

*Static Three Quark Potential
in the Quenched Lattice QCD*

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c.f. hep-lat/0006005

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Introduction (1)

Static potential between quarks inside hadrons

- $Q-\bar{Q}$ system (\sim meson)
- Well described with short range Coulomb term
+ long range confining linear term.

$$V(r) \simeq -\frac{A}{r} + \sigma r$$

String tension: $\sigma \simeq 1 \text{ GeV/fm}$

- $3Q$ system (\sim baryon)
- Short range: two body Coulomb potential
- Long range - string picture

(“Y-type” ansatz)

$$V_{3Q} = \sum_{(i,j)} V_{ij} + V_{123}$$

(“ Δ -type” ansatz)

$$V_{3Q} = \sum_{(i,j)} V'_{ij}$$

Introduction (2)

□ Flux tube picture prefers Y -type.

○ L_{min} (minimum length of Y -type string) $\simeq \frac{1}{2} l_{\Delta}$ ($\frac{1}{2} l_{\Delta} \leq l_{min} \leq \frac{1}{\sqrt{3}} l_{\Delta}$)
 \Rightarrow classically Y -type configuration is more stable.

○ Regge Pole analysis

\rightarrow Universality of string tension: $\sigma_{Q\bar{Q}} \simeq \sigma_{3Q}$

○ Strong coupling expansion

○ Quark Models (Nonrelativistic/Relativistic)

e.g. *Capstick and Isgur, PRD34(1986)2809:*

$$H = \sum_i \sqrt{\vec{p}_i^2 + m^2} - \sum_{(ij)} (V_{ij}^{Coulomb} + V_{ij}^{spin}) + \sigma L_{min}$$

well describes baryon spectra.

Δ -type ansatz with $\sigma_{\Delta} \simeq \sigma/2$ also acceptable.

□ Lattice results

(1) *Sommer and Wosiek, PLB149(1984)497, NPB267(1986)531*

(2) *Thacker, Eichten and Sexton, in Lattice'87*

(3) *Kammesberger et al.*, in Proc. of "Few-Body Problems in Particle, Nuclear, Atomic and Molecular Physics" (1987) 529

Without smearing technique (ground state enhancement)

Not sufficiently large time separation

(4) *Bali, hep-lat/0001312*

(1),(2) and (4) claim that their data support Δ -type ansatz.

However, their data also seem to be consistent

with Y -type ansatz.

— Still not conclusive.

Our Goals

- (1) Extract 3 quark potential from lattice QCD simulation.
 - Determine the parameters A and σ nonperturbatively.
 - discuss which of Δ and Y type is appropriate picture.
- (2) Apply obtained result to phenomenological study.
 - 3-body potential \rightarrow effective one-body potential
 - Model calculation \rightarrow baryon spectroscopy
- (3) $T > 0$ study
 - Near T_c , in confined/deconfined phase
 - Universality of $\sigma_{Q\bar{Q}}$ and σ_{3Q} ?
 - Baryon spectra at $T > 0$ (model calculation)
 - Comparison with lattice calculation (c.f. Umeda's talk)

Lattice Simulation (1): Baryonic Wilson loop

Static potential (potential between infinitely heavy quarks)

- Q - \bar{Q} potential \leftarrow Wilson loop

$$W(r, t) \equiv \text{Tr}$$

$$\longrightarrow c \cdot \exp[-V_{Q\bar{Q}}(r)t] \quad \text{as } t \rightarrow \infty$$

- $3Q$ potential \leftarrow Baryonic Wilson loop:

$$W(\vec{r}_1, \vec{r}_2, \vec{r}_3, t) \equiv \frac{1}{3!} \epsilon_{abc} \epsilon_{a'b'c'} U_1^{aa'} U_2^{bb'} U_3^{cc'}$$

=

$$\longrightarrow c' \cdot \exp[-V_{3Q}(\vec{r}_1, \vec{r}_2, \vec{r}_3)t] \quad \text{as } t \rightarrow \infty$$

In this limit, spatial path of links to connect three quark positions is unimportant.

Lattice Simulation (2): Simulation Parameters

□ Lattice

With standard Wilson action, in quenched approximation

○ $12^3 \times 24$, $\beta = 5.7$ (#conf = 210)

$a^{-1} \simeq 1.0 \text{GeV}$ (from static potential)

○ $16^3 \times 32$, $\beta = 6.0$ (#conf = 60, preliminary)

$a^{-1} \simeq 2.0 \text{GeV}$

Smearing – enhance the ground state contribution

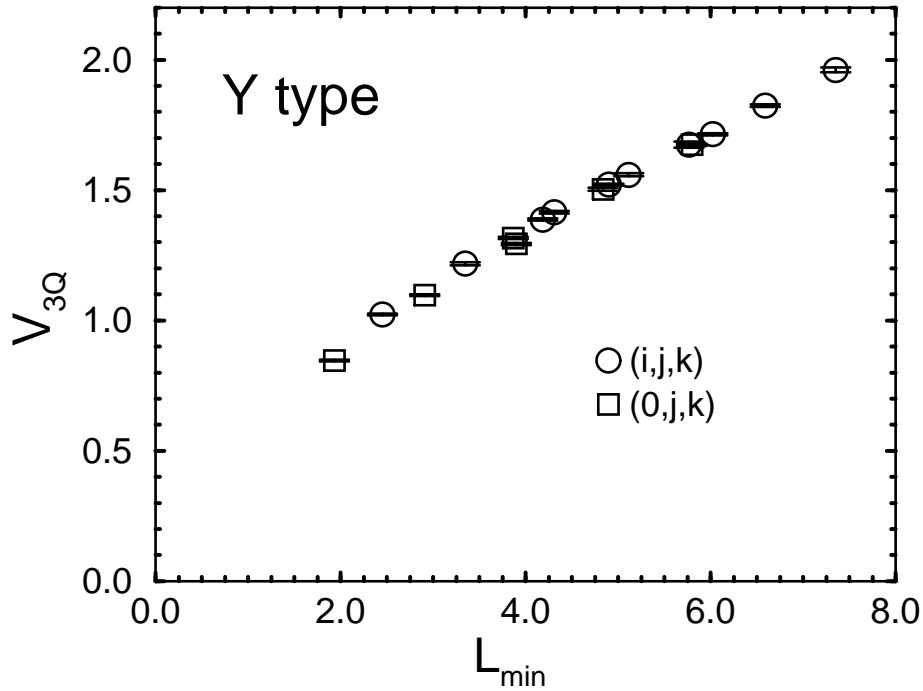
Iterate

— incorporate long range effect.

$\alpha = 2.3$, 20 sweeps for $Q-\bar{Q}$ / 12 sweeps for $3Q$ ($\beta = 5.7$)

Position of 3 quarks:

Lattice Simulation (3): $3Q$ Potential



Smearing: 12 sweeps

Fit:

$$V(\vec{r}_1, \vec{r}_2, \vec{r}_3) = \text{const.} - A \left(\frac{1}{r_{12}} + \frac{1}{r_{23}} + \frac{1}{r_{31}} \right) + \sigma l_{min}$$

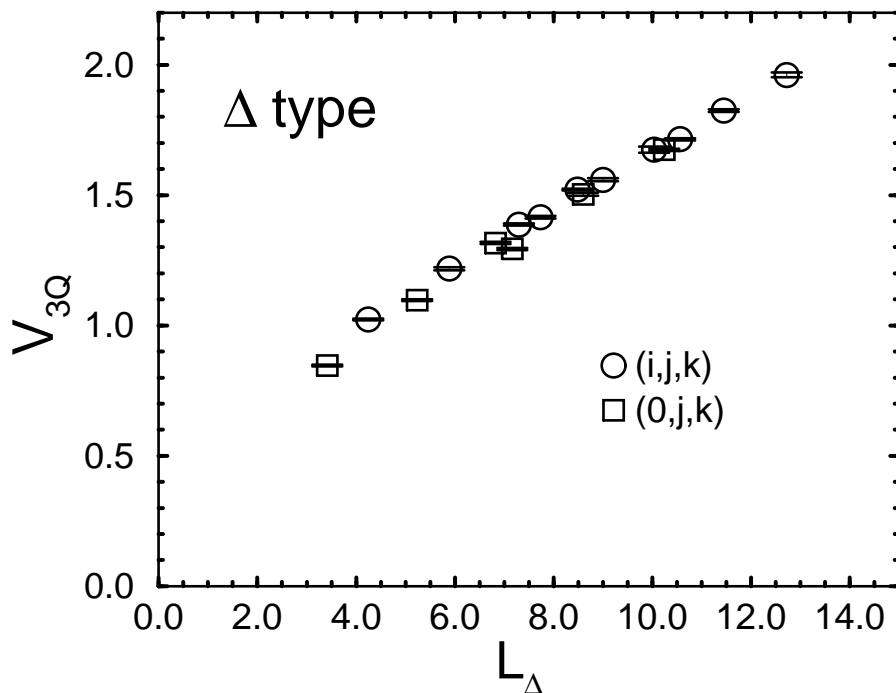
Fit range	const.	A	σ	χ^2/N_{dof}
all r	0.9140(20)	0.1316(62)	0.1528(20)	51.9/13
c.f. $Q\text{-}\bar{Q}$ ($r \geq 2$)	0.696(24)	0.395(35)	0.1561(38)	60.0/18

Lattice cutoff: $a^{-1} \simeq 1$. GeV

Well described by Y-type form

$\sigma_{3Q} \simeq \sigma_{Q\bar{Q}}$: universality of string tension holds.

Lattice Simulation (4): Δ type ansatz



Smearing: 12 sweeps

Fit:

$$V(\vec{r}_1, \vec{r}_2, \vec{r}_3) = \text{const.} - A \left(\frac{1}{r_{12}} + \frac{1}{r_{23}} + \frac{1}{r_{31}} \right) + \sigma(r_{12} + r_{23} + r_{31})$$

Fit range	const.	A	σ	χ^2/N_{dof}
all r	0.934(20)	0.1405(60)	0.0858(15)	142./13
c.f. $Q-\bar{Q}$ ($r \geq 2$)	0.696(24)	0.395(35)	0.1561(38)	60.0/18

Lattice cutoff: $a^{-1} \simeq 1$. GeV

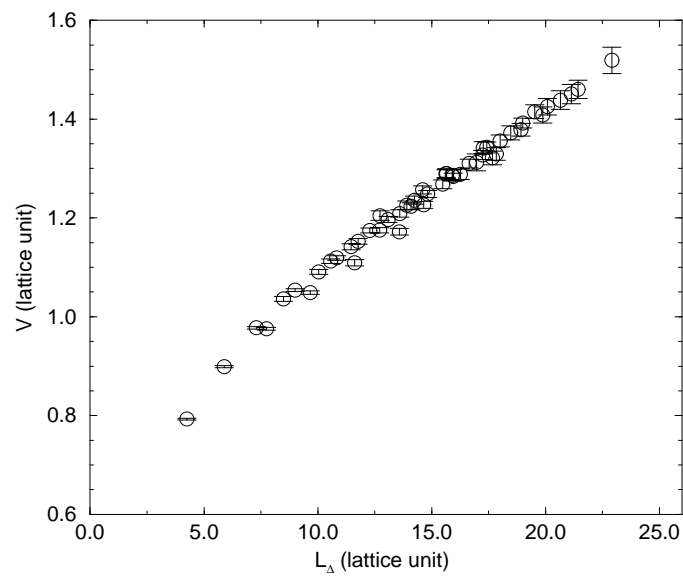
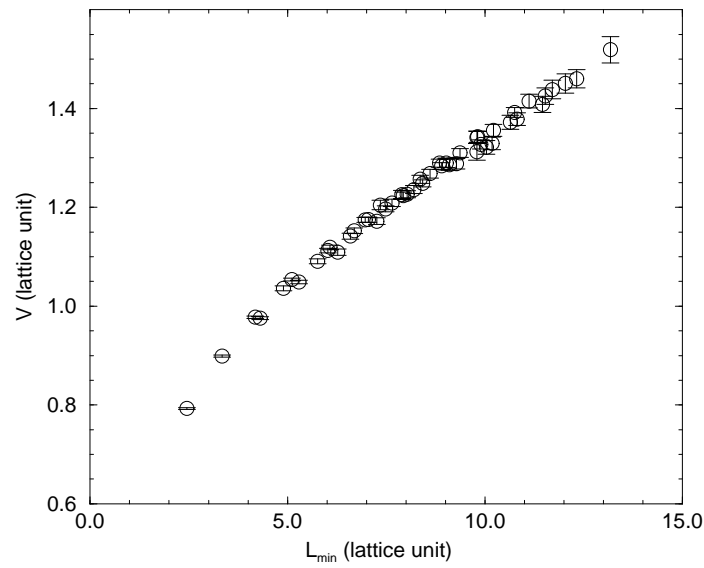
Fit seems not so bad.

$\sigma_{\Delta} \simeq \frac{1}{2}\sigma_{Q\bar{Q}}$ — No clear explanation.

(color factor like as Coulomb term ?)

○ χ^2 is much larger than Y-type fit.

Lattice Simulation (5): Preliminary result at $\beta = 6.0$



Smearing: 42 sweeps ($3Q$) / 36 sweeps ($Q\bar{Q}$)

Fit	const.	A	σ	χ^2/N_{dof}
Y	0.933(12)	0.1203(46)	0.0475(14)	54.1/46
Δ	0.932(12)	0.1195(45)	0.02685(75)	25.9/46
$Q\bar{Q}$ (on-axis)	0.6247(63)	0.2674(50)	0.0532(16)	0.51/3

Lattice cutoff: $a^{-1} \simeq 2$. GeV

Lattice Simulation (6): Summary and Outlook

We calculated the three quark potential with the baryonic Wilson loop, making use of smearing technique.

- Which of Y - and Δ -type is appropriate description ?

□ Fit result at $\beta = 5.7$:

$$\chi^2 (Y\text{-type})/N_{dof} = 51.9/13 = 3.99$$

$$\chi^2 (\Delta\text{-type})/N_{dof} = 142./13 = 10.9$$

$$\chi^2 (Q\bar{Q}; r > 2)/N_{dof} = 60.0/18 = 3.33 \quad (\rightarrow \text{guide of finite } a \text{ effect})$$

Y -type is preferable, but Δ -type is not strongly excluded.

□ String tension: $2\sigma(\Delta) \simeq \sigma(Y) \simeq \sigma_{Q\bar{Q}}$

– In the fit, this is because geometrically $l_{min} \simeq \frac{1}{2}l_{\Delta}$ roughly holds in most cases. ($\frac{1}{2}l_{\Delta} \leq l_{min} \leq \frac{1}{\sqrt{3}}l_{\Delta}$)

String: nonperturbative picture

— Physical reason of $\frac{1}{2}$ is unclear.

Our conclusion: Y -type ansatz is better for numerical result, although Δ -type ansatz is not completely excluded yet.

\Rightarrow Simulation at larger β ($\beta = 6.0$ is in progress)

Outlook

- Precise results / extrapolation to the continuum limit
- 4 quark system ($QQ\bar{Q}\bar{Q}$)
 - Difference of Y -type and Δ -type ansatz are more clear.
- Excited states of string \leftarrow Variational analysis
- With dynamical quarks – String breaking effect

At Finite Temperature (1): Polyakov Loop Correlators

Finite temperature lattice QCD

We use anisotropic lattice:

$$a_\sigma > a_\tau \quad (\xi \equiv a_\sigma/a_\tau: \text{anisotropy})$$
$$N_\sigma a_\sigma \gg N_\tau a_\tau = 1/T$$

Polyakov loop

$$P(\vec{x}) = \text{Tr} \prod_{t=0}^{N_\tau-1} U_4(\vec{x})$$

– $\langle P \rangle$ is the order parameter of deconfining transition

Static quark potentials

◦ Q - \bar{Q} potential:

$$P_2(\vec{r}) = \langle P(0)P^\dagger(\vec{r}) \rangle \simeq c \cdot \exp(-V_{Q\bar{Q}}(\vec{r})N_\tau)$$

◦ $3Q$ potential:

$$P_3(\vec{r}_1, \vec{r}_2, \vec{r}_3) = \langle P(\vec{r}_1)P(\vec{r}_2)P(\vec{r}_3) \rangle$$
$$\simeq c' \cdot \exp(-V_{3Q}(\vec{r}_1, \vec{r}_2, \vec{r}_3)N_\tau)$$

At Finite Temperature (2): Simulation

□ Lattice

Gauge field action: Anisotropic Symanzik (tree) action

$$16^2 \times 24 \times N_t, \beta = 4.56, \gamma = 3.45, \text{ quenched}$$

- Cutoff: $a_\sigma^{-1} = 1.61(1) \text{ GeV}$ ($a_\sigma \sim 0.125 \text{ fm}$),
 $a_\tau^{-1} = 6.36(5) \text{ GeV}$

- Anisotropy: $\xi \equiv a_s/a_t = 3.95(2)$

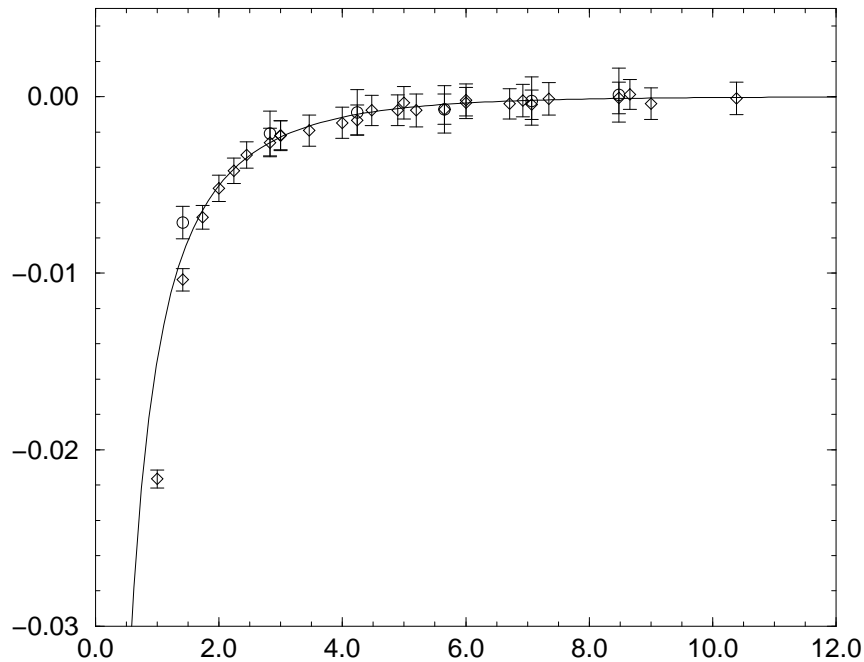
(from the ratio of Wilson loops)

[Engels, Karsch and Scheideler (1997), Klassen (1998)]

$$N_t = 96 (T \simeq 0), 28, 26 (T < T_c), 20 (T > T_c) : T = 1/N_t a_t$$

N_t	T	#conf	observable	$\sqrt{\sigma_{Q\bar{Q}}} [\text{MeV}]$
96	$\simeq 0$	120	Wilson loop	420
28	$0.87T_c$	440	Polyakov loop	~ 300
26	$0.93T_c$	450	Polyakov loop	~ 260
20	$1.22T_c$	60	Polyakov loop	-

At Finite Temperature (3): Above T_c



- $Q\text{-}\bar{Q}$ system — Debye screening

$$V_{Q\bar{Q}} = V_{singlet} + V_{octet}$$

$$\rightarrow c \cdot \exp(-mr)/r \quad : \text{ Yukawa potential}$$

[Gao, PRD41(1990)626]

$m \sim$ screening mass $\sim 0.5\text{-}1$ GeV

- $3Q$ system:

— Sum of two-quark ($Q\text{-}Q$) potential ?

Well fitted to the form

$$V_{3Q} = \sum c \frac{\exp(-mr_{ij})}{r_{ij}}.$$

$V_{3Q}/3$ vs r_{ij} for equilateral triangle configuration

→ Almost same behavior as $Q\text{-}\bar{Q}$ system.

At Finite Temperature (4): Summary

□ At $T > T_c$

We calculated the static three quark potential above T_c with equilateral constellations.

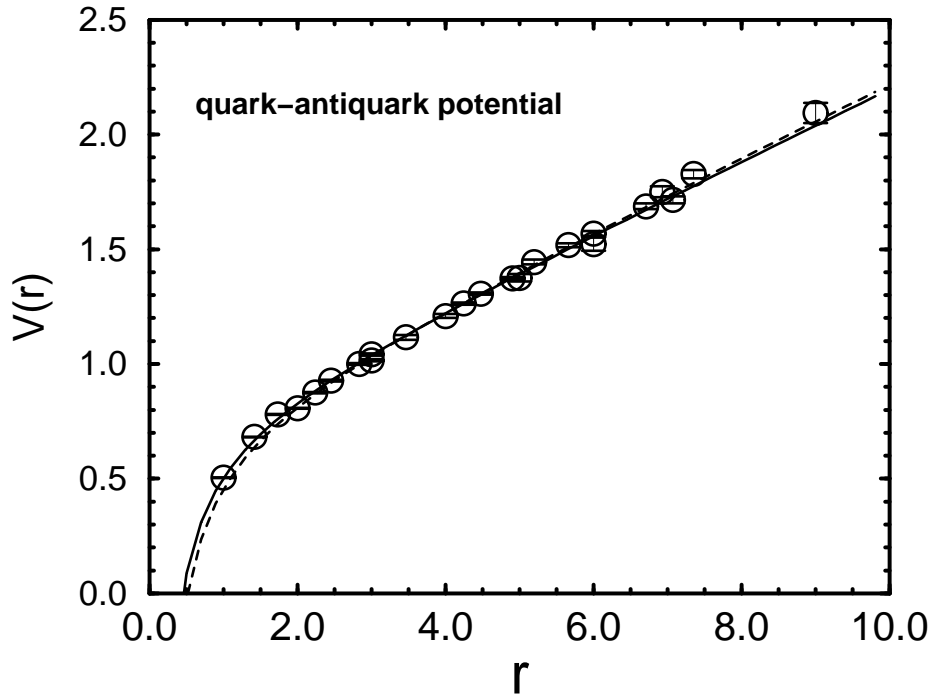
- Debye screening observed.
- $V_{3Q}/3$ shows almost same behavior as $Q\bar{Q}$ potential.

□ At $T < T_c$

$3Q$ potential from Polyakov loop correlator suffer from large statistical fluctuations.

— More study is needed.

Appendix (1): $Q\text{-}\bar{Q}$ potential ($T = 0$)



Smearing: 20 sweeps

Fit:

$$V(r) = \text{const.} - \frac{A}{r} + \sigma r$$

Fit range	const.	A	σ	χ^2/N_{dof}
on-axis	0.629 (16)	0.2793(12)	0.1629(47)	1.77/ 3
all r	0.6859(71)	0.3351(55)	0.1545(20)	341./ 21
$r \geq 2$	0.696 (24)	0.395(35)	0.1561(38)	60.0/18 ~ 3

Lattice cutoff: $a^{-1} \simeq 1$. GeV

- Large χ^2 signals large $O(a^2)$ discretization effect
(Rotationally symmetry is broken)