

Lattice QCD simulation with exact chiral symmetry

--- Not yet at finite temperature ---

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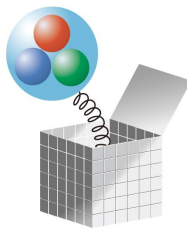
High Energy Accelerator Research Organization (KEK)

2008年9月3-5日 基研研究会「熱場の量子論とその応用」



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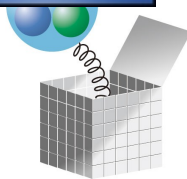
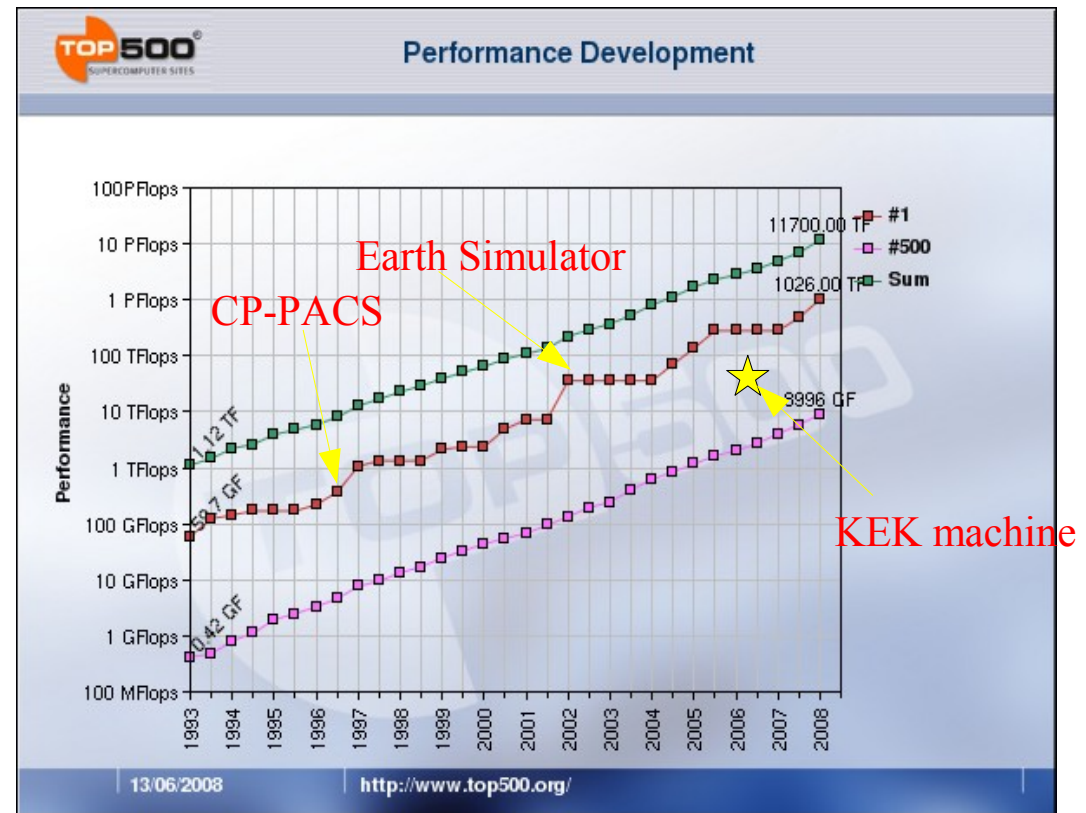
- Introduction
- Lattice fermion actions
- Overlap fermion
- JLQCD's dynamical overlap project
- Toward simulations at finite temperature
- Summary/outlook





Introduction

- Lattice QCD has been an application which requires most large computational resources
- Growth of computational power → new era of lattice simulations
- Around 2010, ~10PFlops expected
 - Japan: Next generation supercomputer project (10TFlops in 2011)
 - Two projects in USA

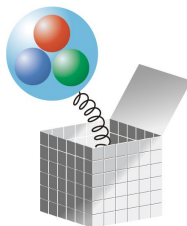




Introduction

- 1980's
 - Dawn of lattice simulations
- 1990's
 - Large quenched simulations (w/o dynamical quark effect)
 - Dynamical simulations (exploratory)
 - Matrix elements with 10% accuracy: systematic errors from continuum/chiral limit, quenching
- 2000's
 - Large dynamical simulations
 - Chiral symmetry (domain-wall, overlap,...)
 - Matrix elements with a few % accuracy
- 2010's ?
 - Nuclear physics
 - Thermodynamics

With accuracy!





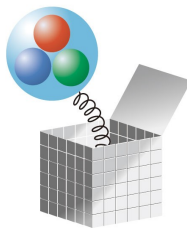
Introduction

Our choice (JLQCD Collaboration): overlap fermion

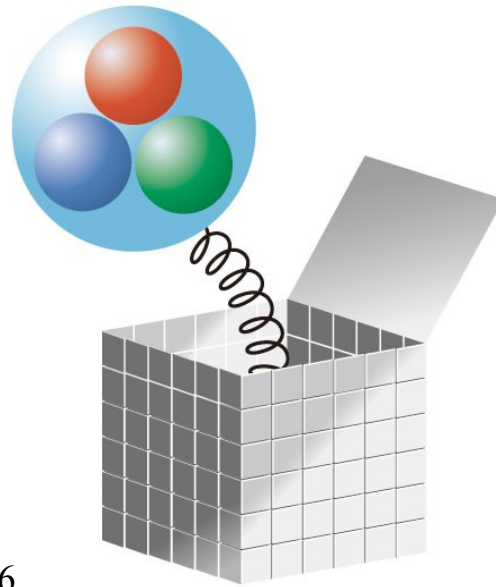
- Exact chiral symmetry (theoretically clean)
- Large simulation costs (numerically challenging)
- First large-scale dynamical simulation
- Now rich results are coming

Toward finite temperature (and density) simulation

- What kind of thermal properties to be explored?
- Cost/interest



Lattice fermion actions



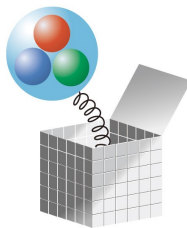
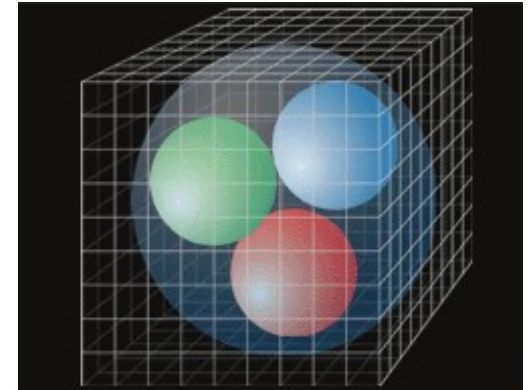
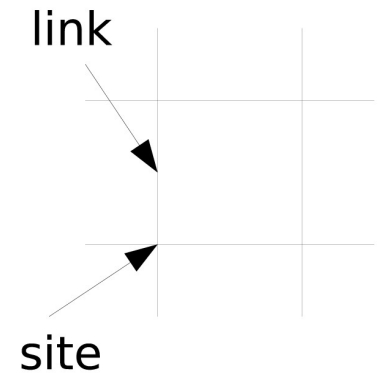
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Lattice QCD

Lattice QCD: gauge theory on 4D Euclidean lattice

- Regularized by lattice
 - gluon: SU(3) link variables
 - quark: Grassmann field on sites
- Continuum limit ($a \rightarrow 0$) \rightarrow QCD
- Numerical simulation by Monte Carlo
- Nonperturbative calculation
- Quantitative calculation has become possible
 - Matrix elements for flavor physics
 - Phase structure/plasma properties
- Chiral extrapolation is still a source of largest systematic errors!





Femion doubling/Wilson fermion

Fermion doubling

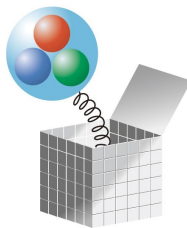
- naïve discretization causes 16-fold doubling
- Nielsen-Ninomiya's No-go theorem
 - Doublers appear unless chiral symmetry is broken

$$D\gamma_5 + \gamma_5 D = 0$$

Wilson fermion

- adds Wilson term to kill 15 doublers
- **breaks chiral symmetry explicitly** → additive mass renorm.
- Improved versions, twisted mass versions are widely used

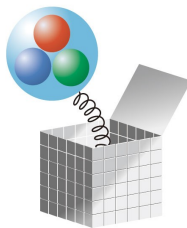
$$S_F = \frac{1}{2} \sum_{x,s} \bar{\psi}(x) \left[\gamma_\mu U_\mu(x) \psi(x + \hat{\mu}) - \gamma_\mu U_\mu(x - \hat{\mu}) \psi(x - \hat{\mu}) - \frac{r}{2} \Delta^{(2)} \psi(x) \right]$$





Staggered fermion

- Staggered fermion
 - $16=4$ spinors x 4 flavors ('tastes')
 - Remnant U(1) symmetry
 - Fourth root trick: still debated
 - Numerical cost is low: popular at finite temperature/density





Domain-wall fermion

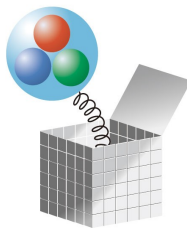
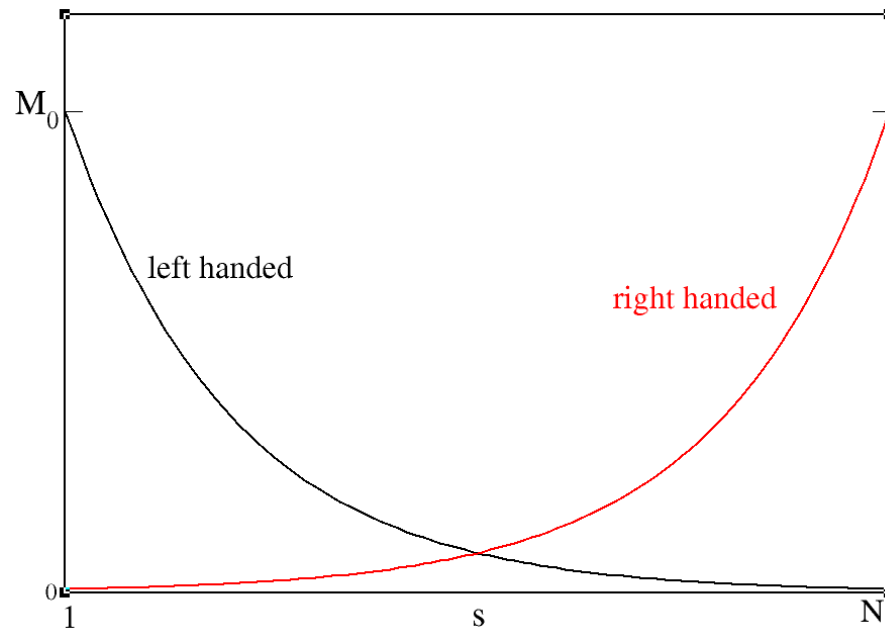
- **Domain-wall fermion** (Kaplan 1992, Shamir 1993)

- 5D formulation, light modes appear at the edges

- Symmetry breaking effect $m_{res} \rightarrow 0$ as $N_5 \rightarrow \infty$

$$S_F = \frac{1}{2} \sum_{x,s} \bar{\psi}(x,s) [D_W(x,s; -M_0)\psi(x,s) + (1 + \gamma_5)\psi(x,s+1) + (1 - \gamma_5)\psi(x,s-1) - 2\psi(x,s)]$$

- Costs O(20) times than Wilson fermions





Chiral symmetry on lattice

Ginsparg-Wilson relation (1982)

$$\gamma_5 D + D \gamma_5 = a R D \gamma_5 D$$

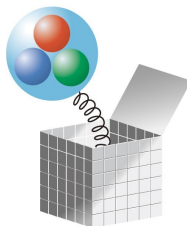
- Exact chiral symmetry on the lattice (Luscher 1998)

$$\delta\psi = \gamma_5 \left(1 - \frac{aR}{2} D \right) \psi \quad \delta\bar{\psi} = \bar{\psi} \left(1 - D \frac{aR}{2} \right) \gamma_5$$

- Satisfied by
 - Overlap fermion (Neuberger, 1998)
 - Fixed point action (Bietenholz and Wiese, 1996)

Caution:

Recently Mandula pointed out that Ginsparg-Wilson relation contains infinite number of lattice chiral transformation. What does this cause? [arXiv:0712.0651]



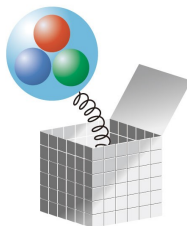


Overlap fermion

$$D = \frac{1}{Ra} \left[1 + \gamma_5 \text{sign}(H_W(-m_0)) \right]$$

H_W : hermitian Wilson-Dirac operator
(Neuberger, 1998)

- Theoretically elegant
 - Satisfies Ginsparg-Wilson relation
 - *Infinite* N_s limit of Domain-wall fermion (No m_{res})
- Numerical cost is high
 - Calculation of sign function
 - Discontinuity at zero eigenvalue of H_W
- Has become feasible with
 - Improvement of algorithms
 - Large computational resources

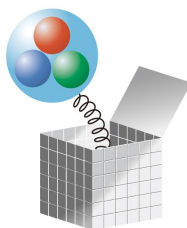




Costs

Large-scale dynamical simulation projects

Fermion	Chiral symmetry	flavor structure	cost	Collab.
Wilson-type	explicitly broken	simple	modest	PACS-CS, etc
Twisted mass	explicitly broken	simple	modest	ETM
Staggered	remnant U(1)	complex	low	MILC, etc
Domain-wall	good	simple	high	RBC, UKQCD
Overlap	best	simple	very high	JLQCD



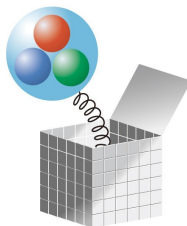


Why overlap fermion?

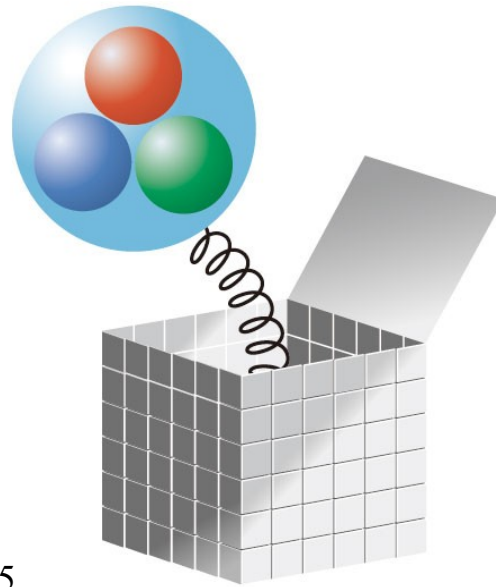
Exact chiral symmetry

- Exploring the chiral dynamics
 - Confirming the chiral symmetry breaking scenario
 - Epsilon regime
 - Phase transition at $T > 0$
- Matrix elements with controlled chiral extrapolation
 - no unwanted operator mixing
- Testing the effective chiral Lagrangian predictions
 - with exact chiral symmetry, continuum ChPT applies

Let us get rid of large cost by large computer power and **improved algorithms** !



Overlap fermion



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Overlap fermion

$$D(m) = \left(M_0 + \frac{m}{2}\right) + \left(M_0 - \frac{m}{2}\right) \gamma_5 \text{sign}(H_W)$$

H_W : hermitian Wilson-Dirac operator

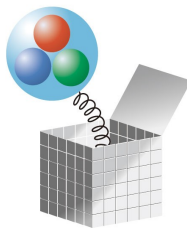
- Sign function means:

$$\text{sign}H_W \cdot v = \sum_{\lambda} \text{sign}(\lambda) (\psi_{\lambda}, v) \psi_{\lambda}$$

$(\lambda, \psi_{\lambda})$: eigenvalue/vector of H_W

In practice, all eigenmodes cannot be determined

- Reasonable solution:
 - Eigenmodes determined at low frequency part
 - Approximation formula for high mode part
 - ex. Chebychev polynomial, partially fractional, etc.



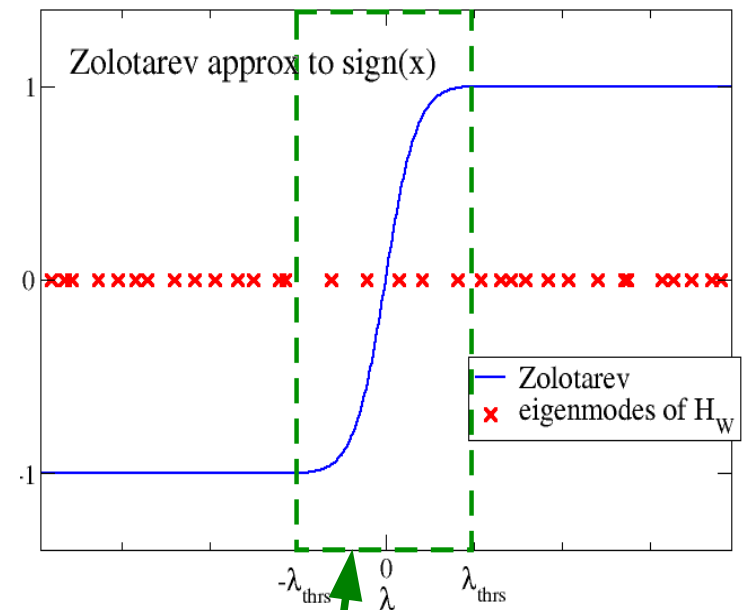


Sign function

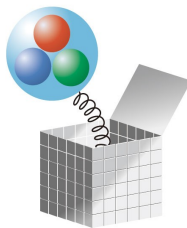
Zolotarev's Rational approximation

$$\text{sign}(H_W) = \frac{H_W}{\sqrt{H_W^2}} = H_W \left(p_0 + \sum_{l=1}^N \frac{p_l}{H_W^2 + q_l} \right)$$

- $(H_W^2 + q_l)^{-1}$: calculable simultaneously
- Valid for $|\lambda|$ (eigenmode of H_W) $\in [\lambda_{thrs}, \lambda_{max}]$
- Projecting out low-modes of H_W below $\lambda_{thrs} \rightarrow \text{sign}(\lambda)$ ($\lambda < \lambda_{thrs}$)
- Cost depends on the low-mode density
($\lambda_{thrs}=0.045$, $N=10$ in this work)



Explicitly calculated





Locality

Fermion operator should be local

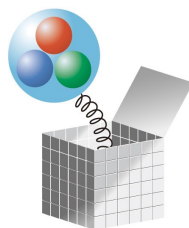
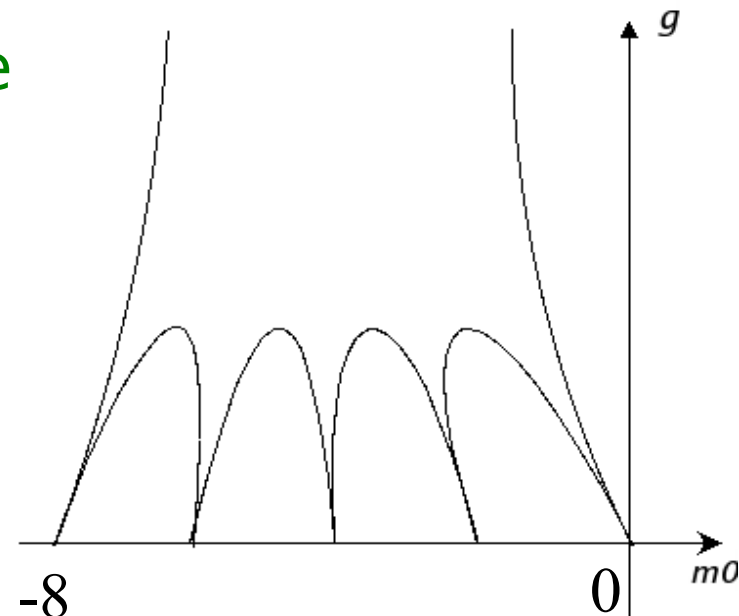
- Overlap operator is exponentially local, if
 - No low mode of H_W below some threshold

Hernandez, Jansen, Luscher, 1999

Near-zero mode is itself exponentially local

Golterman, Shamir, 2003; Golterman, Shamir, Svetitsky, 2003

⇔ Out of Aoki phase
(parity broken phase)

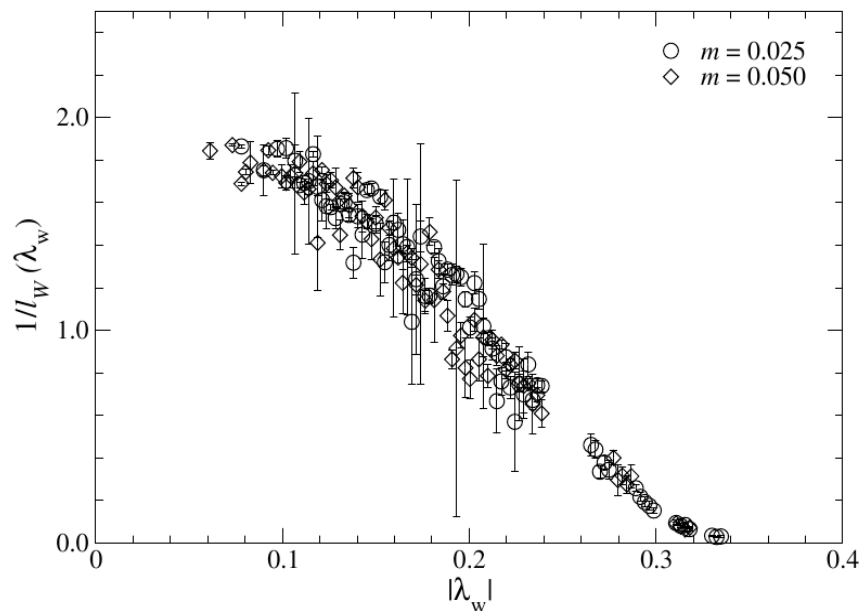




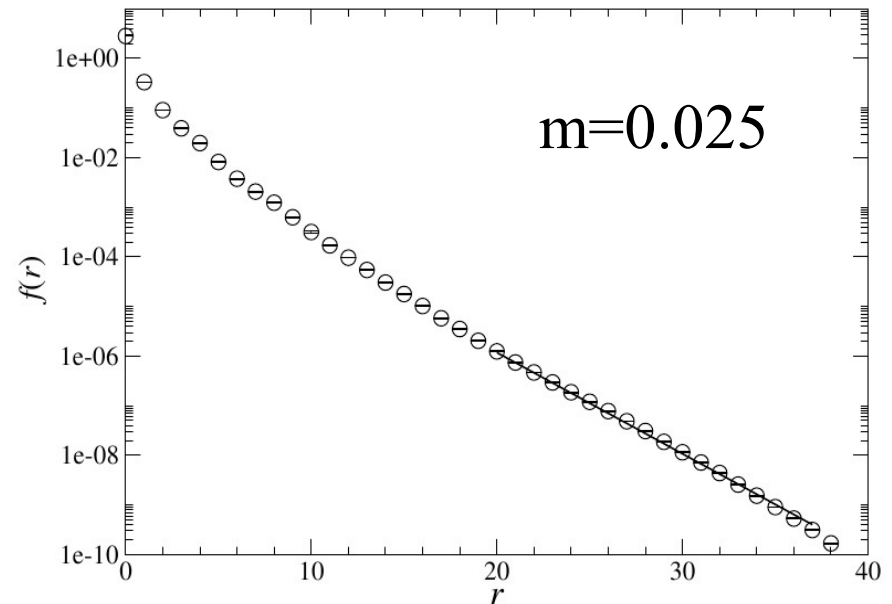
Locality

JLQCD, 2008; JLQCD (Yamada et al.), Proc. of Lattice 2006

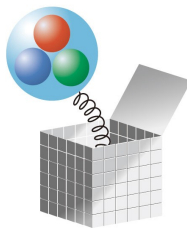
- At beta=2.3 (Nf=2)



Localization length of low eigenmodes of H_W



Localization of overlap operator
 $l = 0.25\text{fm} (\sim 2a)$

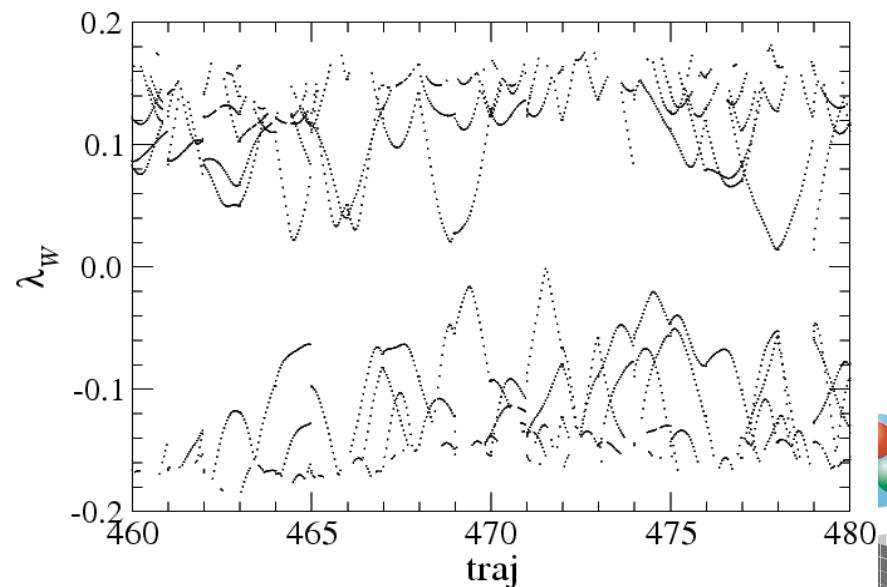
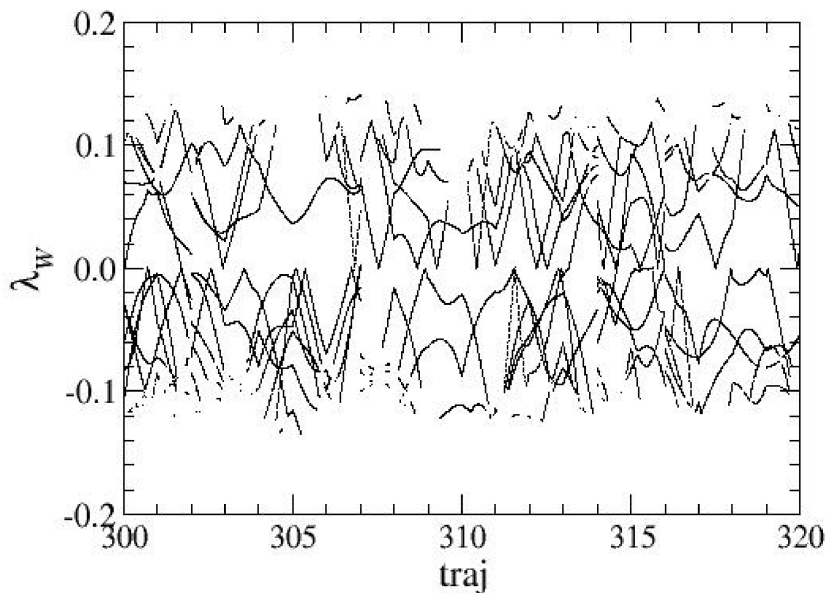




Simulation at fixed topology

- Suppressing near-zero modes of H_W
 - Locality of overlap operator
 - Cost of dynamical simulation reduced
- Achieved by extra-Wilson fermion term
 - Twisted mass ghost: suppress high mode effect

$$\det \left(\frac{H_W^2}{H_W^2 + \mu^2} \right) = \int \mathcal{D}\chi^\dagger \mathcal{D}\chi \exp[-S_E]$$



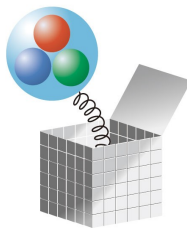


Simulation at fixed topology

Forbidding near-zero mode of $H_W \Leftrightarrow$ fixing topology

Is fixing topology a problem ?

- In the infinite V limit,
 - Fixing topology is irrelevant
 - Local fluctuation of topology is active
- In practice, V is finite
 - Topology fixing \Rightarrow finite V effect
 - $\theta=0$ physics can be reconstructed
 - Finite size correction to fixed Q result (with help of ChPT)
 - Must check local topological fluctuation
 - \Rightarrow topological susceptibility, η' mass
 - Remaining question: Ergodicity ?





Physics at fixed topology

One can reconstruct fixed θ physics from fixed Q physics
(Bowler et al., 2003, Aoki, Fukaya, Hashimoto, & Onogi, 2007)

- Partition function at fixed topology

$$Z_Q = \frac{1}{2\pi} \int_{-\pi}^{\pi} Z(\theta) \exp(i\theta Q) \quad \Leftrightarrow \quad Z(\theta) = \sum_Q Z_Q \exp(-i\theta Q)$$

- For $Q \ll \chi_t V$, Q distribution is Gaussian

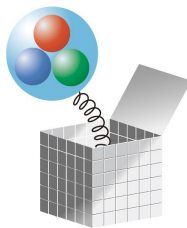
- Physical observables

- Saddle point analysis

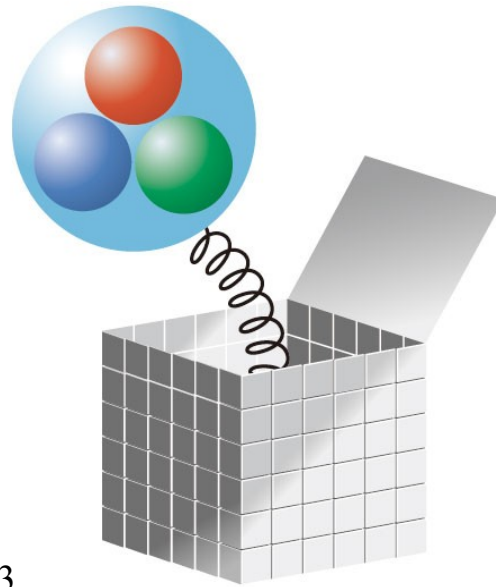
$$\Rightarrow \langle O \rangle_\theta = \langle O \rangle_Q + (\text{finite } V \text{ correction}) \quad \text{for } Q \ll \chi_t V$$

- Example: pion mass

$$m_\pi^Q = m_\pi(\theta = 0) + \frac{1}{2V\chi_t} \left(1 - \frac{Q^2}{V\chi_t} \right) \frac{\partial^2 m_\pi(\theta)}{\partial \theta^2} \Big|_{\theta=0} + O(V^{-2})$$



Large-scale simulation of overlap fermion



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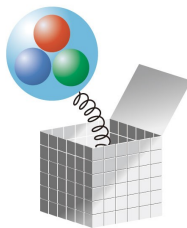
Collaboration

- JLQCD Collaboration

- S.Hashimoto, T.Kaneko, J.Noaki, E.Shintani, N.Yamada, K.Takeda, H.Ikeda, H.M. (KEK)
- S.Aoki, K.Kanaya, Y.Kuramashi, Y.Taniguchi, A.Ukawa, T.Yoshie (Tsukuba Univ)
- H.Fukaya (Niels Bohr Institute)
- T.Onogi, H.Ohki (YITP, Kyoto Univ)
- K-I.Ishikawa, M.Okawa (Hiroshima Univ)

- TWQCD Collaboration

- T-W.Chiu, K.Ogawa (Natl.Taiwan Univ)
- T-H.Hsieh (RCAS, Academia Sinica)





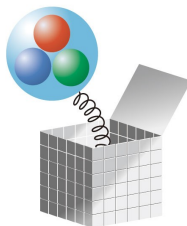
Project goals

Dynamical overlap simulations

- Exact chiral symmetry
- Small enough quark masses, ε -regime
- Since 2006 --- on new KEK system (x50 upgraded)
- Let us try new formalism!

Goals:

- Exploring the chiral regime
 - Confirming the chiral symmetry breaking scenario
 - Testing the effective chiral Lagrangian predictions
- Matrix elements with controlled chiral extrapolation
 - Without artificial operator mixing
 - Precision computation for flavor physics





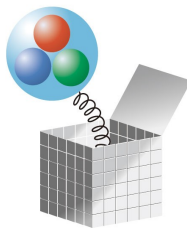
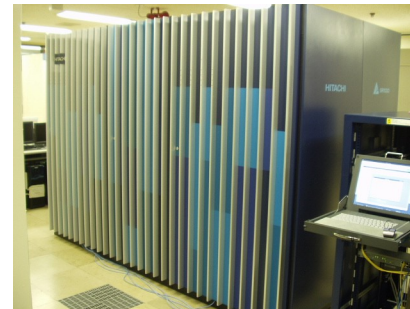
Machines

Main machine: IBM Blue Gene/L at KEK

- 57.6 Tflops peak (10 racks)
- 8x8x8(16) torus network



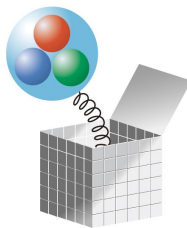
- Hitachi SR11000 (KEK)
 - 2.15TFlops/0.5TB memory
- NEC SX8 (YITP, Kyoto)
 - 0.77TFlops/0.77TB memory





Runs

- Nf=2: 16³x32, a=0.12fm (production run finished)
 - 6 quark masses covering (1/6~1) m_s
 - 10,000 trajectories with length 0.5
 - 20-60 min/traj on BG/L 1024 nodes
 - Q=0, Q=-2, -4 ($m_{sea} \sim m_s/2$)
 - ϵ -regime ($m_{sea} \sim 3\text{MeV}$)
- Nf=2+1 : 16³x48, a=0.11fm (production run finished)
 - 2 strange quark masses around physical m_s
 - 5 ud quark masses covering (1/6~1) m_s
 - 2500 trajectories with length 1
 - About 2 hours/traj on BG/L 1024 nodes
- Nf=2+1 : 24³x48 (just started)



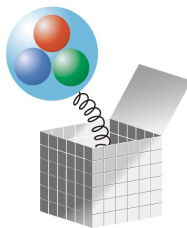
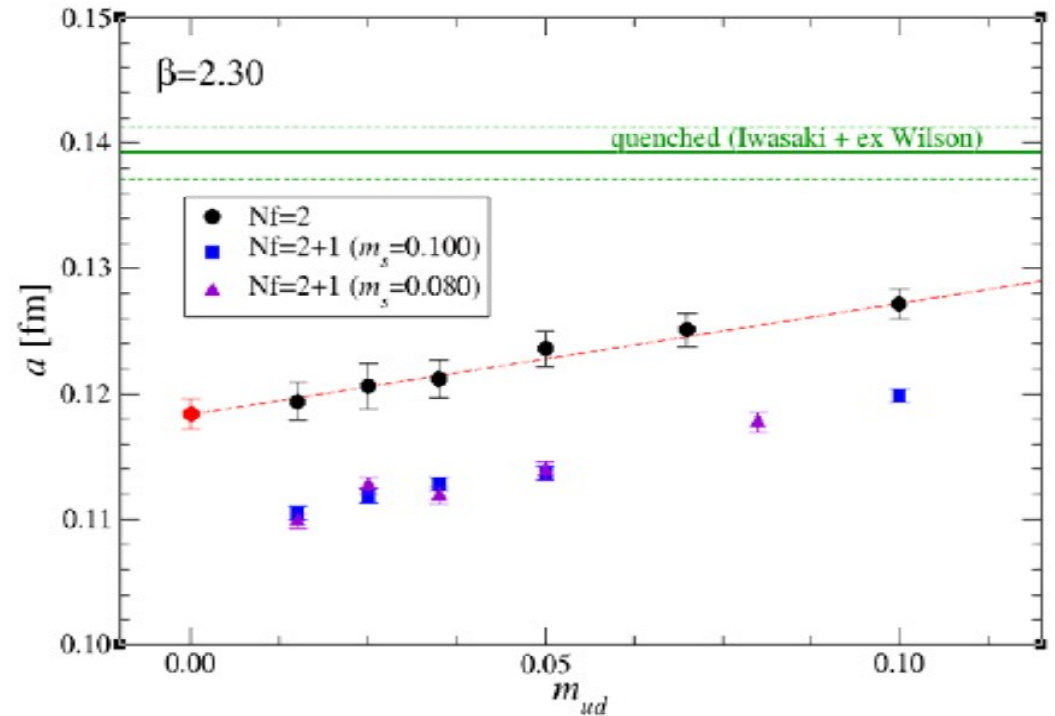
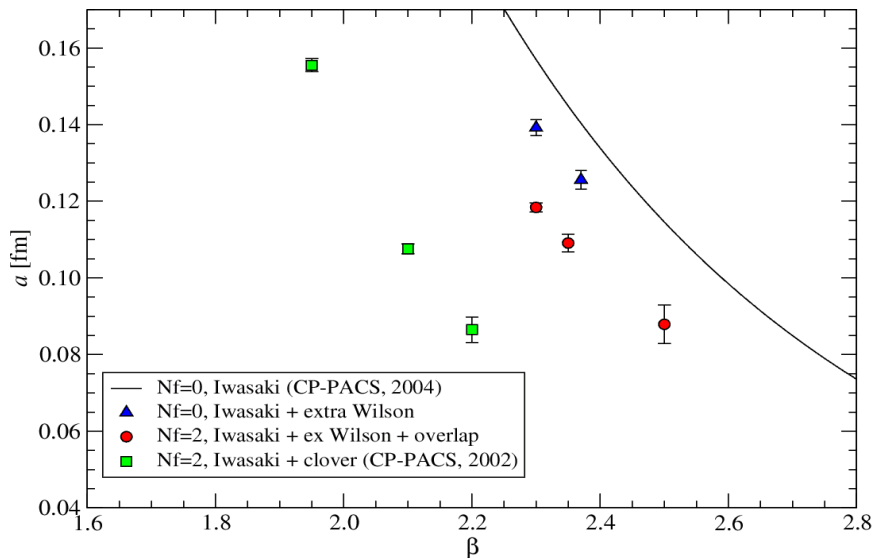


Lattice scale

- Scale: set by $r_0 = 0.49\text{fm}$
 - Static quark potential

$$r^2 \left. \frac{\partial V(r)}{\partial r} \right|_{r=r_0} = 1.65$$

- Milder β -shift than Wilson-type fermion





Chiral condensate

- Banks-Casher relation (Banks & Casher, 1980)

$$\Sigma = \langle \bar{q}q \rangle = \lim_{m \rightarrow 0} \lim_{V \rightarrow \infty} \frac{\pi \rho(0)}{V}$$

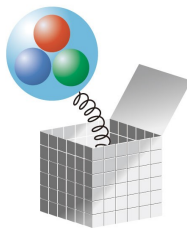
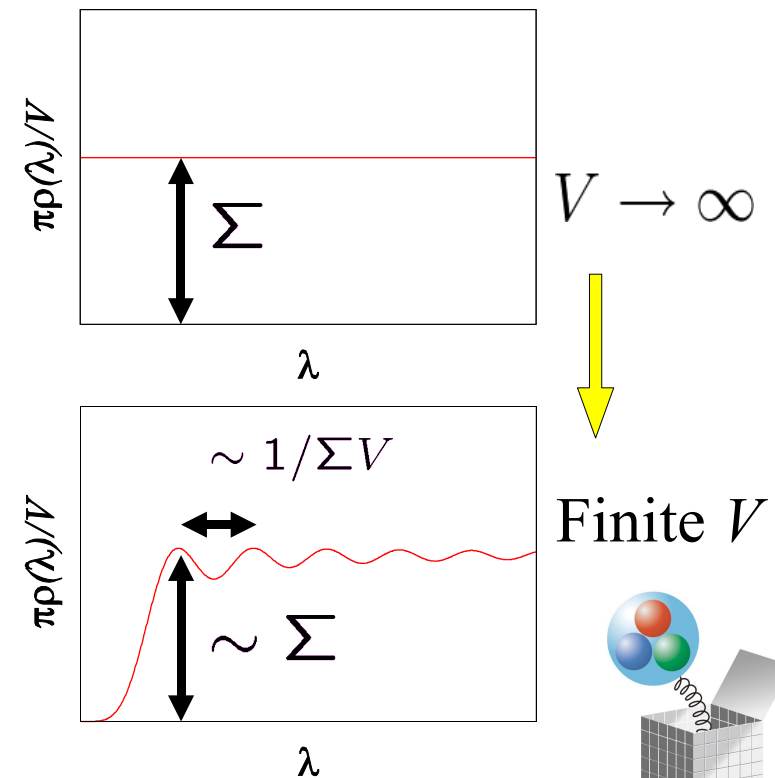
$$\rho(\lambda) = \sum_k \langle \delta(\lambda - \lambda_k) \rangle : \text{spectral density of } D$$

- Accumulation of low modes \iff Chiral SSB
- $V \rightarrow \infty$, then $m \rightarrow 0$

- ϵ -regime: $m \ll 1/\Sigma V$ at finite V

$$1/\Lambda_{QCD} \ll L \ll 1/m_\pi$$

- Low-energy effective theory
- Q -dependence is manifest
- Random Matrix Theory (RMT)





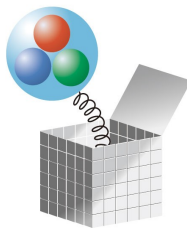
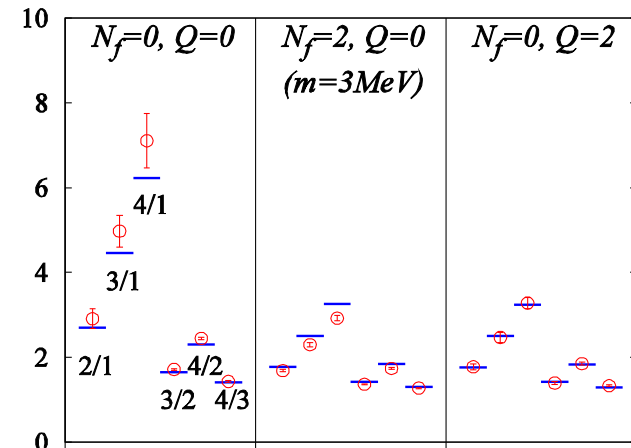
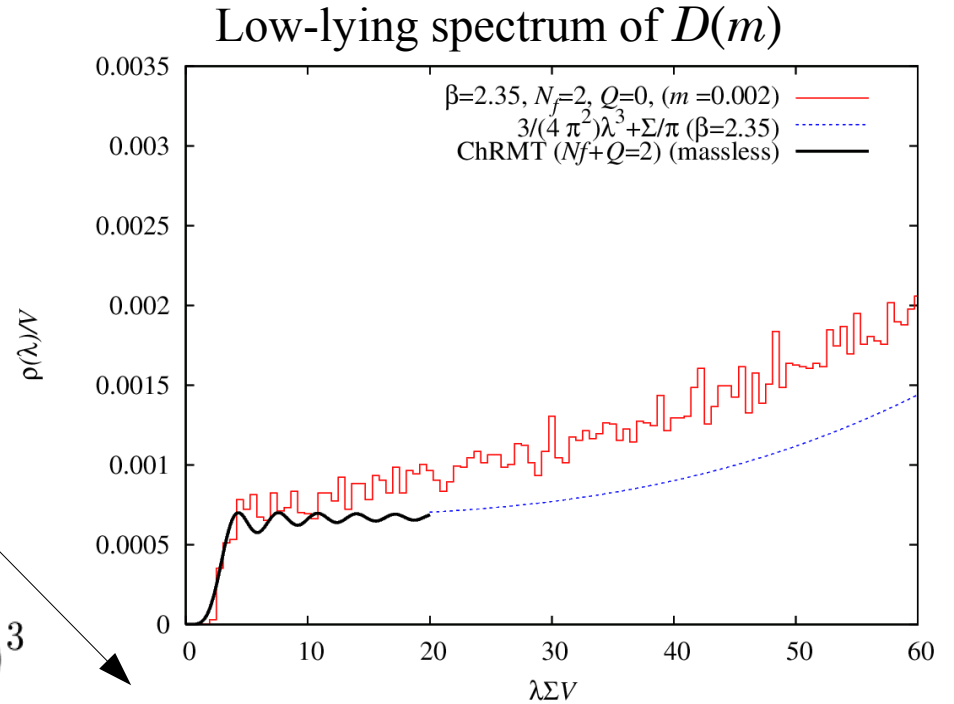
Result in the ε -regime

(JLQCD, 2007, JLQCD and TWQCD, 2007)

- $N_f=2, 16^3 \times 32, a=0.11\text{fm}$
- $m \sim 3\text{MeV}$
- Good agreement with RMT
 - lowest level distrib. $\rightarrow \Sigma$
 - Flavor-topology duality
- Chiral condensate:
 - Nonperturbative renorm.

$$\Sigma^{\overline{MS}}(2\text{ GeV}) = (251 \pm 7(\text{stat}) \pm 11(\text{syst}) \text{ MeV})^3$$

$O(\varepsilon^2)$ effect: correctable by meson correlator





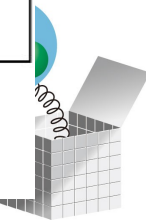
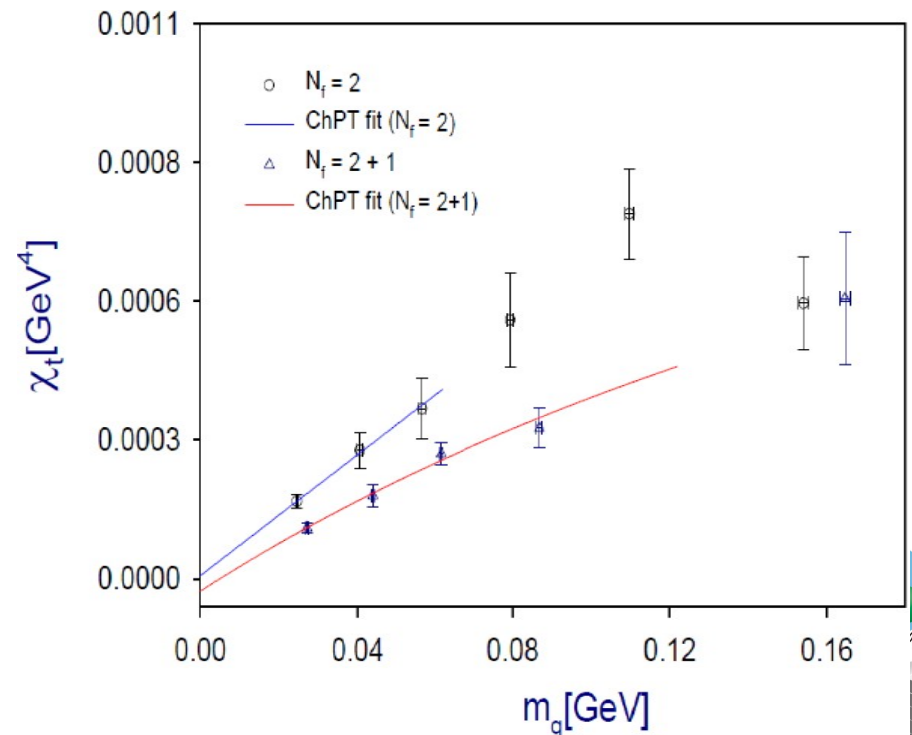
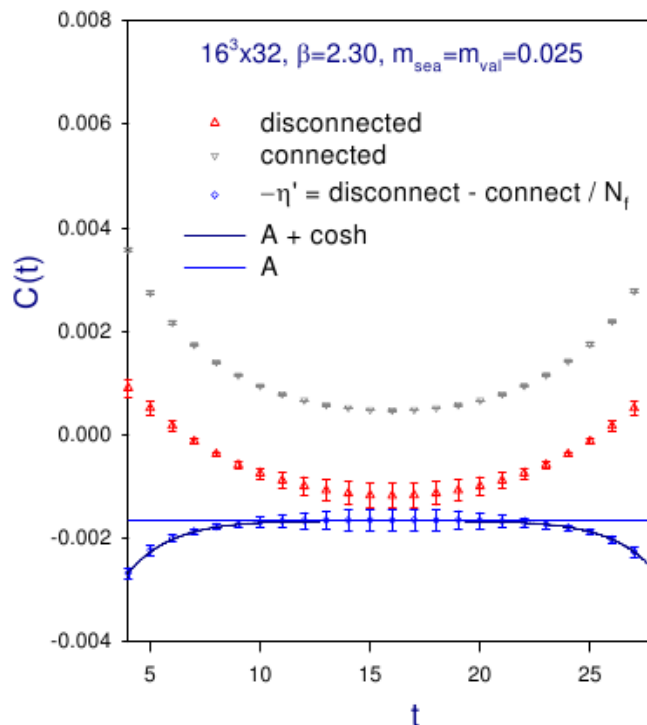
Topological susceptibility

Is local topological fluctuation sufficient? (JLQCD-TWQCD, 2007)

- Topological susceptibility χ_t can be extracted from correlation functions (Aoki et al., 2007)

$$\frac{1}{L^3} \langle m P^0(\vec{x}, t) m P^0(\vec{0}, 0) \rangle_Q \xrightarrow{t \gg 1} \frac{1}{V} \left[\frac{Q^2}{V} - \chi_t - \frac{c_4}{2\chi_t V} \right] + O(V^{-3}) + O(e^{-m_\eta t})$$

where
$$P^0 \equiv \frac{1}{N_f} \sum_{f=1}^{N_f} \bar{q}_f \gamma_5 \left(1 - \frac{D}{2m_0} \right) q_f$$





Meson spectrum: finite volume effect

Finite volume correction

JLQCD, arXiv:0804.0894

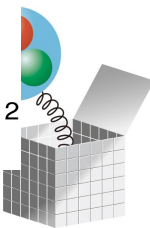
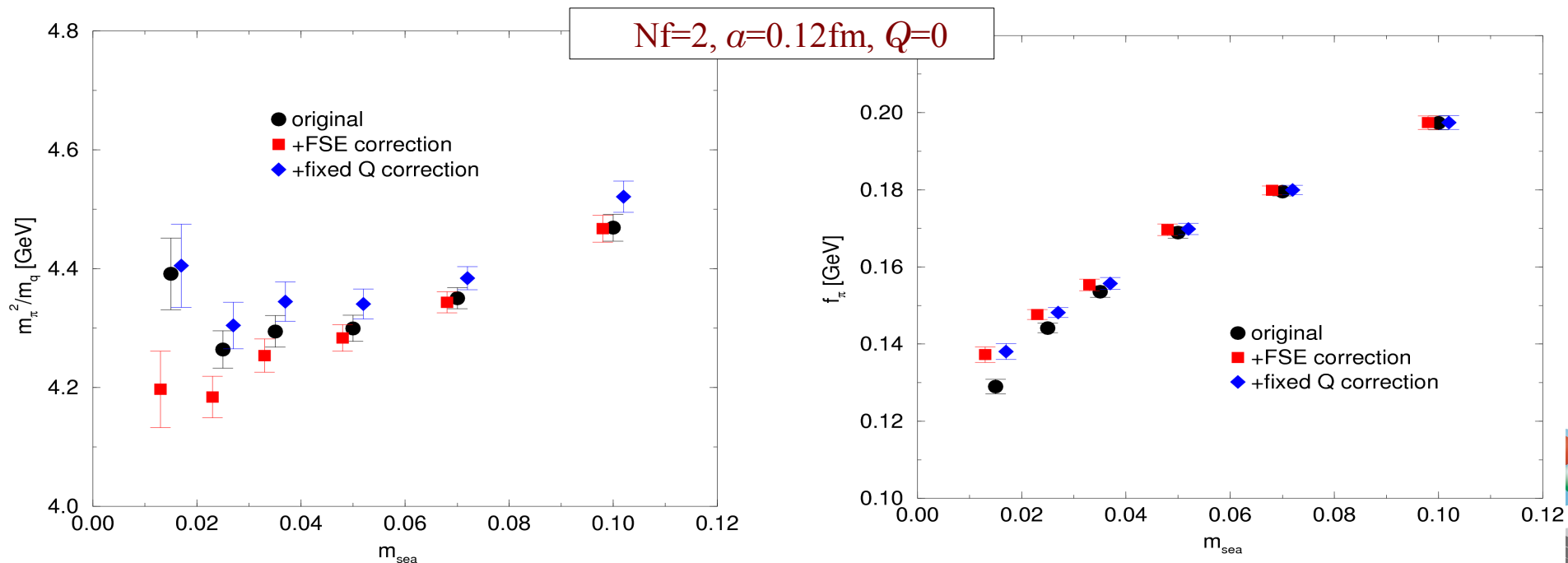
$$(m_\pi^2)^{\text{corrected}} = \frac{m_\pi^2}{(1 + R_m)^2(1 + T_m)^2}, \quad (f_\pi)^{\text{corrected}} = \frac{f_\pi}{(1 + R_f)(1 + T_f)}$$

- ***R***: ordinary finite size effect

Estimated using two-loop ChPT (Colangelo et al, 2005)

- ***T***: Fixed topology effect (Aoki et al, 2007)

– At most 5% effect --- largely cancel between R and T



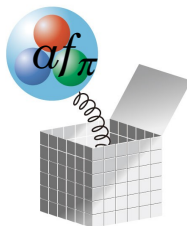
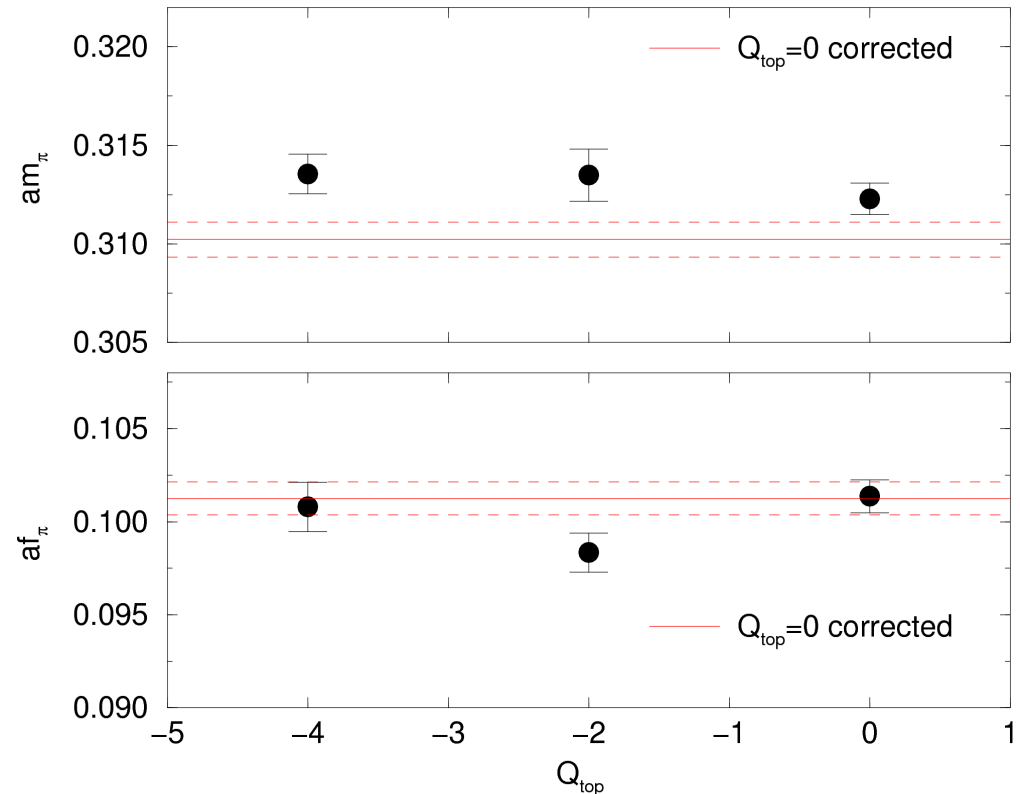


Meson spectrum: Q-dependence

$$N_f=2, m_{sea} = 0.050$$

$$Q = 0, -2, -4$$

- No large Q-dependence
(consistent with expectation)





Meson spectrum: ChPT test

Chiral expansion

$N_f=2, \alpha=0.12\text{fm}, Q=0$

JLQCD, arXiv:0804.0894

- Continuum ChPT formula applicable
- The region of convergence

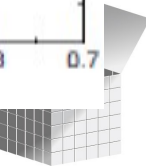
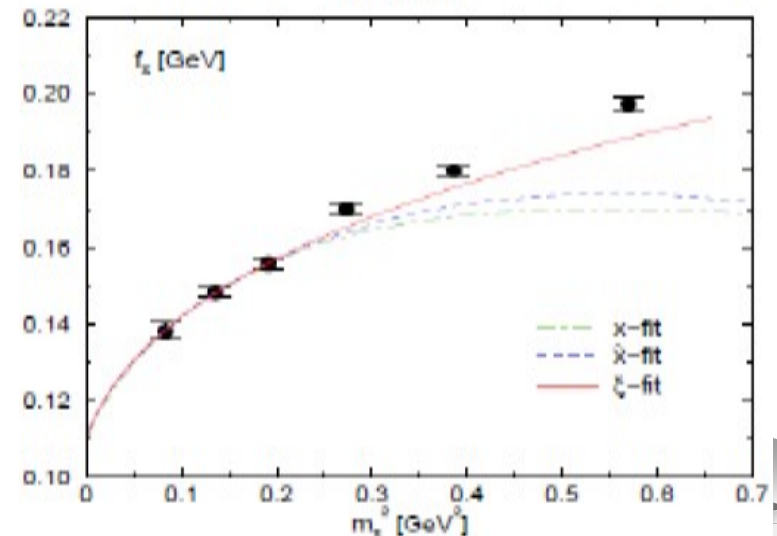
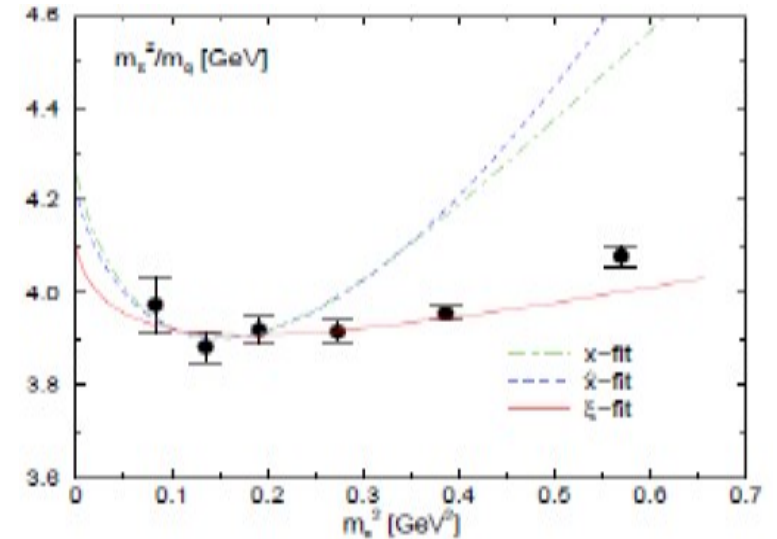
$$\frac{m_\pi^2}{m_q} = 2B[1 + x \ln(x) + c_3x + O(x^2)]$$

$$f_\pi = f[1 - 2x \ln(x) + c_4x + O(x^2)]$$

Expand either

$$x \equiv \frac{m^2}{(4\pi f)^2}, \quad \hat{x} \equiv \frac{m_\pi^2}{(4\pi f)^2}, \quad \xi \equiv \frac{m_\pi^2}{(4\pi f_\pi)^2}$$

ξ extends the region significantly



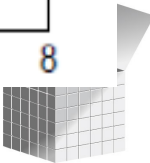
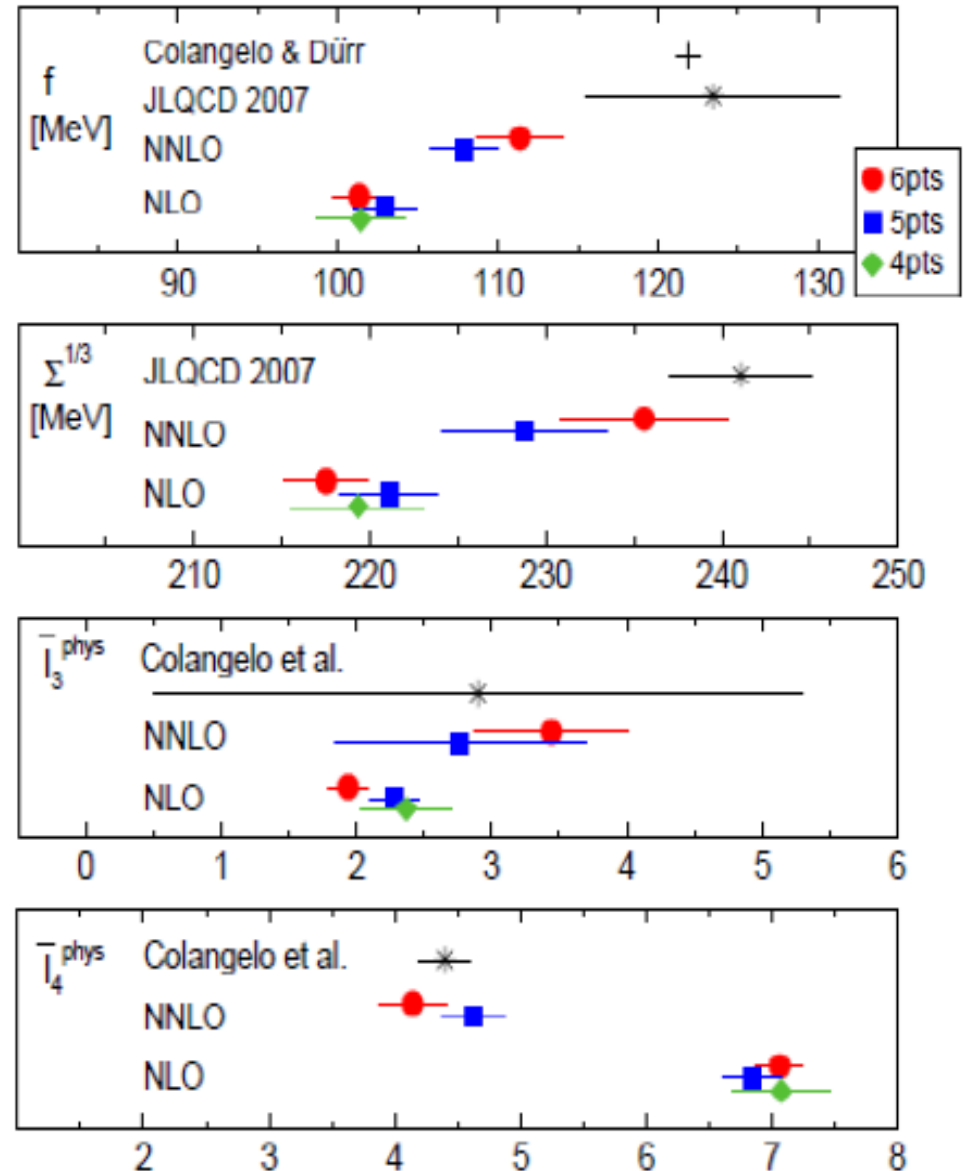


Meson spectrum: low energy consts

Low energy constants

--- NNLO fit with ξ

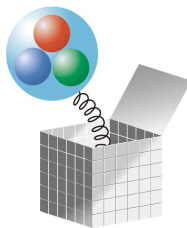
For reliable extraction of low energy constant, NNLO terms are mandatory



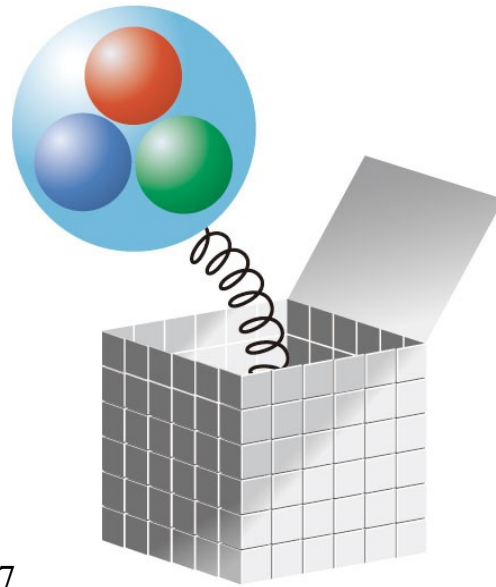


Other results

- Physics publication: Cf. <http://jlqcd.kek.jp/>
 - ε -regime
 - Topological susceptibility
 - Meson spectroscopy and ChPT test
 - $\pi^+-\pi^0$ mass difference
 - Coupling const
 - B_K (Kaon bag parameter)
 - Nucleon sigma term
 - Ps-NG boson mass/S-parameter
 - Pion form factors (in preparation)
 - Nonperturbative renormalization (in preparation)



Toward overlap simulation at finite temperature

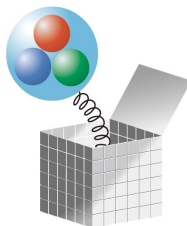


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Lattice QCD at $T > 0$

- Phase diagram: still not conclusive
 - $N_f=2$: expected $O(4)$ scaling not observed in KS
 - All recent works are with KS, $N_t=4$
 - Work with Wilson fermion is not at enough small quark mass
 - $N_f=2+1$ (physical point)
 - Recently only with KS fermions, still large uncertainty
- Finite chemical potential
 - various methods proposed, but limited applicability
 - still waiting a breakthrough!
- Thermodynamics (equation of state, etc.)
- Excitation spectrum
- Viscosity





Toward overlap at $T > 0$

Where overlap should be applied ?

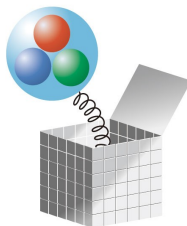
- Chiral symmetry is essential
- KS and Wilson (and others) are inconsistent

Phase structure of $N_f=2, 2+1$

- Vicinity of phase transition --- critical behavior ?
- Observables concerning chiral symmetry and topological charge

For other topics,

- Overlap costs too much
- Other formulations are sufficient
- Technical development is demanded





Phase diagram

Now popular phase diagram, but;

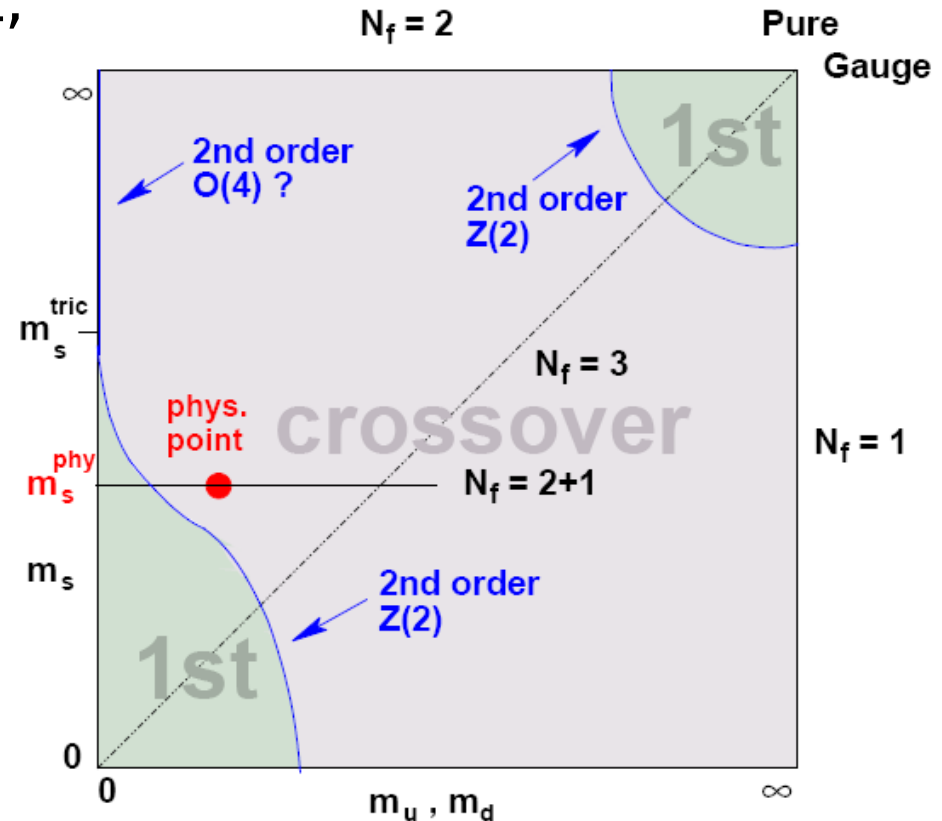
- How reliable?
- Consistency check enough?

$N_f=2$

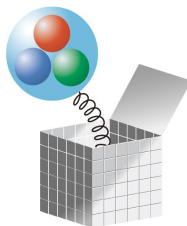
- KS fermion does not exhibit expected $O(4)$ scaling
- Wilson quark shows $O(4)$, but at rather heavy masses
- Most recent works are by KS, $N_t=4$.

$N_f=2+1$ (physical point)

- Really crossover? [(old) Wilson result is of 1st order]
- Recently only with KS fermions, still large uncertainty



(DeTar, Lattice 2008)





Phase diagram

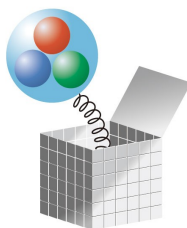
Nf=2 scaling study

- Needs finite size scaling
- Good reference work:

[D'Elia, Di Giacomo, Pica, Phys. Rev. D72 \(2005\) 114510](#)

Lattice simulation with

- Staggered (unimproved)
- $L_t=4, L_s=12, 16, 20, 24, 32 \rightarrow$ finite size scaling
- $am=0.01335 - 0.307036$
- Measure specific heat, susceptibility
- Comparison with $O(4), O(2),$ 2st order scaling

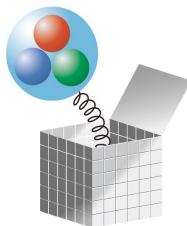
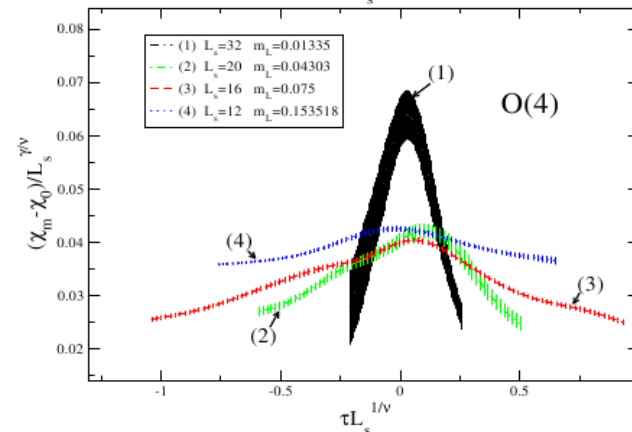
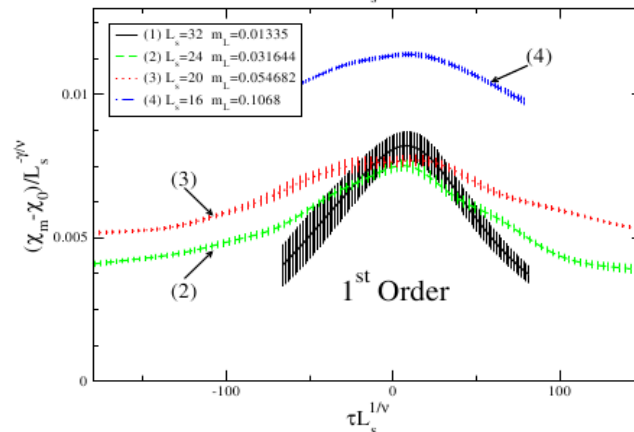
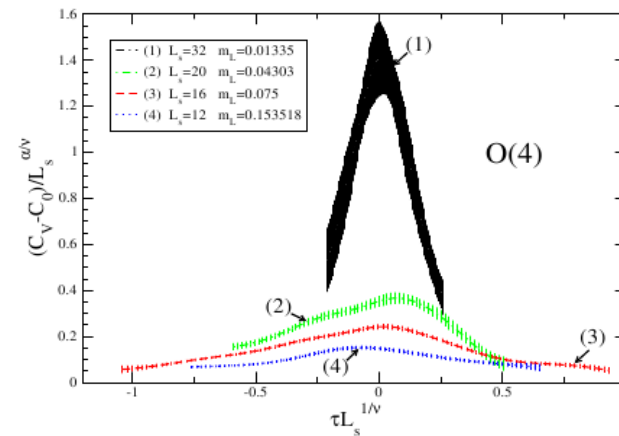
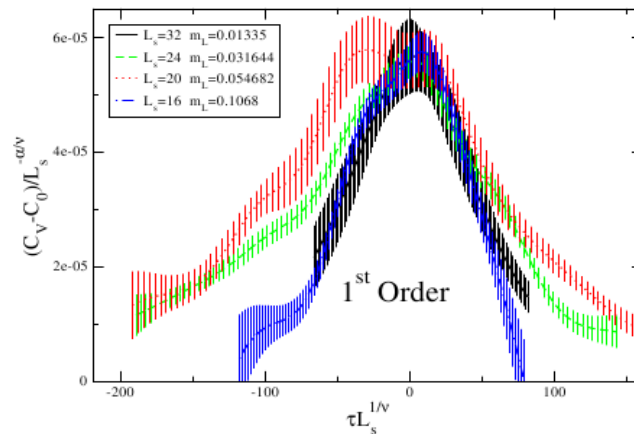




Nf=2 scaling study

Result (Cossu, D'Elia, Di Giacomo, Pica, PoS(LAT2007) 219):

- Inconsistent with O(4) nor O(2)
- Prefers 1st order (but not all quantities are consistent)

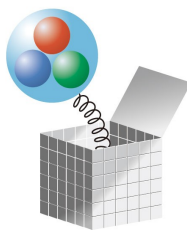




Toward overlap at $T > 0$

To perform the same level of work with overlap fermion,

- Suppose that simulation is possible at $Nt=6$
 - Locality must be kept
 - Not clear whether possible or not
 - At least, around $Ls=32$ $O(10000)$ trj x $O(20)$ param. sets
 - 1-2 years on Blue Gene @KEK (57.3TFlops)
 - Whole project 2-3 years (**optimistic estimate**)
- **At present, not realistic**
 - If performance is 5 times improved, may possible
 - Next generation computer? (time passes so fast!)
 - Within a few years, become feasible

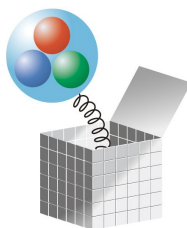




Toward overlap at $T > 0$

Preparation should be started

- How $1/a$ can be reduced keeping locality ?
- Fixed topology effect
 - For matrix elements, finite size effect
 - How about thermodynamic quantities, in particular critical exponent ?
- Survey in parameter space (beta, quark mass, volume)
- Fundamental thermodynamic quantities
- Improved algorithms (always called for!)





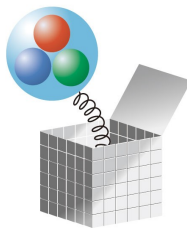
Beyond QCD

Motivation:

- What is relation between confinement and chiral symmetry breaking?
- Phase structure of gauge theories
- Advantages of exact chiral symmetry

Beyond QCD applications

- Large N_f
- Non-fundamental representations (adjoint, etc.)
- N_c not 3
- Confinement and broken chiral symmetry may not occur simultaneously





Summary/outlook

- Overlap fermion has elegant chiral structure
- Numerical cost is high
- JLQCD is performing large dynamical overlap project at $N_f=2$ and $2+1$
- Rich physics results are being produced

Outlook

- Simulation at finite temperature is challenging
- Scaling study ?
- Should start exploratory study
- Beyond QCD simulations are interesting

