

# Finite temperature hadron physics on the lattice

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## The QCD-TARO Collaboration

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### Physical goals:

- Finite Temperature QCD
- Finite Chemical Potential
- MCRG and Improved Actions

## Introduction

Finite temperature hadron properties:

- What happens on hadrons  
or what new effect above  $T_c$  ?
- Change of masses and width near  $T_c$

*Example of effective theory:*

NJL model analysis

*c.f. Hatsuda and Kunihiro, Phys. Rep. 247 (1994) 221*

- Chiral restoration
- parity partner degeneracy
- Above  $T_c$  : Soft modes

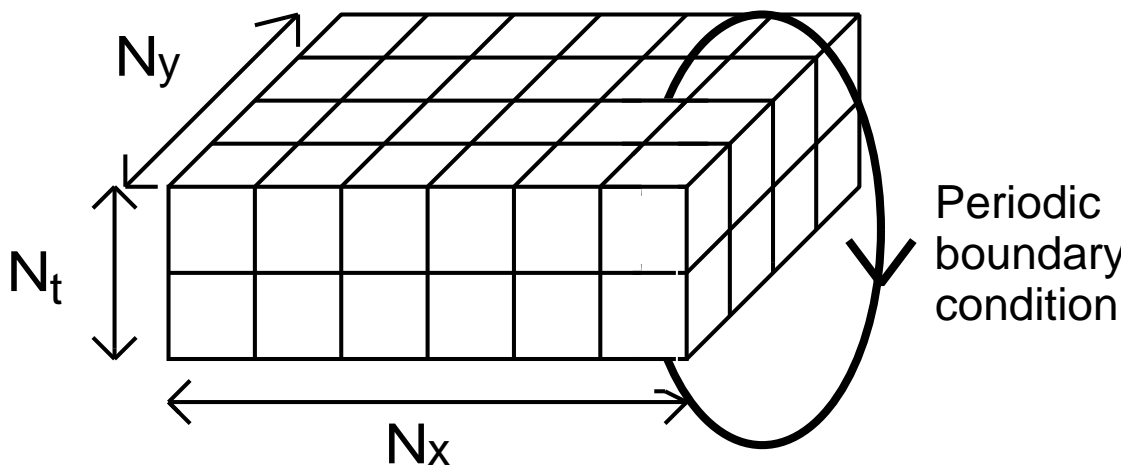
## Lattice QCD

Discretization of space-time ( Euclidean space )  
with **exact local gauge invariance**



Path integral formalism has a well-defined meaning  
⇒ Monte Carlo Method

## Finite Temperature lattice QCD



$$T = \frac{1}{N_t a} \quad ( a : \text{lattice spacing} )$$

## Our Approach

Temporal correlator ;

⇒ Pole mass (mass of temporal direction)



Screening mass (mass of spatial direction)

⇒ Wave function

Is there bound state at  $T > T_c$  ?

How can we show that ?

⇒ Spectral function

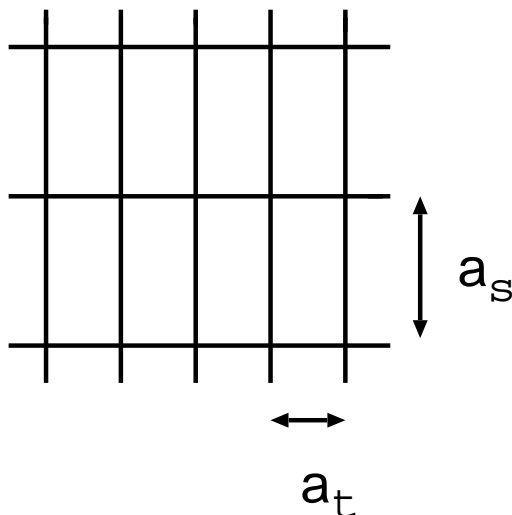
Need detailed information

in temporal(temperature) direction

→ Anisotropic Lattice

Karsch ('82)

Burgers, Karsch, Nakamura and Stamatescu ('88)



anisotropy  $\xi$

$$\xi = a_s / a_t$$

$$\xi_F = m_s / m_t \quad (= \xi)$$

↑ calibration

## (I) Meson with light quark

- Temperature dependence of mass  
( Pole mass vs Screening mass)
- Hadronic bound state above  $T_c$  (?)
- Spectral function

### Simulation parameters

*Gauge field:*

$\beta = 5.68, \gamma = 4.0$ , quenched, 60conf.

- lattice size :  $12^3 \times N_t$

$$N_t = 72 \ ( T \simeq 0 ) \quad : \ T = 1/N_t a_t$$

$$20 \ ( T \simeq 0.93T_c )$$

$$16 \ ( T \simeq 1.15T_c )$$

$$12 \ ( T \simeq 1.5T_c )$$

- Anisotropy :  $\xi \equiv a_s/a_t = 5.3(1)$

*Engels, Karsch and Scheideler (1997), Klassen (1998)*

- Cutoff :  $a_s^{-1} = 0.85(3) \text{ GeV}, a_t^{-1} = 4.5(2) \text{ GeV}$   
from heavy quark potential

*Quark field* : Anisotropic Wilson action

$$m_q \sim 0.17 \text{ GeV} (m_{P_s} \sim 0.81 \text{ GeV})$$

$$0.12 \text{ GeV} ( \quad \quad 0.68 \text{ GeV})$$

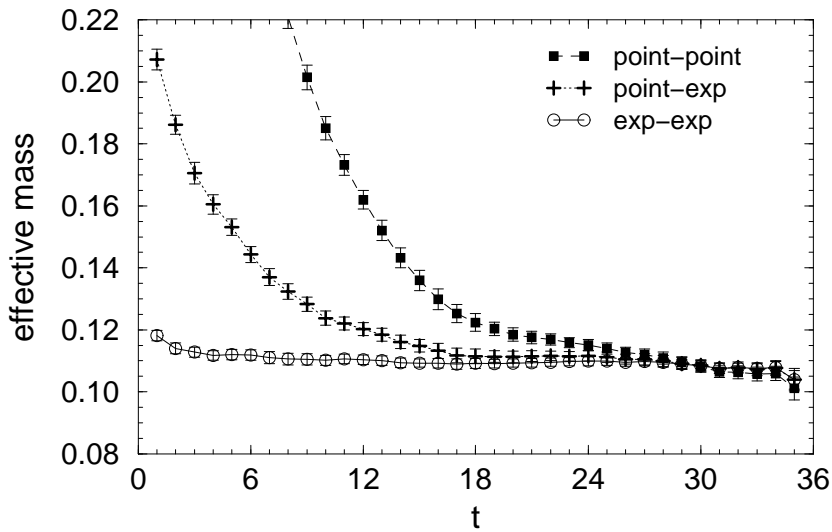
$$0.10 \text{ GeV} ( \quad \quad 0.61 \text{ GeV})$$

## Effective mass

Effective mass :  $m_{\text{eff}}$

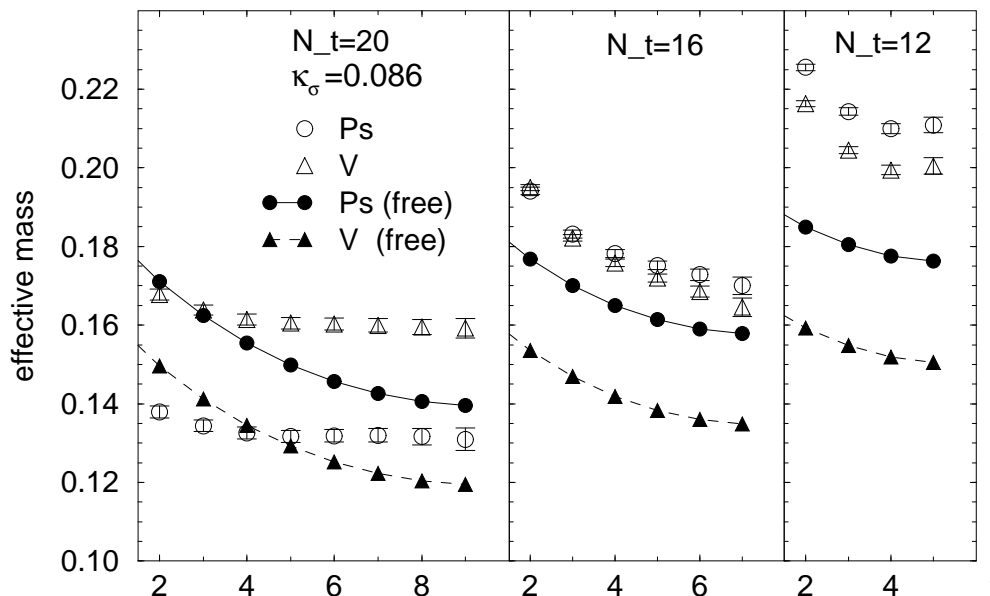
$$\frac{G_M(\vec{x} = 0, t)}{G_M(\vec{x} = 0, t + 1)} = \frac{\cosh [m_{\text{eff}}(N_t/2 - t)]}{\cosh [m_{\text{eff}}(N_t/2 - t - 1)]}$$

○ Source dependence

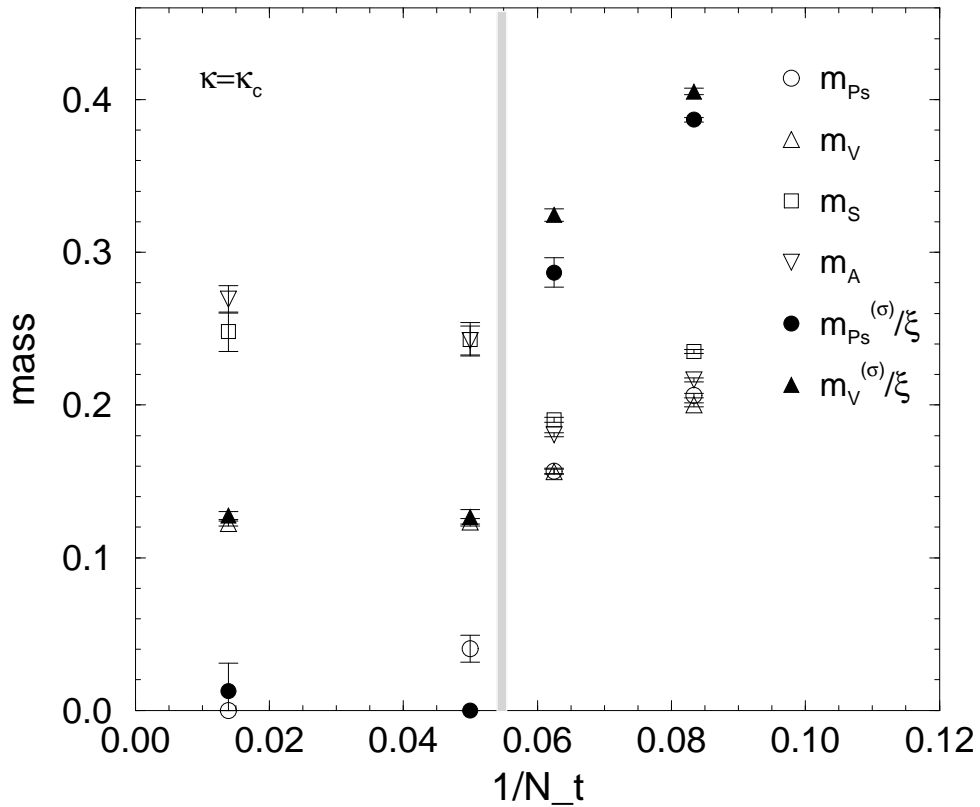
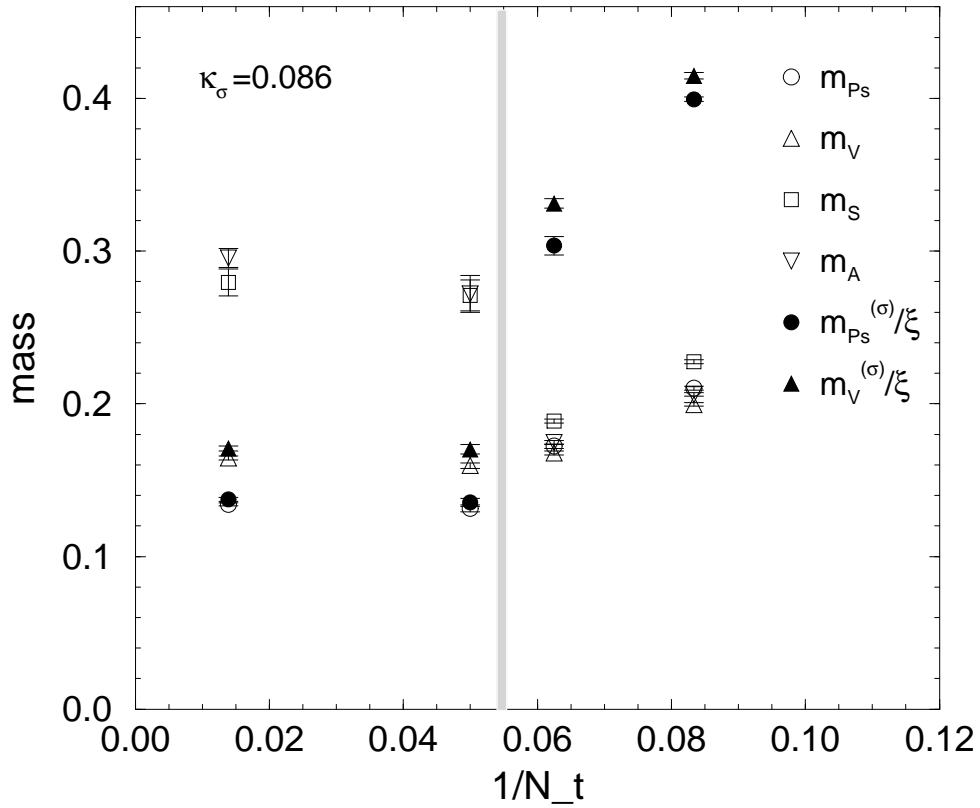


⇒ We adopt the **exp-exp source**

○ effective mass at  $T > 0$



## Temperature dependence of mass

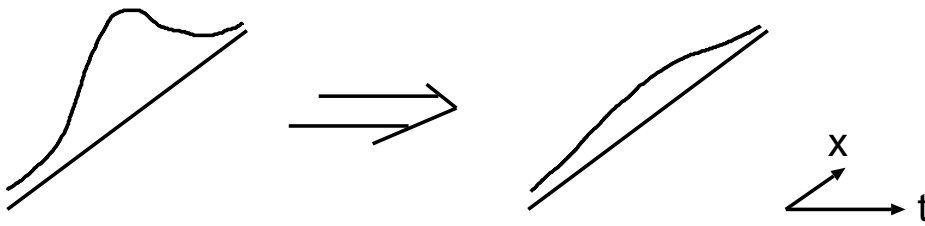




## t-dependence of the wave function

$$w_\Gamma(r, t) = \sum_{\vec{x}} \langle \bar{q}(\vec{x} + \vec{r}, t) \Gamma q(\vec{x}, t) O^\dagger(0) \rangle$$

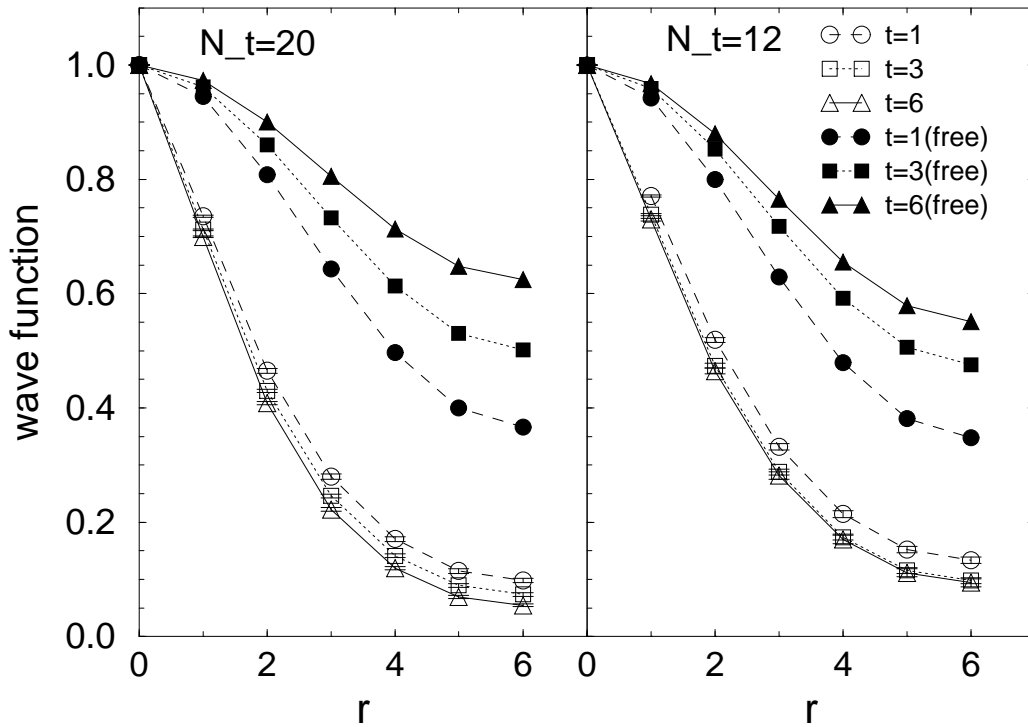
If there is no bound state (like free quark case), wave function becomes broader as  $t$ .



In this case,  $\phi_\Gamma(r, t) = w_\Gamma(r, t) / w_\Gamma(r = 0, t)$   
(normalized at spatial origin)

increase as  $t$ .

Ps, exp-exp source

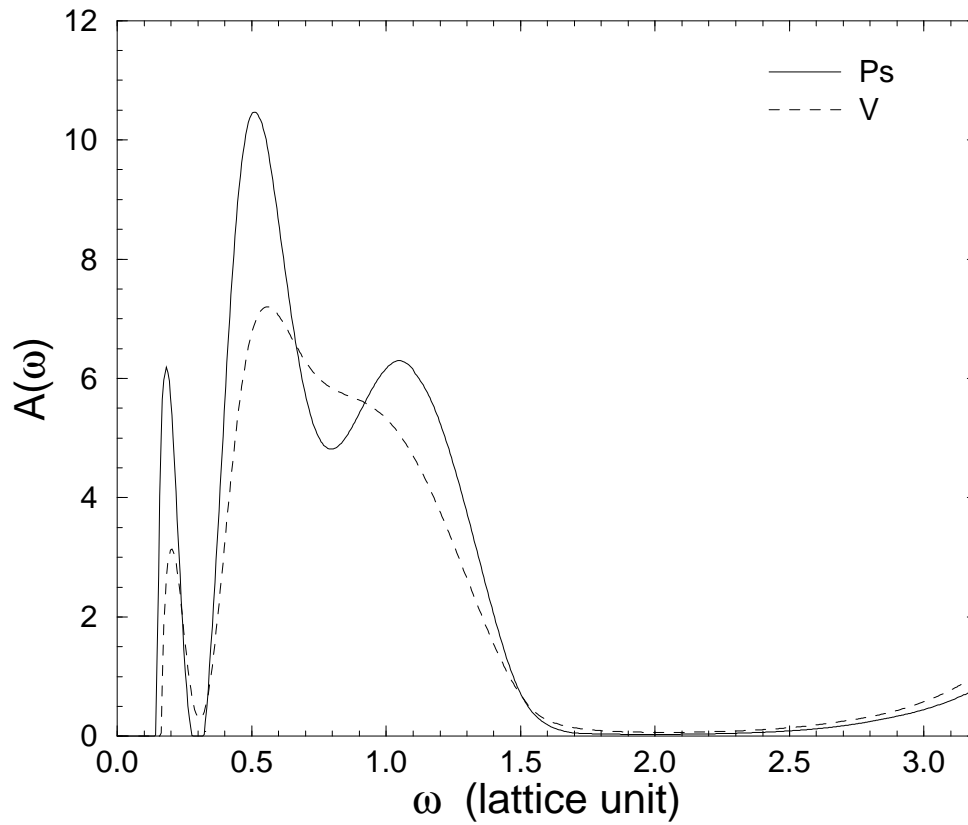


## Spectral function

The two point correlation function of meson :  $C(t)$

$$C(t) = \int K(t; \omega) A(\omega) d\omega$$
$$K(t; \omega) = e^{-\omega t} + e^{-\omega(N_t - t)}$$

$A(\omega)$  : spectral function



## Summary and Conclusion (I)

- Effective mass
  - Significant change above  $T_c$
- Wave functions
  - Narrower than free  $q - \bar{q}$  even at  $N_t = 16, 12$ .  
Still prefer to stay together.
- Spatial-mass (screening mass)  
vs temporal-mass (pole mass)
  - Grows above  $T_c$ .
  - Stronger growing of spatial-masses.
  - Consistent with NJL model analysis.

## (II) Meson with heavy quark

(Problem)

Spatial lattice spacing is coarse

$$(a_s^{-1} \sim 0.85 \text{ GeV})$$

Charmonium mass  $\sim$  a few GeV



- Heavy quark action : **Fermilab action**  
(  $O(a)$  improved ( Clover action ) )

$\implies$  Heavy quarkonium physics

Clover action : Sheikholeslami, Wohlert ('85)

Fermilab action : El-Khadra, Kronfeld, Mackenzie ('97)

We study the Charmonium at  $T \geq 0$

- mass shift near  $T_c$   
T.Hashimoto et al, ('86)
- $J/\psi$  suppression above  $T_c$   
T.Matsui and H.Satz, ('86)

## Simulation parameters

On the same configuration as (I) ( see p.5 )

Target meson mass is  $J/\psi$  mass :  $m_V \sim 3.1 \text{ GeV}$

## Effective mass at $T > 0$

— Heavy quark field—

- parameter

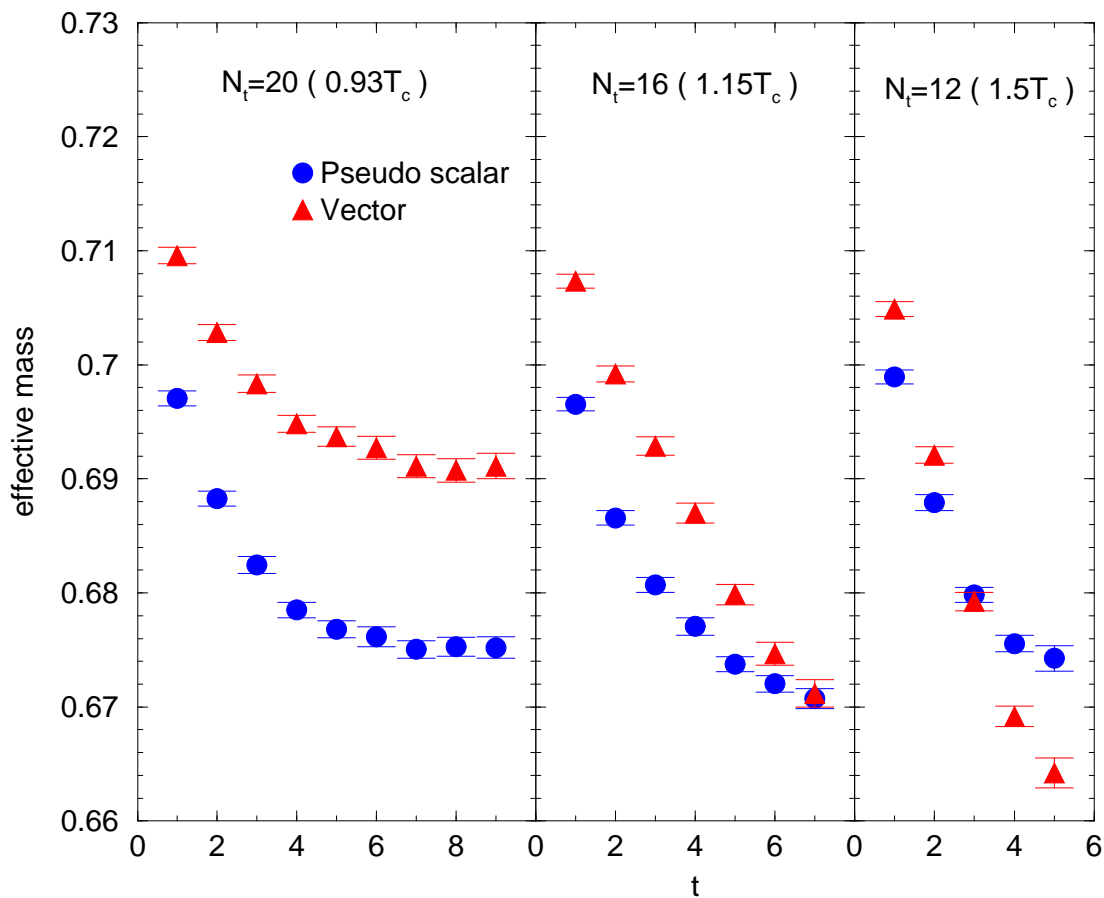
$$1/\kappa = 9.690, \quad \gamma_F = 3.67$$

- results at  $T = 0$

$$m_{P_s} = 3.04(14)\text{GeV}, \quad m_V = 3.10(14)\text{GeV}$$

- Finite Temperature

### effective mass at $T > 0$



## Summary and Conclusion (II)

- (1) We use the anisotropic Fermilab action for heavy quark system.
- (2) Meson correlator
  - No significant change at  $N_t = 20$
  - Significant change at  $N_t = 16, 12$

The Temperature dependence is different from that of light mesons .  
Especially behavior of vector meson change.

## Out look

- Precise determination of mass shift

- variational analysis
- Spectral function
- Fine lattice

⇒ Now under progress

Lattice size :  $16^2 \times 24 \times N_t$

$N_t = 96$  ( at  $T \sim 0$  )

$\beta = 4.56$  ( Symanzik action (tree) )

Anisotropy :  $\xi \sim 4.0$

Cutoff :  $a_s^{-1} \sim 1.6$  GeV (  $a_s \sim 0.12$  fm )

- More detailed analysis of Heavy quarkonium :  $J/\psi, \dots$
- Other channels : baryons etc.
- Relation with topological quantities.
- With dynamical quarks.