

# *Numerical study of staggered fermion on anisotropic lattices*

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## Introduction(1)

### Anisotropic lattice:

Fine resolution in temporal direction  $a_\tau < a_\sigma$

### Advantages

□ Large number of d.o.f. in  $t$  direction

⇒ Large number of Matsubara frequencies

\* Finite temperature

Equation of state

cf. Y.Namekawa et al., Phys. Rev. D64 (2001) 074507

Analysis of correlation functions

cf. QCD-TARO Collaboration, Phys. Rev. D63 (2001) 054501

cf. T.Umeda, K.Nomura, H.Matsufuru, hep-lat/0211003

□ Large cutoff in  $t$  direction

\* Heavy quark systems

cf. J.Harada et al., Phys. Rev. D64 (2001) 074501

□ Easy to follow the change of correlators

\* Glueball spectrum

cf. C.J.Morningstar et al., Phys. Rev. D60 (1999) 034509

cf. N.Ishii et al., Phys. Rev. D66 (2002) 014507

\* Negative parity baryon spectrum

cf. Y.Nemoto et al., hep-lat/0302013

## Introduction(2)

### Disadvantages

- Need to tune anisotropy parameters nonperturbatively (  $\xi \equiv a_\sigma/a_\tau$  )

Observable suffer from calibration uncertainties.

$\implies$  Must be tuned with high precision

We investigate which calibration procedure provides good precision and good scaling behavior.

Our goal is to tune anisotropy parameters within 1 % statistical error on dynamical lattices.

Target anisotropy:  $\xi = 4$

## Form of action

### □ Wilson gauge + staggered fermion actions

Previous work ( $\gamma_G = \gamma_F$ )

L.Levkova et al., Nucl.Phys.Proc.Suppl.106 (2002)218

⇒ Precision is not enough

Gauge sector:

$$S_G = \beta \sum_x \left[ \sum_{i < j}^3 \frac{1}{\gamma_G} \left( 1 - \frac{1}{3} \text{Re tr} W_{ij}(x) \right) + \sum_i^3 \gamma_G \left( 1 - \frac{1}{3} \text{Re tr} W_{i4}(x) \right) \right]$$

$\gamma_G$ : bare gauge anisotropy parameter

Quark sector:

$$S_F = \sum_{x,y} \bar{\chi}(x) K(x,y) \chi(y)$$

$$K(x,y) = \delta_{x,y} - \gamma_F \kappa_\sigma \eta_{x4} \left[ U_4(x) \delta_{x,y-\hat{4}} - U_4^\dagger(x-\hat{4}) \delta_{x,y+\hat{4}} \right] \\ - \kappa_\sigma \sum_i \eta_{x\mu} \left[ U_i(x) \delta_{x,y-\hat{i}} - U_i^\dagger(x-\hat{i}) \delta_{x,y+\hat{i}} \right]$$

$$\kappa_\sigma = 1/2m_0$$

$\gamma_F$ : bare fermionic anisotropy parameter

$m_0$ : bare quark mass (in spacial lattice units)

## Tuning of anisotropy parameters (1)

In a simulation on the anisotropic lattice,  
we need to find  $\gamma_G = \gamma_G^*$  and  $\gamma_F = \gamma_F^*$  which satisfy

$$\xi_G(\gamma_G^*, \gamma_F^*) = \xi_F(\gamma_G^*, \gamma_F^*) = \xi$$

for each  $(\beta, m_0)$ .

Renormalized anisotropies

$\xi_G$  : defined with gauge observable

$\xi_F$  : defined with fermion observable

$\xi_G, \xi_F$  must be determined in good accuracy.

□ In quenched simulation

One can first tune  $\gamma_G$  independently of  $\gamma_F$ .

□ In dynamical simulation

We have to tune  $\gamma_G$  and  $\gamma_F$  simultaneously.

## Tuning of anisotropy parameters (2)

We adopt following definitions of  $\xi_G$  and  $\xi_F$

- $\xi_G$ : defined with hadronic radius  $r_0$   
 $\Rightarrow$  Good scaling behavior

Precise static potentials in  $x$  and  $t$ (fine) directions

- Quench: Lüscher-Weisz noise reduction technique

cf. Lüscher and Weisz JHEP 0109 (2001) 010

- Dynamical: smearing technique in  $(x, t)$  plane

$$\Rightarrow r_0 : r_0^2 F(r_0) = 1.65 \quad \Rightarrow \quad \xi_G = r_0^{(t)} / r_0^{(x)}$$

- $\xi_F$ : defined with ratio of meson masses  
in  $x$  and  $t$  directions  $\xi_F = m_h^{(z)} / m_h^{(t)}$

— Less statistical error

Dispersion relation:  $E^2(\vec{p}) = m^2 + \vec{p}^2 / \xi_F^2$

— used to check the systematic uncertainties

## Quenched lattice simulation

□ Gauge parameters (  $\xi = 4$  )

$\beta$	$\gamma_G$	size	$a_\sigma^{-1}$ (GeV)	$N_{conf}$
5.75	3.136	$12^2 \times 24 \times 96$ (fine)	1.100(6)	224

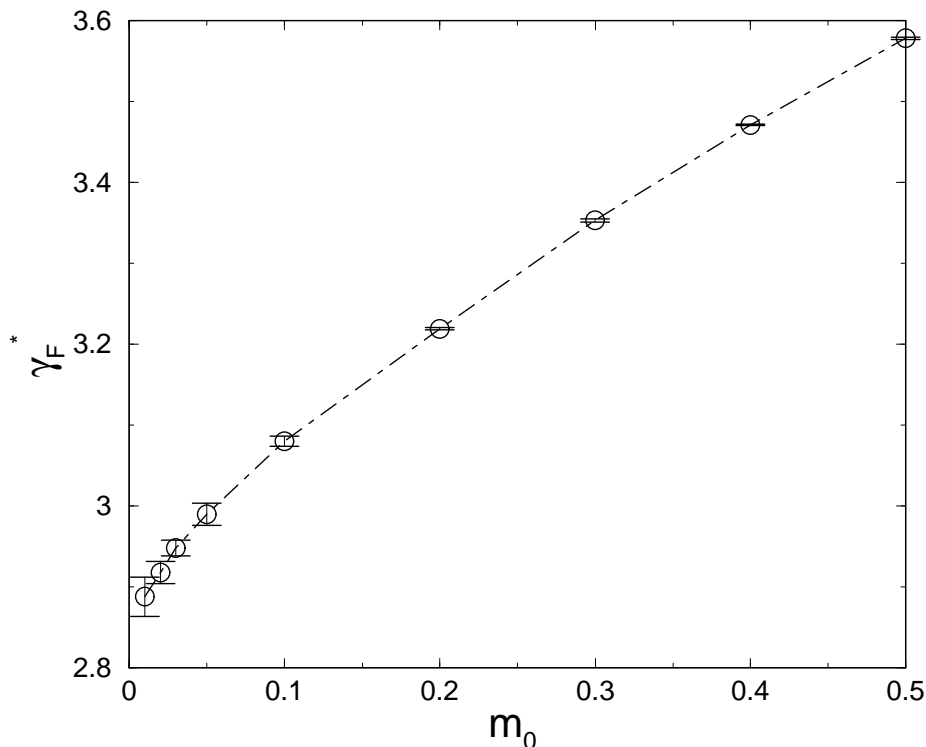
$\gamma_G$  determined with 0.3% error from hadronic radius  $r_0$

□ Quark parameters

$m_0=0.01\sim 0.5$  ( $m_{PS}/m_V = 0.33 \sim 0.90$ )

input  $\gamma_F$  : 2 ~ 4 points  $\Rightarrow$  interpolation to  $\gamma_F^*$

□ Simulation result



- Need to compare with free quark analysis
- Need to compare with result of dispersion relation

## $N_f = 2$ full QCD simulation

Lattice:  $10^2 \times 20 \times 80$ (fine),  $\xi = 4$ ,  $N_f = 2$

$\beta = 5.3 \sim 5.45$  ( $a_\sigma^{-1} = 0.69 \sim 0.85$  GeV)

Bare quark mass:  $m_0 = 0.1$  ( $m_{PS}/m_V \sim 0.6$ )

$N_{conf}$ : 300  $\sim$  600 trajectories with Hybrid R algorithm

Determination of anisotropy parameters ( $\gamma_G^*$ ,  $\gamma_F^*$ ):

4 $\sim$ 6 points of input ( $\gamma_G$ ,  $\gamma_F$ ) for each ( $\beta$ ,  $m_0$ )

tuning to  $\Rightarrow \xi = 4$

We suppose the following linear relations

$$\xi_G = a_0 + a_1(\gamma_G - \gamma_G^*) + a_2(\gamma_F - \gamma_F^*)$$

$$\xi_F = b_0 + b_1(\gamma_G - \gamma_G^*) + b_2(\gamma_F - \gamma_F^*)$$

$\chi^2$  fitting  $\Rightarrow (\gamma_G^*$ ,  $\gamma_F^*)$

□  $\beta = 5.3$  result

$$[a_0, a_1, a_2] = [4.000(24), 0.702(92), 0.075(92)]$$

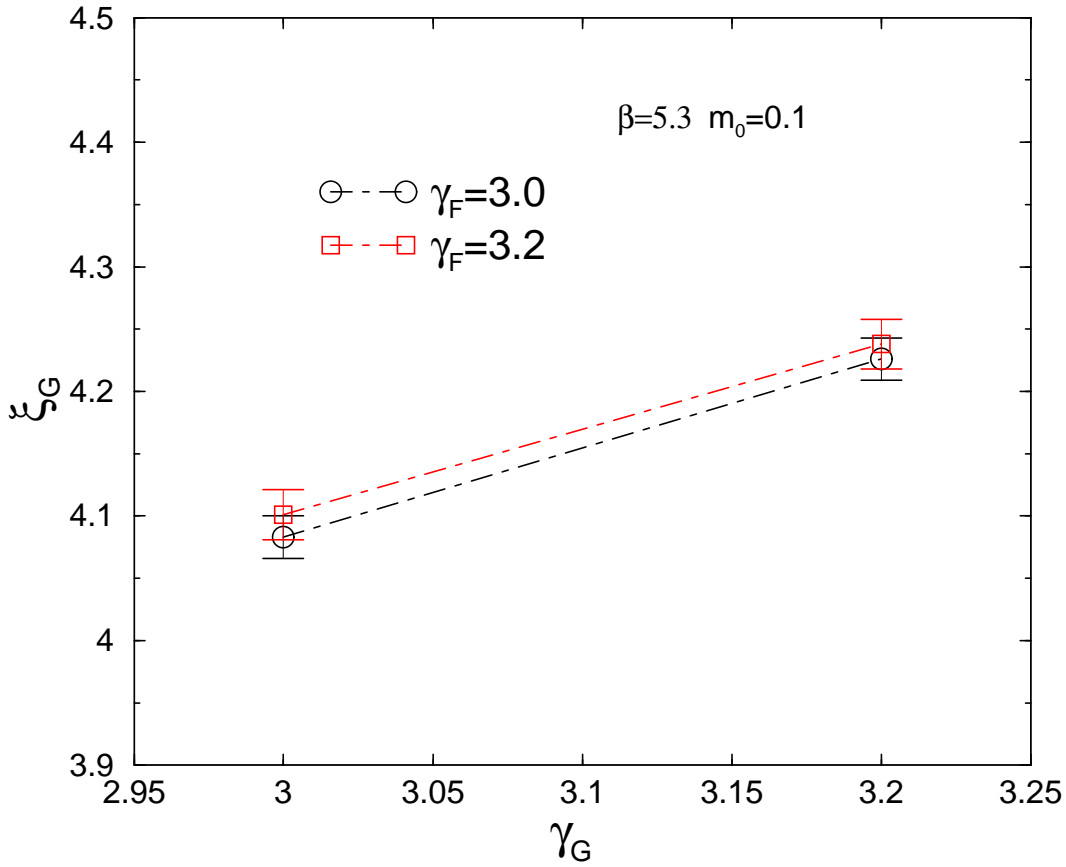
$$[b_0, b_1, b_2] = [4.000(28), 0.220(08), 1.225(08)]$$

$$\gamma_G^* = 2.885(35), \quad \gamma_F^* = 2.955(31)$$

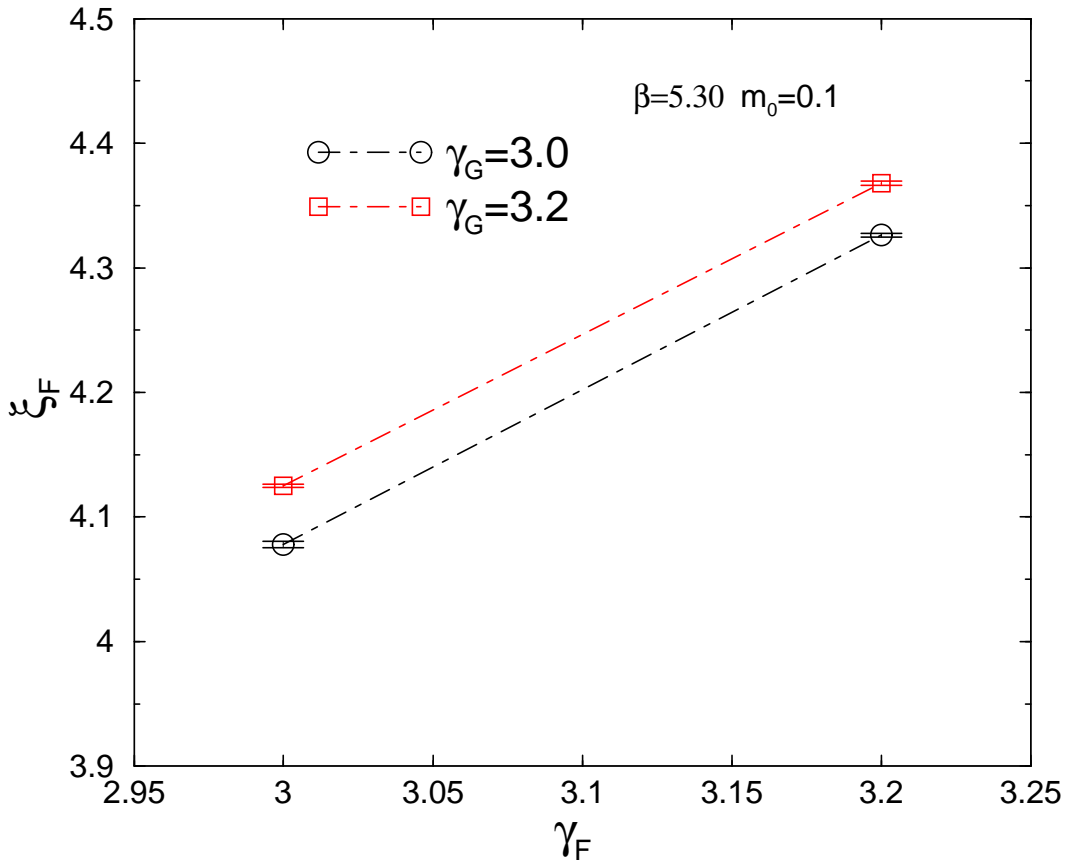
- $\gamma_F$  has small influence on  $\xi_G$
- $\gamma_G$  has small influence on  $\xi_F$
- Errors of  $\gamma_G$  and  $\gamma_F$  are about 1 %



# $\xi_G$ VS $\gamma_G$



# $\xi_F$ VS $\gamma_F$



## Conclusion

- We carried out the quenched and full QCD simulations on the anisotropic lattices with Wilson gauge + staggered fermion action
  
- \* Calibration is carried out using
  - gauge sector: hadronic radius  $r_0$
  - quark sector: ratio of masses in  $x$  and  $t$  directions
  
- Simulation result
  - We can perform calibration within 1 % error on both quenched and dynamical lattice.
  
- Outlook
  - Tuning in wide range of  $\beta$  and  $m_0$
  - Comparison with free quark analysis
  - Comparison with dispersion relation
  - Spectroscopy on anisotropic lattices
  - Application to finite temperature