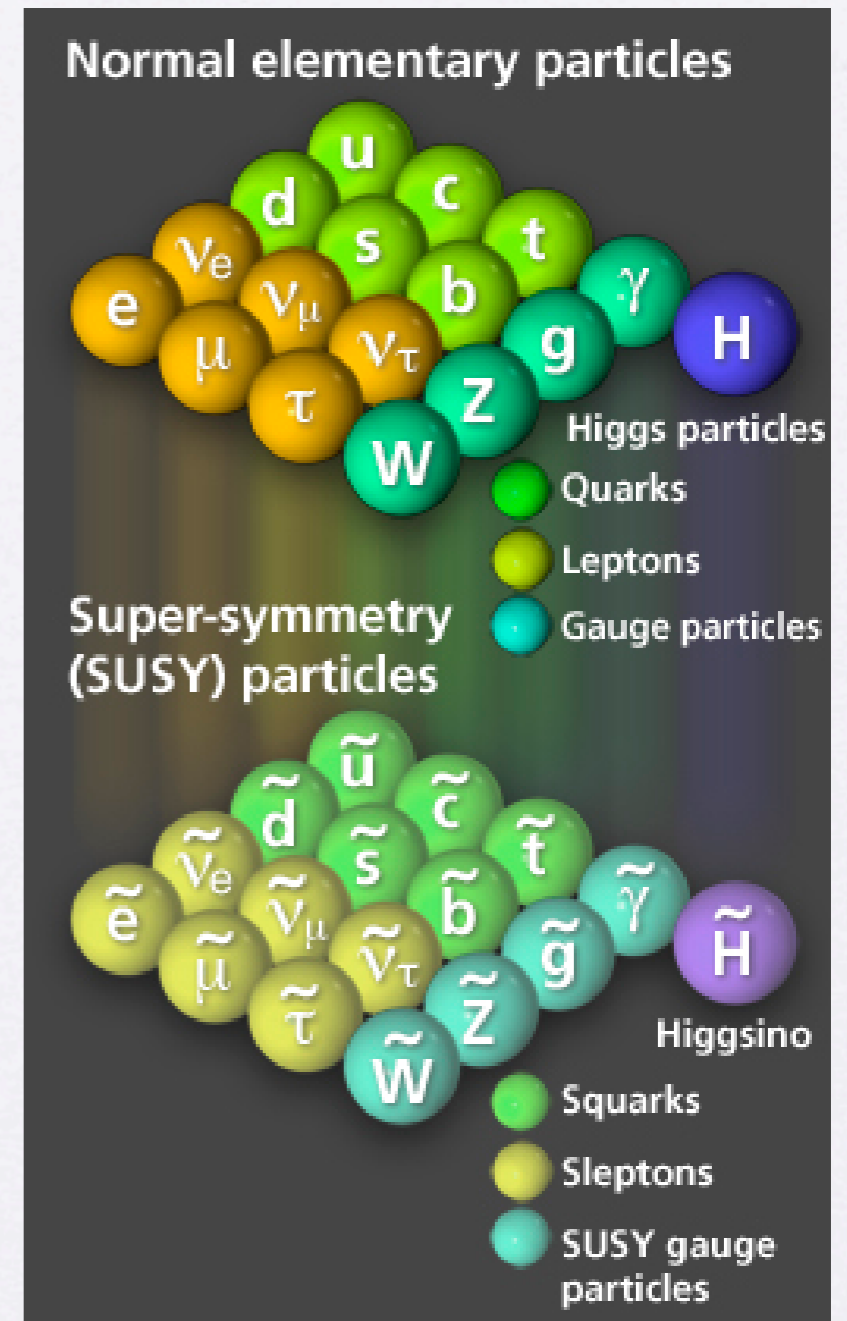
matter content of MSSM

	quark	squarks
$(3, 2)_{1/6}$	$q_L = (u, d)_L$	$\tilde{q}_L = (\tilde{u}_L, \tilde{d}_L)$
$(3^*, 1)_{-2/3, 1/3}$	u^c, d^c	\tilde{u}^c, \tilde{d}^c
	lepton	slepton
$(1, 2)_{1/2}$	$l_L = (\nu, e)_L$	$\tilde{l}_L = (\tilde{\nu}_L, \tilde{e}_L)$
$(1, 1)_1$	$l_R = e_R^c$	\tilde{e}_R^c
	higgsino	higgs
$(1, 2)_{-1/2}$	$(\tilde{H}_1^0, \tilde{H}_1^-)$	$H_1 = (H_1^0, H_1^-)$
$(1, 2)_{1/2}$	$(\tilde{H}_2^+, \tilde{H}_2^0)$	$H_2 = (H_2^+, H_2^0)$

vector multilets

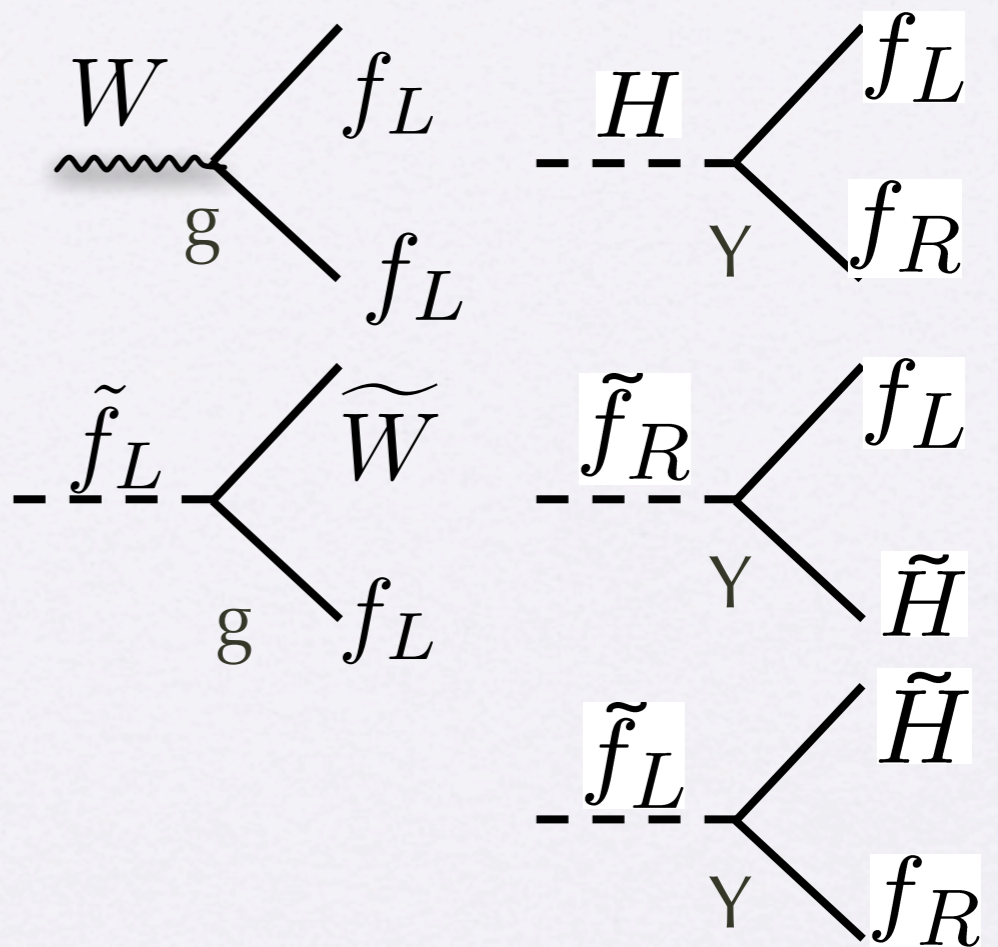
$G_\mu, W_\mu, B_\mu \leftrightarrow \tilde{G}, \tilde{W}, \tilde{B}$ (gluino, wino, bino)

gravity $\tilde{\psi}^\mu$ (gravitino), $g_{\mu\nu}$ (graviton)



R parity, SUSY relation

- R parity conservation
 - SUSY particles will be pair produced.
 - SUSY particles decay into SUSY particles
- There are no new dimensionless couplings
 - gaugino interaction is gauge coupling
 - Higgsino matter interaction is yukawa coupling



Higgsino-gaugino mix
 sfermion left right also
 mix due to scalar 4 point
 couplings

Should we only consider SUSY Dynamical symmetry breaking ?

- Technicolor \rightarrow Little Higgs model
 - Higgs boson is goldstone boson of a large symmetry. $SU(5) \rightarrow SO(5)$
 - Gauge symmetry: $SU(2)_1 \times SU(2)_2 \times U(1)_1 \times U(1)_2$
 (g_1, g_2, g'_1, g'_2)
 - quadratic correction to Higgs sector starts from 2 loop
 - top sector must be extended (extra top quark). after all **top-higgs coupling is the source of fine tuning.**
$$\chi = (b_3, t_3, \tilde{t}) \quad \tilde{t}, \tilde{t}'$$
- However it is rather difficult to make simple Little Higgs model and LEP data consistent .

LEP Anchor

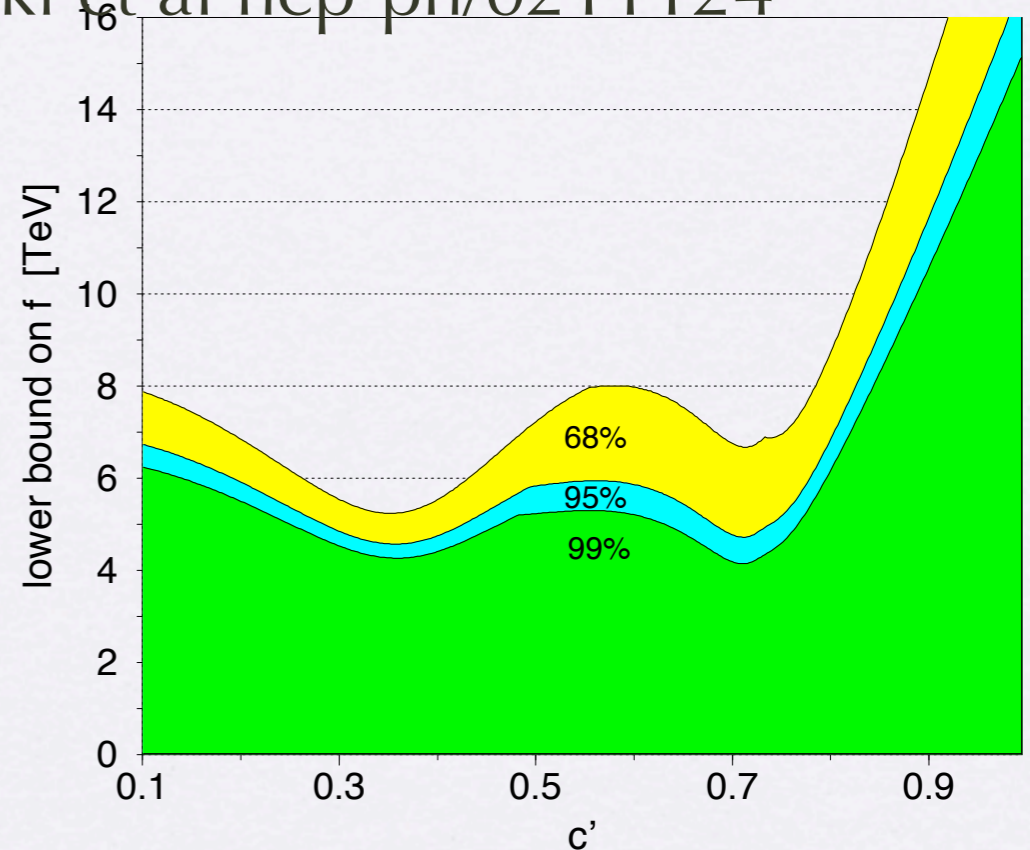
difficulty comes from tree level Heavy-Light mixing

$$L_2 \rightarrow -\frac{g_1 g_2 (g_1^2 - g_2^2)}{4(g_1^2 + g_2^2)} W_{\mu L}^a W_H^{a\mu} h^2 \rightarrow -\frac{g^2 (s^2 - c^2)^2}{8f^2} W_L W_L h^4$$

$$(W_L^a = sW_1^a + cW_2^a)$$

- Various v^2/f^2 corrections. proportional to the coupling difference, $\Delta g = g_1 - g_2$
- $M^2(W_H) = (g_1^2 + g_2^2) f^2/4 \sim (gf/2)^2 > 2.7 \text{ TeV}$
- $f > 4 \text{ TeV}$ $m(t') > 7 \text{ TeV}$, (Hewett et al JHEP, 2003) Fine tuning is reintroduced

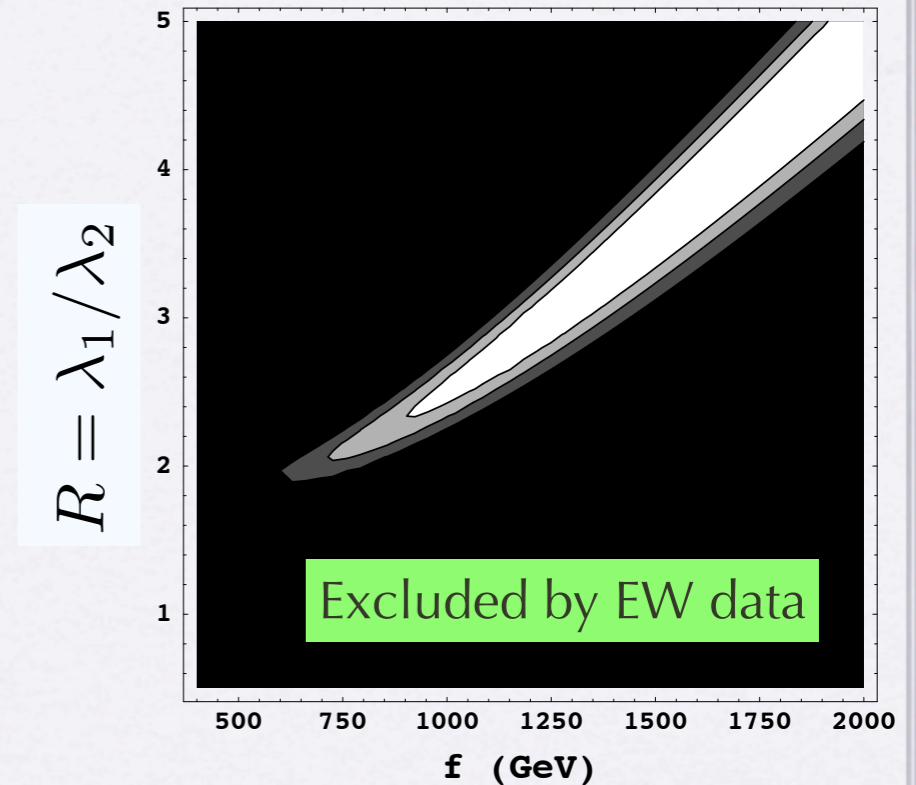
Csaki et al hep-ph/0211124



Little Higgs with T-parity

- gauge groups and matter contents respect T parity. $SU(2)_1 \rightleftharpoons SU(2)_2$ $U(1)_1 \rightleftharpoons U(1)_2$
- T-odd matters are introduced. Looks like SUSY without gluino
- LEP constraint is weaker.
 - Heavy gauge bosons and triplet higgs boson live in T-odd sector. No tree level mixing
- **Need more attempts to construct a model including symmetry breaking sector. (cf. the study of SUSY breaking sector.)**
- UED has similar nature.

Hubisz et al 0411264



$$m_{T_-} = \lambda_2 f$$

The Lesson is

- LEP constraint (small radiative correction)
 - New Physics scale Λ is high, suggesting fine tuning.
 - Need symmetry to cancel divergence
 - top partner \rightarrow top must be involved in the symmetry.
- “DM” and “radiative correction” \rightarrow parity structure

LHC signature: strongly interacting particle decay into DM
(and flavor sector involving b quark.....)

Supersymmetry and DM

SUSY mass spectrum

Radiative correction due to the gauge interactions. Strongly interacting particles are heavy weakly interaction particles are light.

$$\begin{aligned}m_{\tilde{q}_L}^2 &= m_0^2 + 4.5M_{1/2}^2 \\m_{\tilde{l}_L}^2 &= m_0^2 + 0.5M_{1/2}^2 \\m_{\tilde{l}_R}^2 &= m_0^2 + 0.16M_{1/2}^2\end{aligned}$$

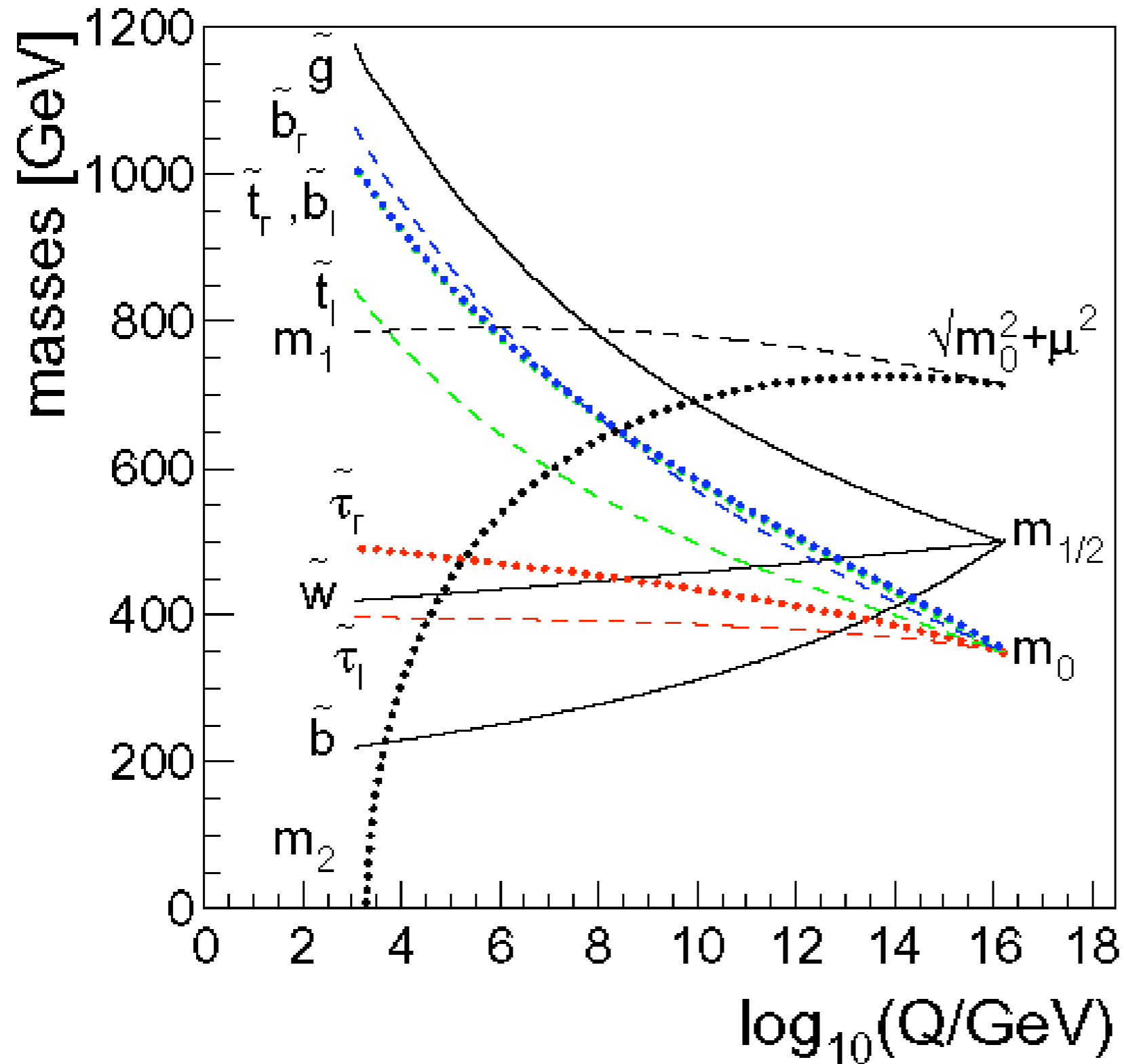
Common scalar mass m_0
common gaugino mass $M_{1/2}$
at Unification scale.

working assumption

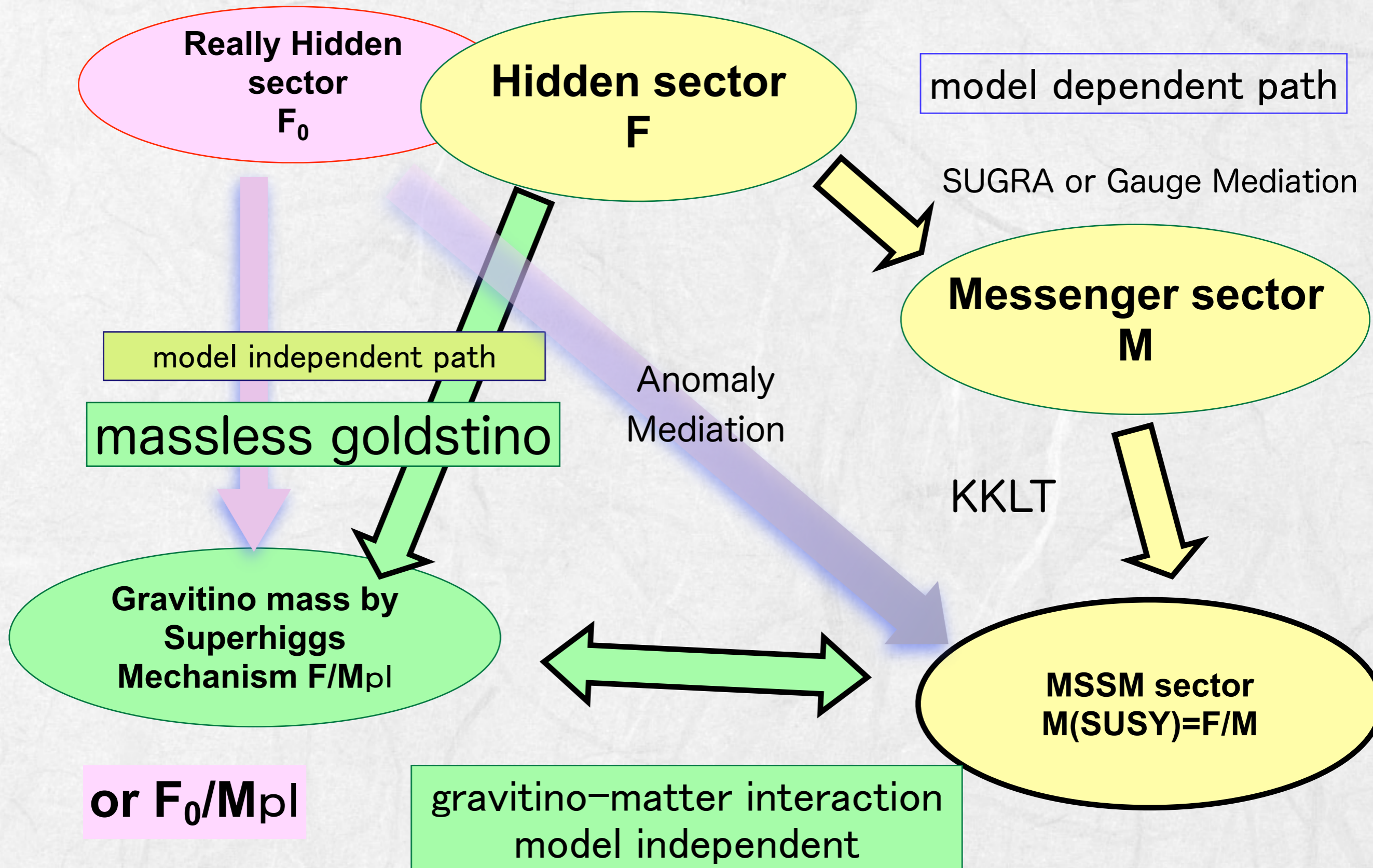
Yukawa interaction
mass difference among
generations , Effects on
generation mixing of SUSY
particles

Low energy mass
spectrum of SUSY
particles

Radiative Symmetry Breaking



Supersymmetry-a picture



SUSY breaking scenarios and mass spectrum

- Low energy phenomenology is not the end of the story .
- Hidden sector break supersymmetry. “flavor and CP” problem
 - gravity mediation, gauge mediation, anomaly mediation(string inspired mixed cases) , “geometric separation”
- Problems (why alternatives are searched for)
 - Light higgs boson (hope and/or worry) little hierarchy
 - DM constraints
 - gravitino, string moduli.....



Rich Field!

DM candidate in SUSY

- ✿ neutralino LSP

- ✿ a neutralino is a mixture of gauginos and Higgsinos
- ✿ $\Omega(\text{th})h^2 \sim 0.1 \Leftrightarrow$ light slepton, Higgs exchange, or gaugino-higgsino mixing, light annihilation.

- ✿ gravitino LSP

- ✿ no prediction on the density.
- ✿ direct detection is not possible
- ✿ need additional trick to explain cosmic ray anomaly. (for example R parity violation --later)

- ✿ sneutrino essentially excluded

But in general, it is good to have
a DM candidate in the model

The nature of the Lightest Neutralino

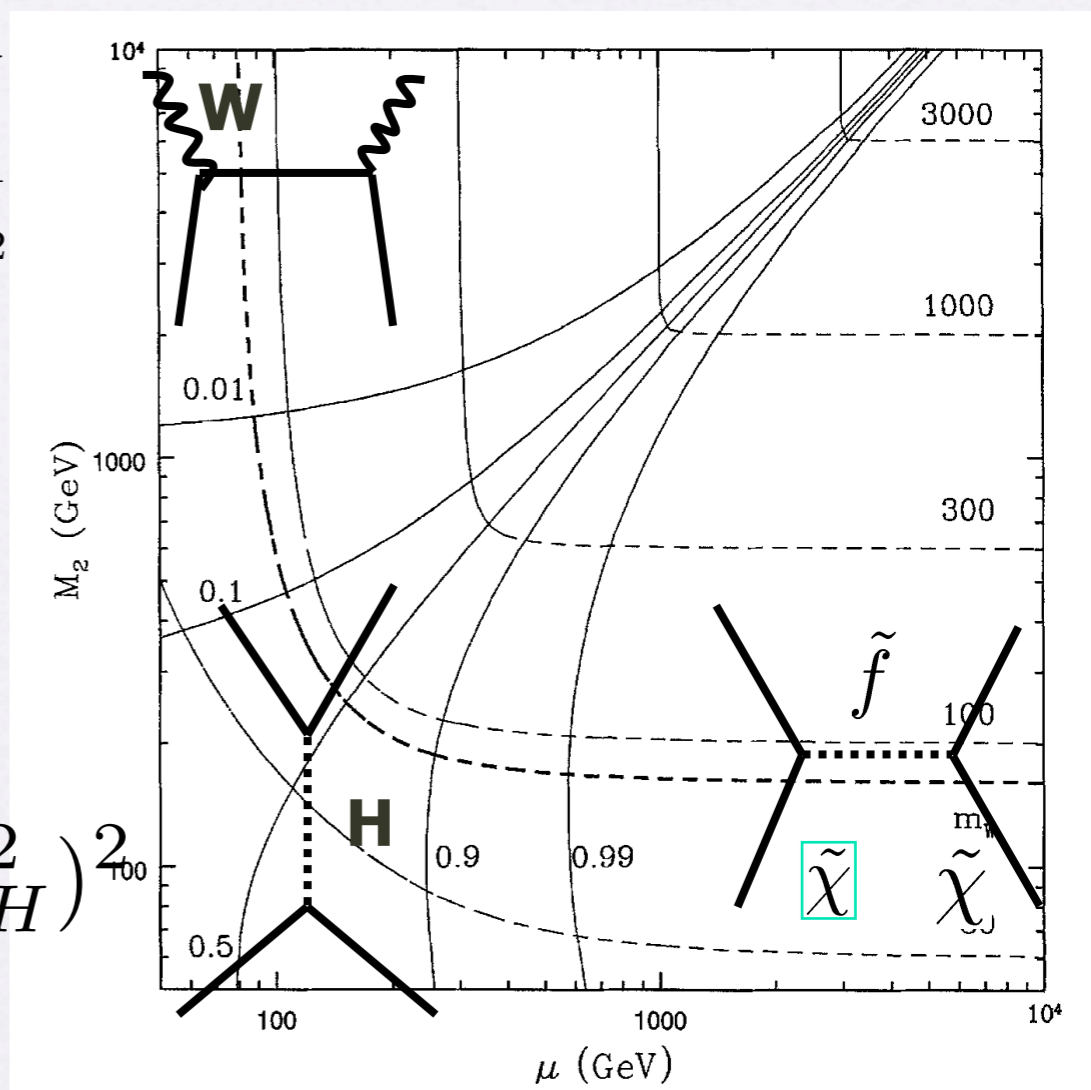
Neutralino mass matrix

$$M = \begin{pmatrix} M_1 & 0 & -m_Z s_W c_\beta & m_Z s_W s_\beta \\ 0 & M_2 & m_Z c_W c_\beta & -m_Z c_W s_\beta \\ & & 0 & -\mu \\ & & -\mu & 0 \end{pmatrix} \begin{pmatrix} \tilde{B} \\ \tilde{W} \\ \tilde{H}_1 \\ \tilde{H}_2 \end{pmatrix}$$

$$M_1 \ll \mu \quad \sigma v \propto m_{\tilde{\chi}}^2 / m_{\tilde{l}}^4$$

$$M_1 \gg \mu \quad \sigma v \propto 1 / m_{\tilde{\chi}}^2$$

$$M_1 \sim \mu \quad \sigma v \propto m_{\tilde{\chi}}^2 / (4m_{\tilde{\chi}}^2 - m_H^2)^2$$



DM density constraint is important in "MSUGRA"

1) bulk: LSP is Bino like.
 Slepton exchanges fix the density ◦

$$\Omega h^2 \propto m_{\tilde{l}}^4 / m_{\tilde{\chi}}^2$$

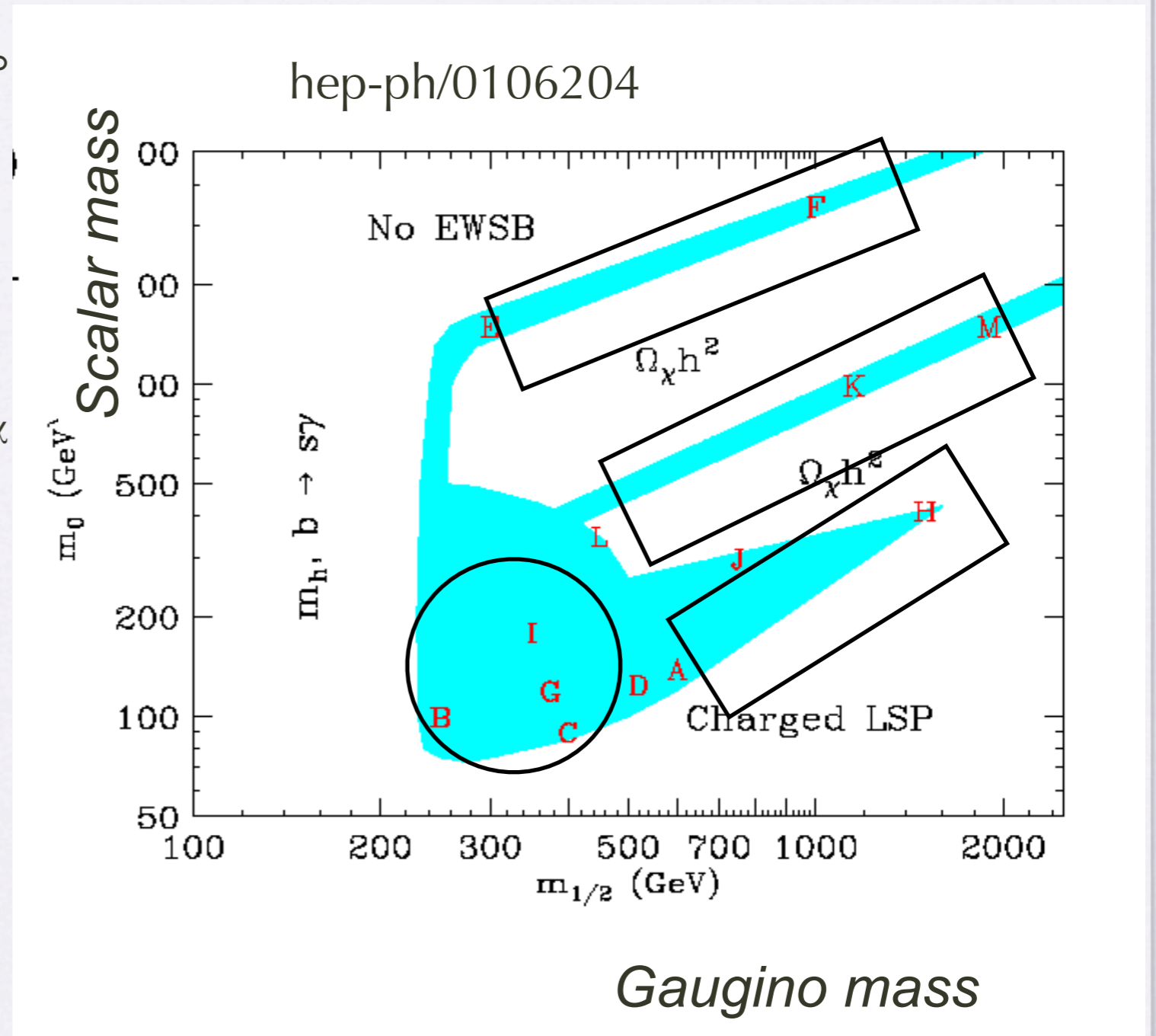
too large mass density

2) Higgs pole effects $m_H = 2m_\chi$

3) coannihilation region

$$\tilde{\tau} \tilde{\chi}$$

4) focus point region:
 higgsino-gaugino mixing



Target at LHC

- finding LSP
- measure gaugino and chargino masses to fix $M_1, M_2, \mu, \tan\beta$
- stau masses.
- measure Higgs mass (especially heavy ones)
- gravitino? (if I have a time)