

Issues on beam dynamics in SuperKEKB: Hard-edge dipole fringe and FMA

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With contributions from

KEK: E. Forest, K. Ohmi, K. Oide

Cornell: D. Sagan

SLAC: Y. Cai

SuperKEKB mini optics meeting, Nov. 27, 2014

Outline

➤ Hard-edge model for dipole fringe field

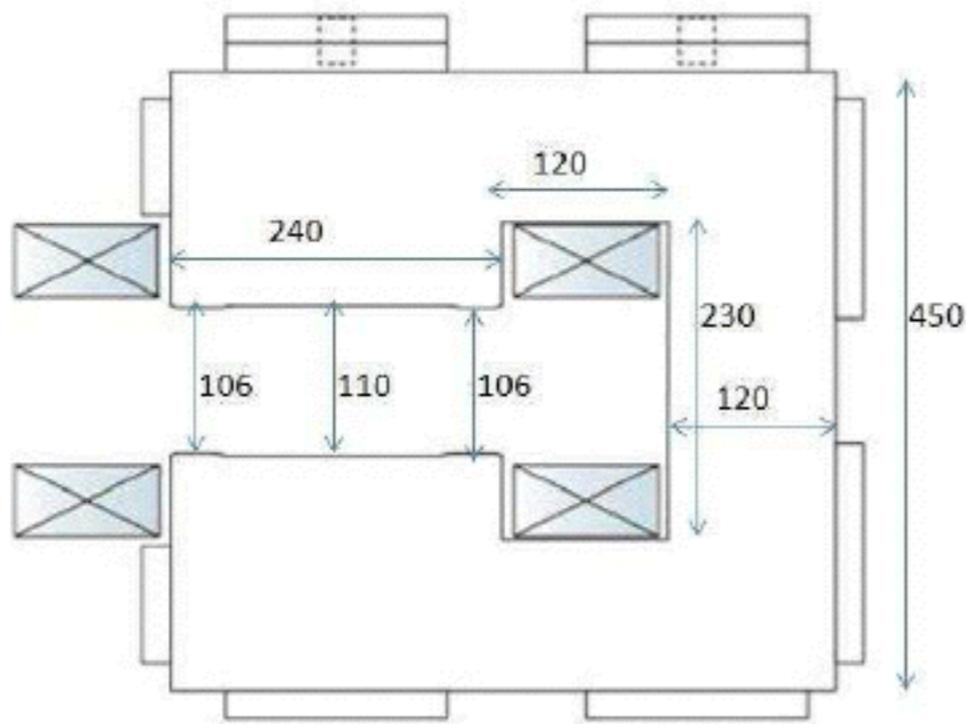
- Resolve the discrepancy of two models: Cai-Nosochkov's and Forest's

➤ Benchmark of SAD and Bmad

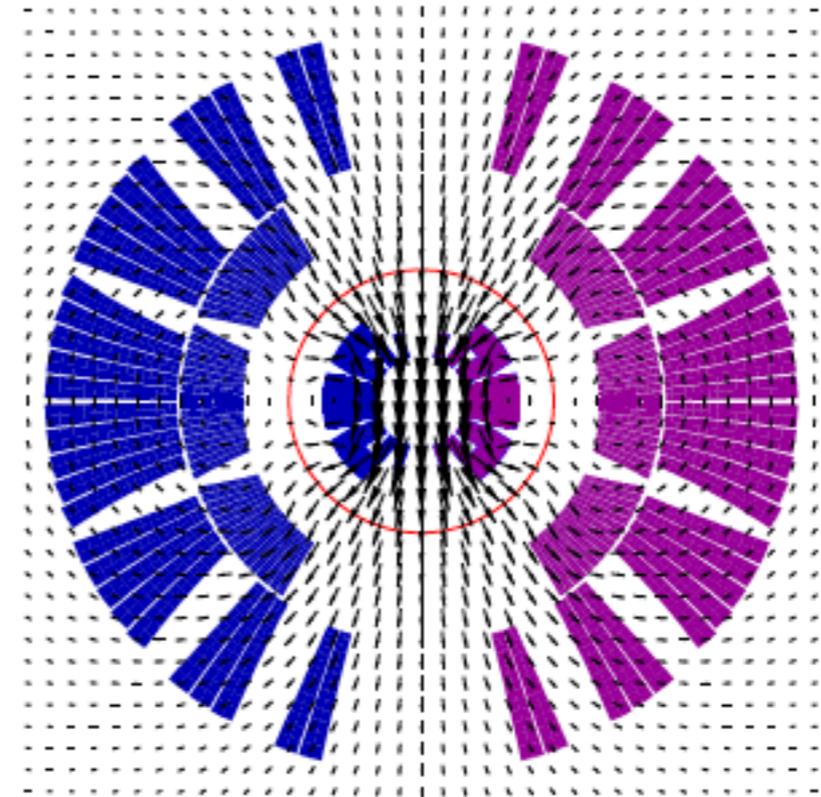
- FMA results

1. Dipole fringe

- Question raised from debugging SAD code
- Two models found for **hard-edge fringe**
 - E. Forest: “Parallel-plate” shape (popular theory)
 - Y. Cai: Round shape (apply for SC magnets?)



Realistic case
(From SuperKEKB TDR)

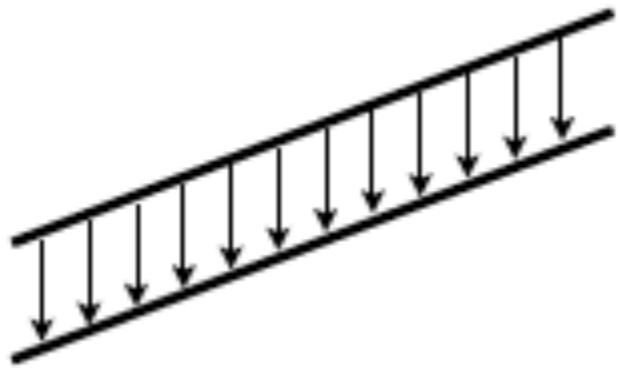


SC magnet
(From S. Russenschuck's
textbook, 2010)

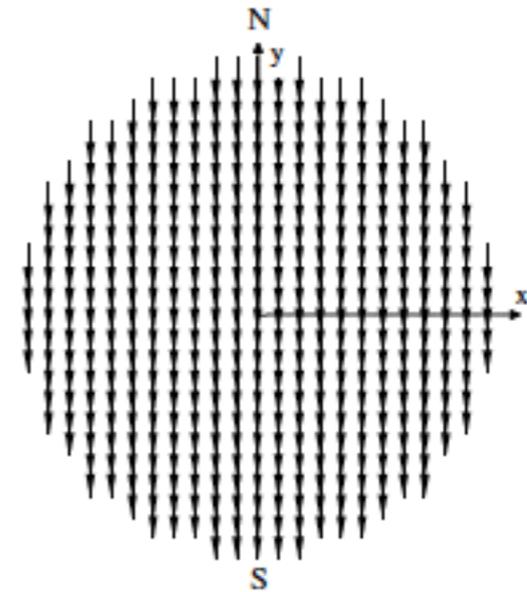
1. Dipole fringe

➤ Maxwellian solution for hard-edge dipole field

- G. Lee-Whiting et al. => E. Forest et al.
- S. Caspi et al. => M. Bassetti et al. => Y. Cai et al.



The model for wide magnet.
The field is confined at region of
 $-b < y < b$ and $-\infty < x < \infty$



The model for harmonics expansion.
The field is confined inside a circle
with $r < r_0$
(From S. Russenschuck's textbook,
2010)

1. Dipole fringe

➤ Maxwellian solution for hard-edge dipole field

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$$A_s = -xB(s) = -xB_0\theta(s).$$

$$\vec{A} = (A_x, 0, A_s)$$

$$\nabla \times \nabla \times \vec{A} = 0$$

$$A_x = B_0 \sum_{n=1}^{\infty} \frac{(-1)^n \theta^{(2n-1)}(s)}{(2n)!} y^{2n}$$

$$A_y = 0$$

$$A_x = \frac{1}{2}(x^2 - y^2) \sum_{p=0}^{\infty} \frac{1}{2+p} G_{1,2p+1}(s)(x^2 + y^2)^p,$$

$$A_y = xy \sum_{p=0}^{\infty} \frac{1}{2+p} G_{1,2p+1}(s)(x^2 + y^2)^p,$$

$$A_s = -x \sum_{p=0}^{\infty} G_{1,2p}(s)(x^2 + y^2)^p.$$

$$G_{n,2p}(s) = (-1)^p \frac{n!}{4^p(n+p)!p!} \frac{d^{2p} G_{n,0}(s)}{ds^{2p}},$$

$$G_{n,2p+1}(s) = \frac{dG_{n,2p}(s)}{ds},$$

$$A_y \neq 0$$

1. Dipole fringe

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Field distribution with hard-edge:

$$B_x = 0,$$

$$B_y(y, s) = B_0 \sum_{n=1}^{\infty} \frac{(-1)^n y^{2n} \theta^{(2n)}(s)}{(2n)!},$$

$$B_s(y, s) = -2B_0 \sum_{n=1}^{\infty} \frac{(-1)^n n y^{2n-1} \theta^{(2n-1)}(s)}{(2n)!}.$$

Field distribution with hard-edge:

$$B_x(x, y, s) = -\frac{1}{2} B_0 x y \sum_{p=0}^{\infty} \frac{(-1)^p (x^2 + y^2)^p \theta^{(2p+2)}(s)}{4^p p! (p+2)!}, \quad (20)$$

$$B_y(x, y, s) = B_0 \theta(s) + B_0 \sum_{p=0}^{\infty} \frac{(-1)^{p+1} (x^2 + y^2)^p \theta^{(2p+2)}(s)}{4^{p+1} (p+1)! (p+2)!} [x^2 + (2p+3)y^2], \quad (21)$$

$$B_s(x, y, s) = B_0 y \sum_{n=1}^{\infty} \frac{(-1)^n (x^2 + y^2)^n \theta^{(2n+1)}(s)}{4^n n! (n+1)!}. \quad (22)$$

1. Dipole fringe: conclusion

- Cai-Nosochkov's model applies for round shape (**SC magnet** is close to this case)
- Forest's model applies for wide shape (SuperKEKB dipole is close to this case?)
- The realistic magnets are between the two models

2. Benchmark of SAD and Bmad

➤ Compare FMA using Bmad and SAD

- Same FMA subroutine by D. Zhou
- Examples: sler_1684 and sher_5755
- FMA results: tracking on-momentum particles by varying initial transverse coordinates, plot DA in X-Y space and footprint in tune space.

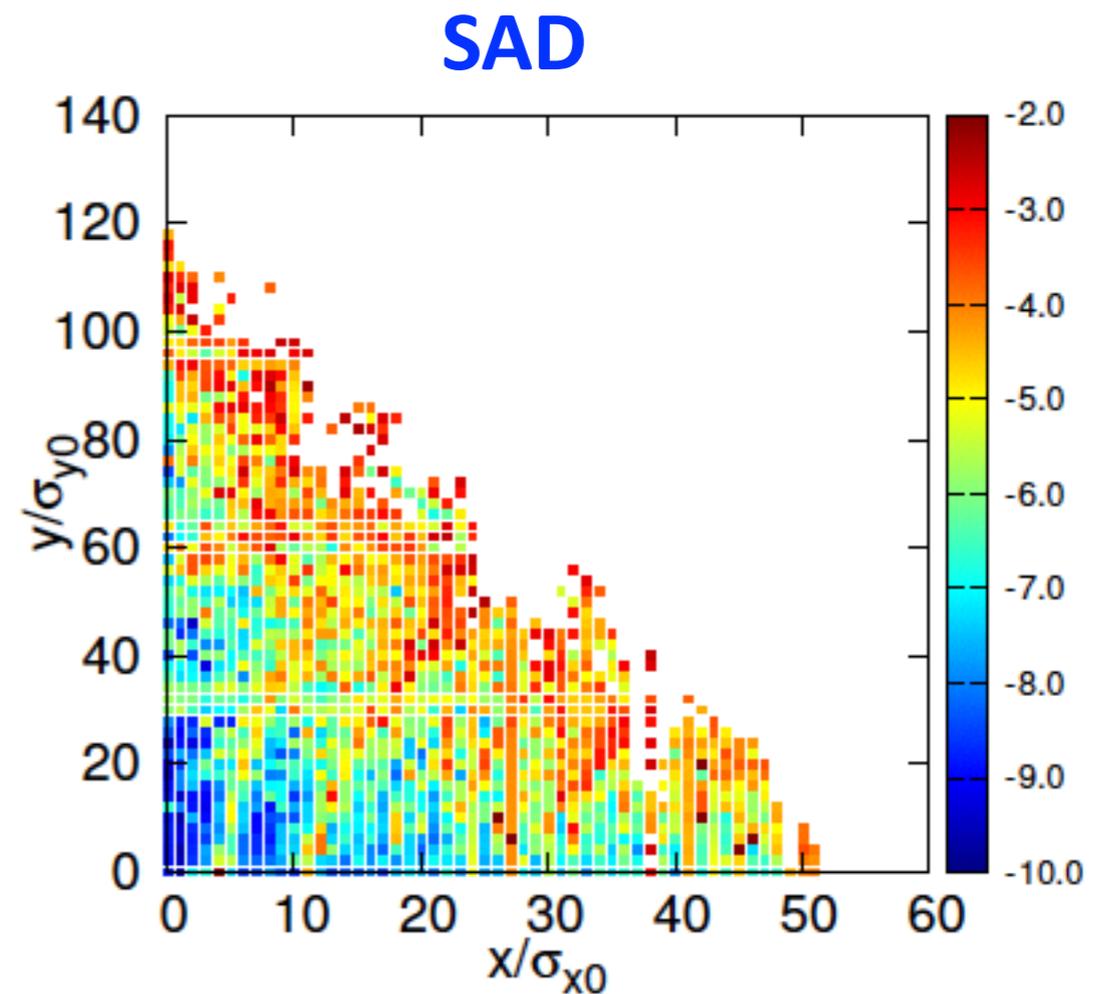
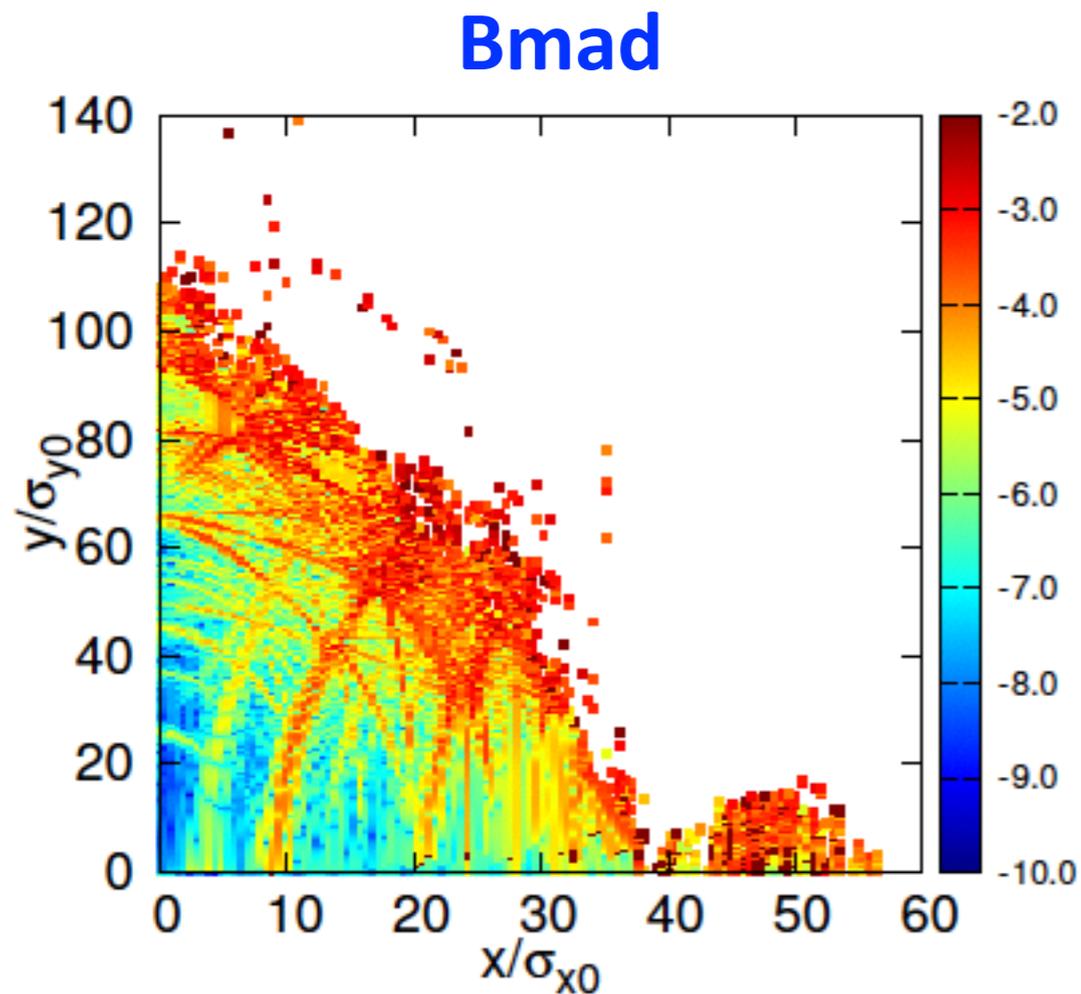
➤ Future work with Bmad and SAD

- Try Bmad's FMA subroutine
- Try E. Forest's PTC on FMA

2. Comparison of FMA: sler_1684

➤ X-Y space

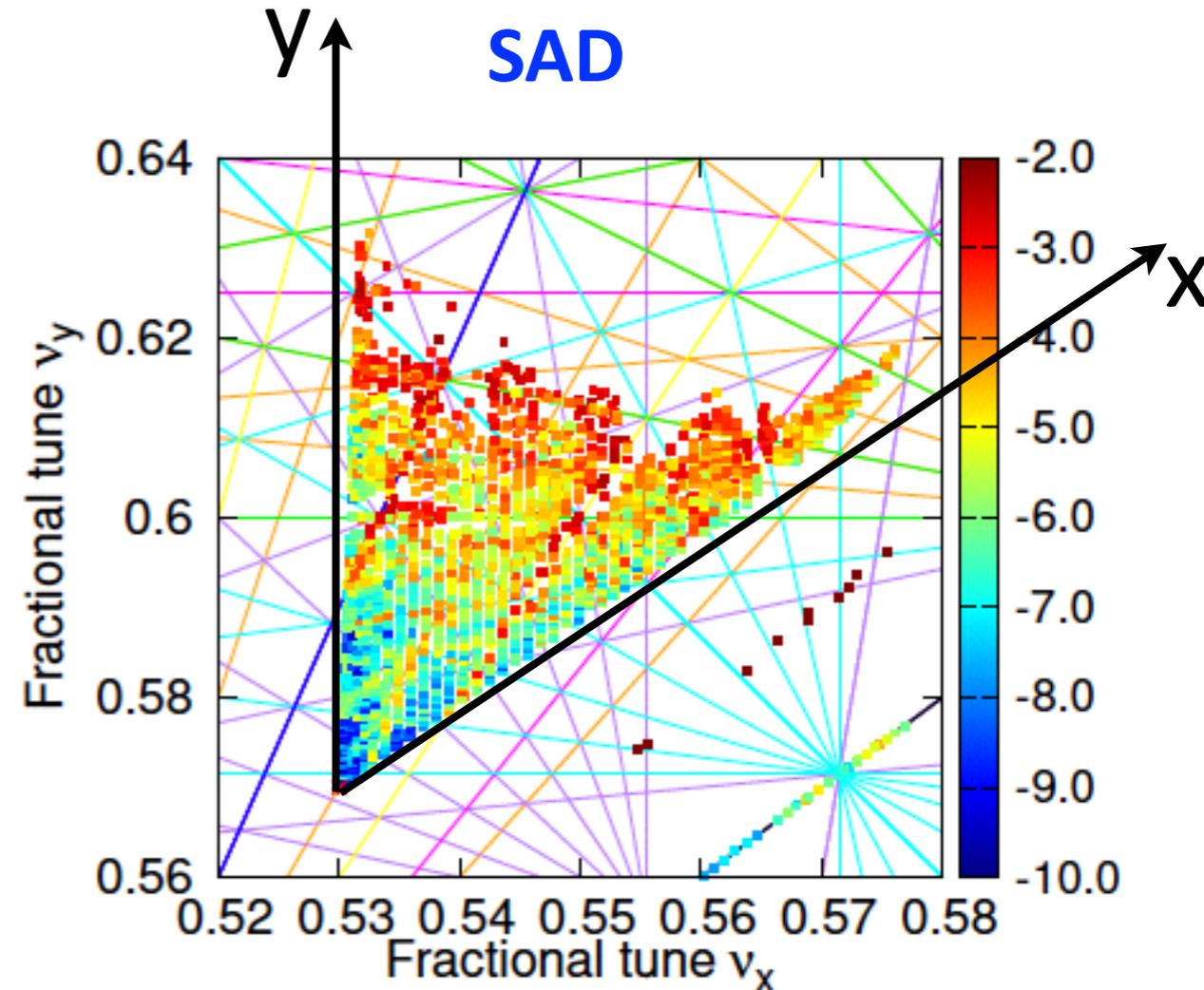
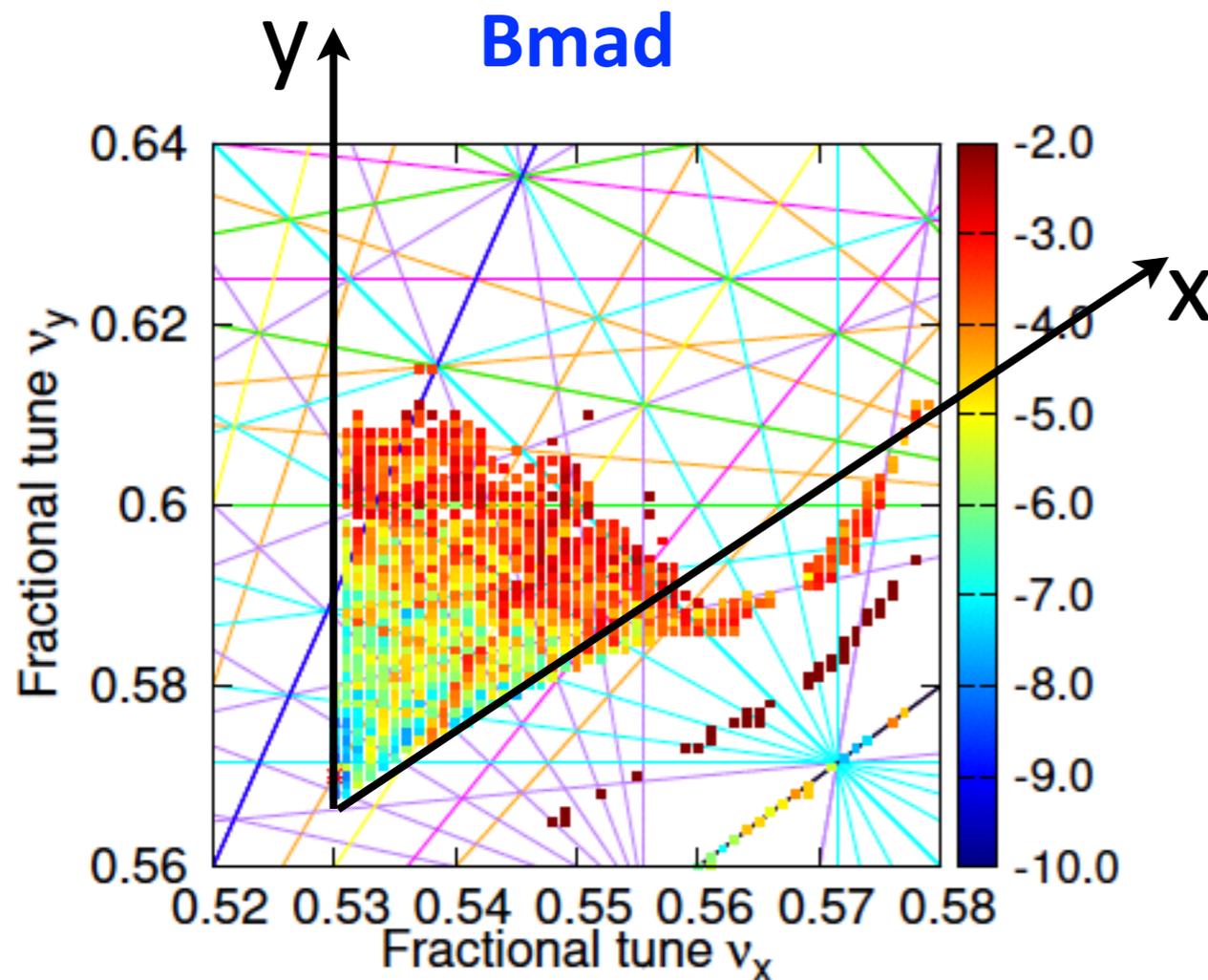
- Bmad and SAD give similar DA in size
- Discrepancy exists in details: indicate different (high-order) nonlinear maps used by Bmad and SAD?



2. Comparison of FMA: sler_1684

► Tune space

- How to evaluate the discrepancy? [David's answer: Check element-by-element tracking data with large amplitude]

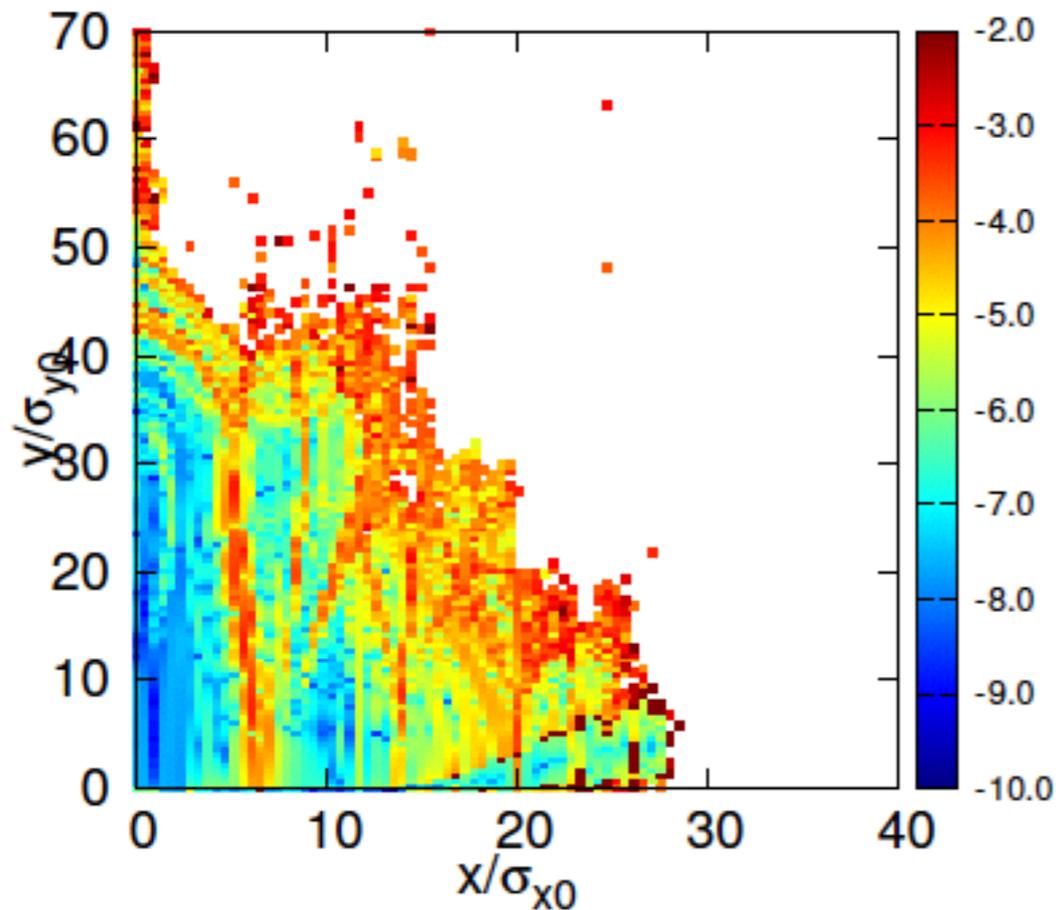


2. Comparison of FMA: sher_5755

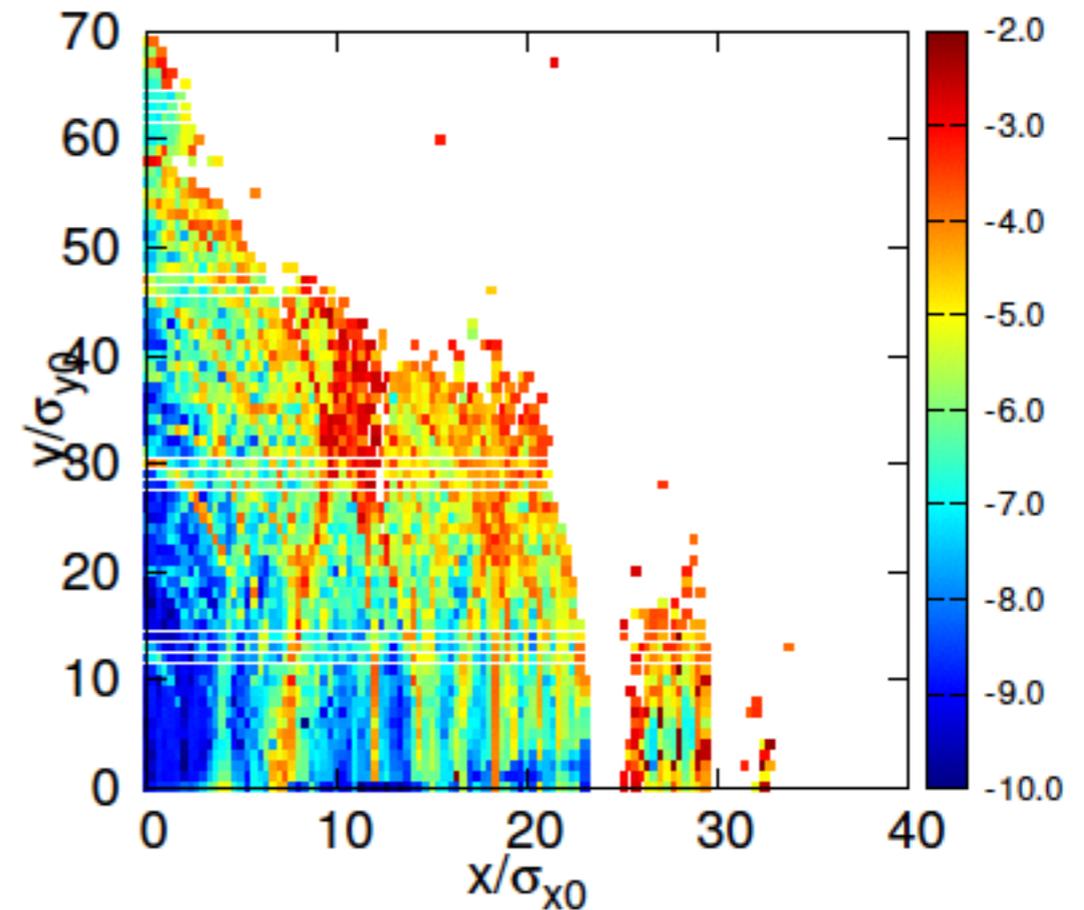
➤ X-Y space

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Bmad



SAD

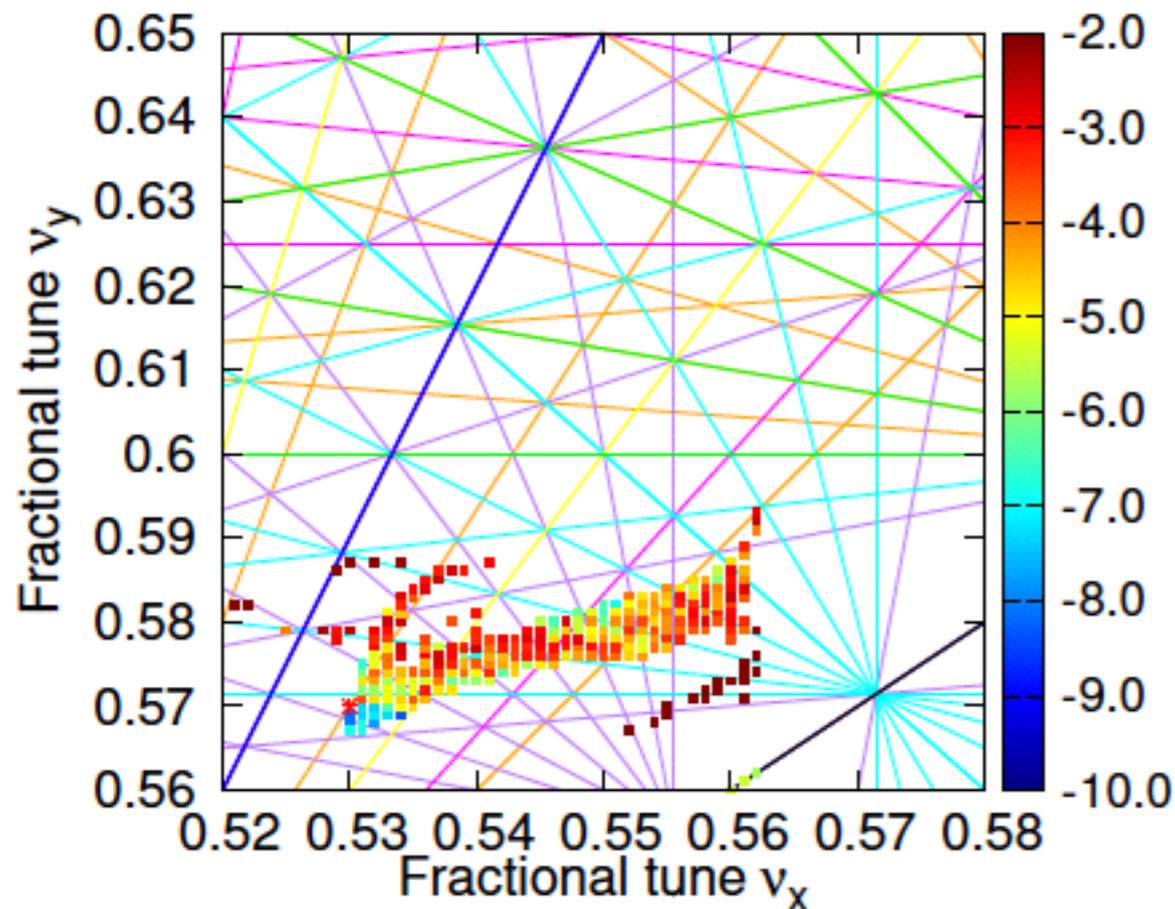


2. Comparison of FMA: sher_5755

► Tune space

- How to evaluate the discrepancy? [David's answer: Check element-by-element tracking data with large amplitude]
- Bmad result looks strange ...

Bmad



SAD

