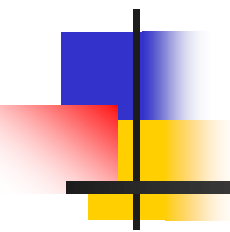
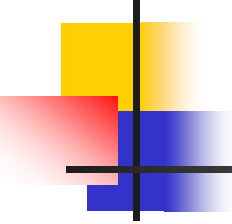


GPDs (and DAs) from (semi-) exclusive DY processes



KEK, March 12, 2015

Oleg Teryaev
JINR, Dubna

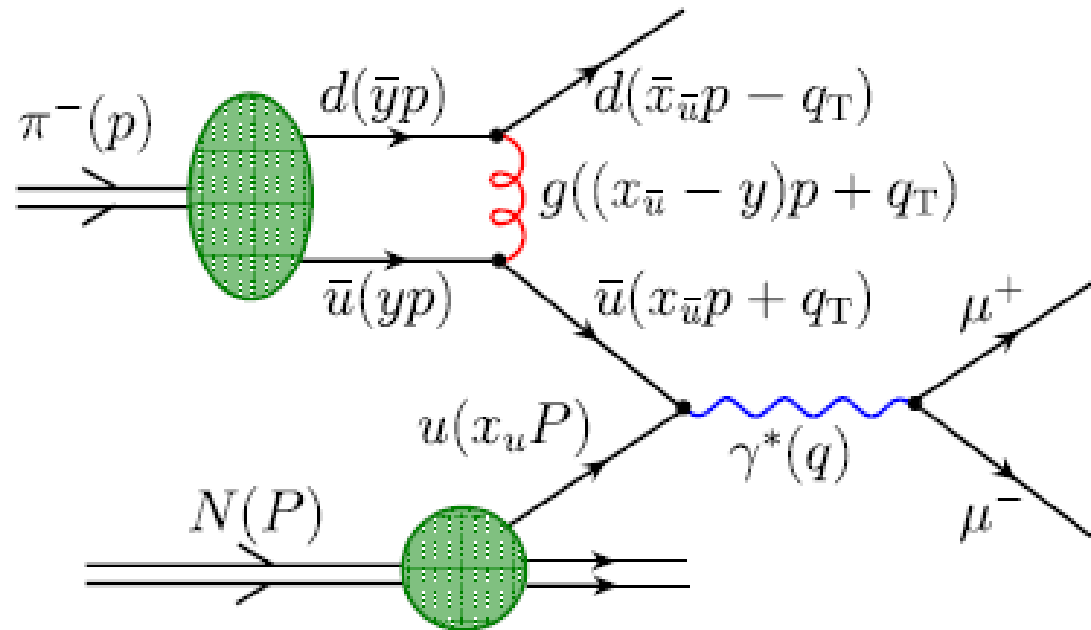


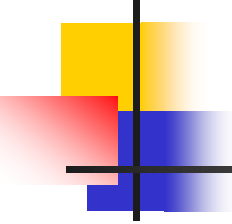
Outline

How GPDs are manifested at large x_F ? Is it possible to separate SE (Brandenburg, Brodsky, Khoze, Muller) from E (Berger, Diehl, Pire)? (thanks to Wen-Chen Chang)

- Semi-exclusive pion-nucleon DY (COMPASS) and pion DA
- Exclusive DY and GPDs
- Role of X_F for E and Q_T for SE
- Role of angular distribution

Semi-Exclusive DY (large x_F) - Pion participates through Distribution Amplitude (Light-cone WF)



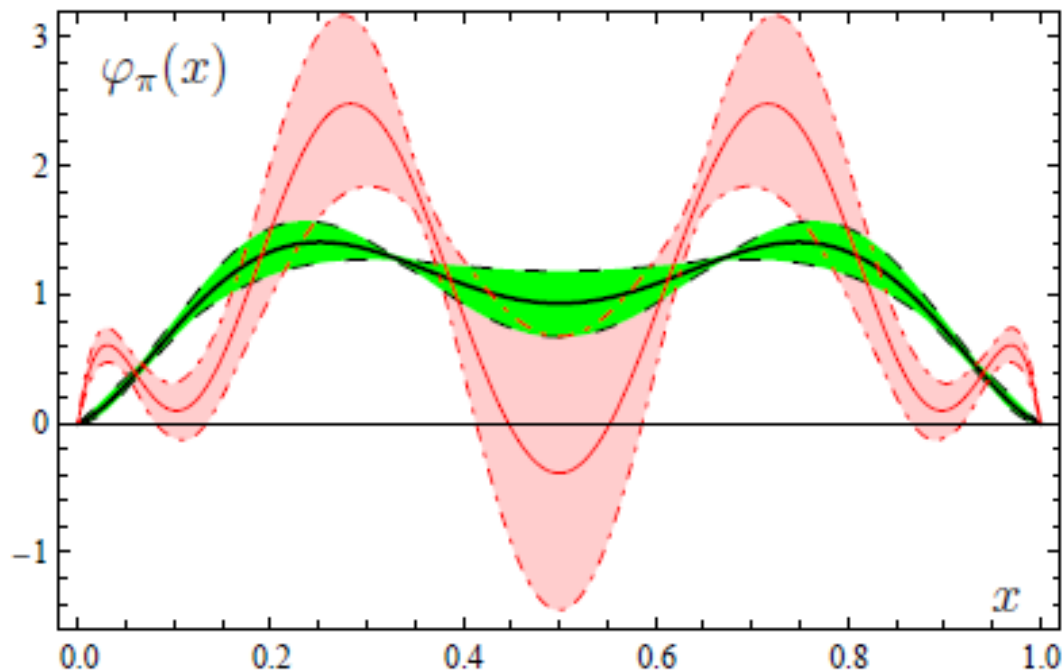


Pion DA

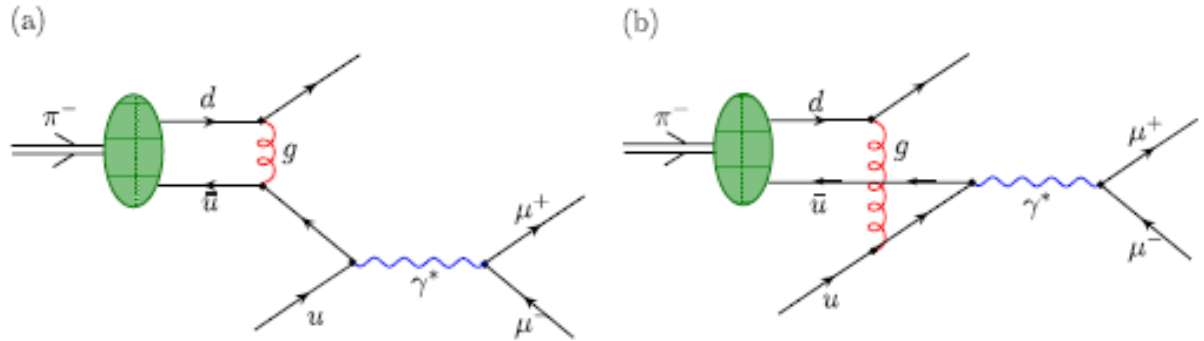
- Element of ERBL factorization
- Describes probability amplitude for the (anti) quark carrying given light-cone momentum fraction
- Interest recently increased due to BaBar/BELLE data for pion-photon transition formfactor-simplest exclusive process

Pion DA

- (Conservative) model of Bakulev, Mikhailov, Stefanis vs (3D) fit



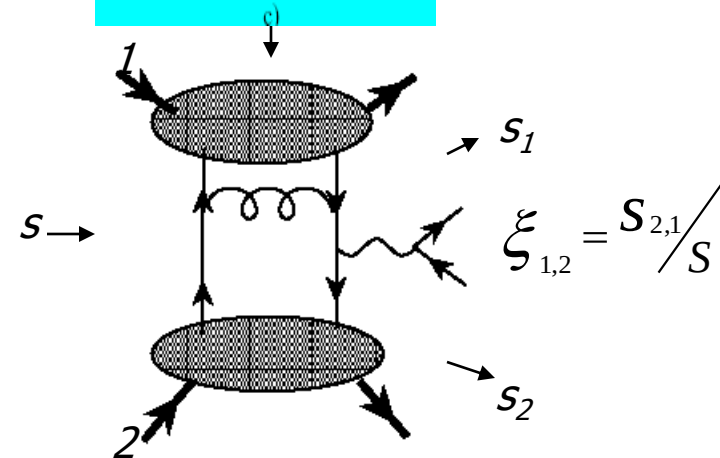
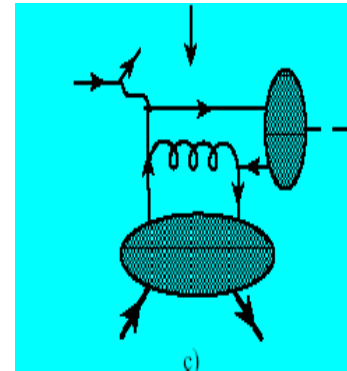
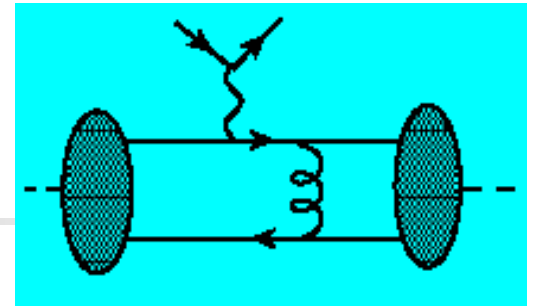
GI ->

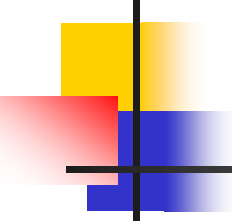


- Colour GI -> second diagram -> phase
- Unpolarized – Brandenburg, Brodsky, Khoze, Mueller(94)
- Longitudinally polarized -> SSA – Brandenburg, Mueller, OT(95)
- Refined DA – Bakulev, Stefanis, OT(07)

From semi- to pure exclusive

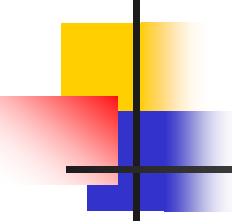
- Simplest case - pion FF(ERBL)
- Change DA to GPD- exclusive electroproduction (Frankfurt,Strikman)
- $M_{DY} \sim M_{DVCS} F_{\pi} g g^*$
- Time from right to left- exclusive DY (DAXGPD)- Berger,Diehl,Pire
- Phase sign change: c.f. Sivers for SIDIS/DY
- Analytic properties $\sim$ factorization (not completely clear yet)
- Second DA->GPD-another mechanism- OT'05 (problems with factorization -analytic continuation to be performed)





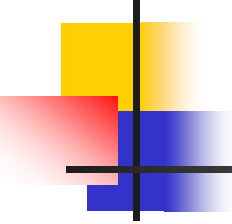
Backward DY –qqq TDA – transition from nucleon to meson

- J.P. Lansberg, B.Pire, K. Semenov-Tian-Shansky and L. Szymanowski, Phys. Rev. D86 (2012) 114033 [arXiv:1210.0126 [hep-ph]].
- Are negative x accessible?



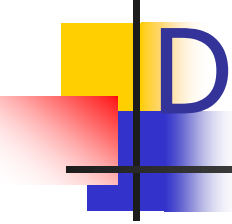
Kinematical regions for exclusive and semi-exclusive DY

- $(x_1 p_1 + x_2 p_2 + Q_T)^2 = Q^2$
- $x_1 x_2 = (Q_T^2 + Q^2)/s$
- $x_1 - x_2 = x_F$
- (Sudakov) Variables:
- $x_F, Q_T^2, Q^2 \leftrightarrow x_1, x_2, Q_T^2/Q^2$
- Semi-exclusive: $x_1 \rightarrow 1$
- Exclusive: $x_1 \rightarrow 1$ ($x_2 \rightarrow \text{skewness}$),
 $Q_T^2(t) \rightarrow 0$



E/SE DY

- SE – may be accessed at relatively large Q_T (supportive, moves x's to larger values)
- E requires selection of low Q_T
- S/SE separation at low Q_T - (tensor) polarization



Dilepton angular distribution

- Angular distribution

$$d\sigma \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi + \rho \sin 2\theta \sin \phi + \sigma \sin^2 \theta \sin 2\phi$$

- Positivity of the matrix (= hadronic tensor in dilepton rest frame)

$$M_0 = \begin{pmatrix} \frac{1-\lambda}{2} & \mu & \rho \\ \mu & \frac{1+\lambda-\nu}{2} & \sigma \\ \rho & \sigma & \frac{1+\lambda+\nu}{2} \end{pmatrix} \quad \begin{aligned} |\lambda| &\leq 1, \quad |\nu| \leq 1 + \lambda, \quad \mu^2 \leq \frac{(1-\lambda)(1+\lambda-\nu)}{4} \\ \rho^2 &\leq \frac{(1-\lambda)(1+\lambda+\nu)}{4}, \quad \sigma^2 \leq \frac{(1+\lambda)^2 - \nu^2}{4} \end{aligned}$$

- + cubic – $\det M_0 > 0$
- 1st line – Lam&Tung by SF method

Angular distributions – probes of DA

■ Unpolarized

$$F = \int_0^1 dy \frac{\varphi(y, \tilde{Q}^2)}{y},$$

$$I(\tilde{x}) = \int_0^1 dy \frac{\varphi(y, \tilde{Q}^2)}{y(y + \tilde{x} - 1 + i\varepsilon)}$$

$$\tilde{x}(x_L, \rho) \equiv \frac{x_L + \sqrt{x_L^2 + 4(1 + \rho^2)\tau}}{2(1 + \rho^2)}.$$

$$\rho \equiv Q_T/Q:$$

$$x_L = 2Q_L/\sqrt{s} < 1:$$

■ Polarized

$$\lambda(\tilde{x}, \rho) = \frac{2}{N} \{ (1 - \tilde{x})^2 [(\text{Im} I(\tilde{x}))^2 + (F + \text{Re} I(\tilde{x}))^2] - (4 - \rho^2) \rho^2 \tilde{x}^2 F^2 \}, \quad (2.19)$$

$$\mu(\tilde{x}, \rho) = -\frac{4}{N} \rho \tilde{x} F \{ (1 - \tilde{x}) [F + \text{Re} I(\tilde{x})] + \rho^2 \tilde{x} F \}, \quad (2.20)$$

$$\nu(\tilde{x}, \rho) = -\frac{8}{N} \rho^2 \tilde{x} (1 - \tilde{x}) F [F + \text{Re} I(\tilde{x})], \quad (2.21)$$

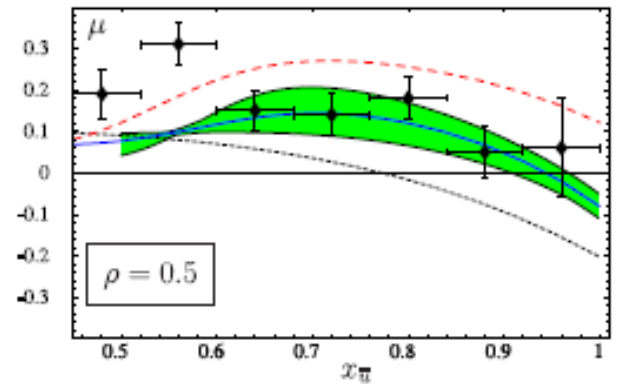
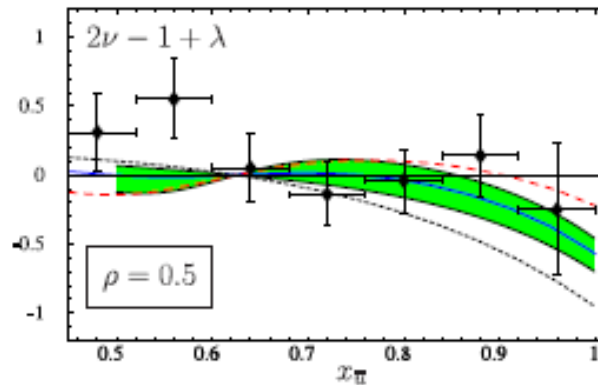
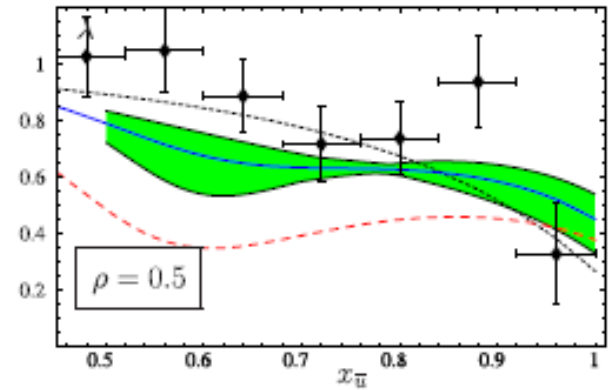
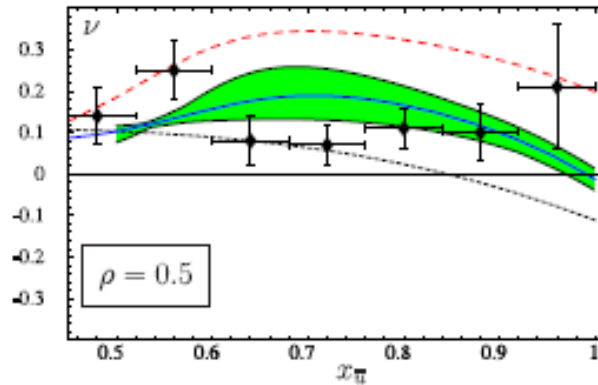
$$N(\tilde{x}, \rho) = 2 \{ (1 - \tilde{x})^2 [(\text{Im} I(\tilde{x}))^2 + (F + \text{Re} I(\tilde{x}))^2] + (4 + \rho^2) \rho^2 \tilde{x}^2 F^2 \} \quad (2.22)$$

$$\bar{\mu}(\tilde{x}, \rho) = \frac{-2\pi s_e \rho \tilde{x} F \varphi(\tilde{x}, \tilde{Q}^2)}{(1 - \tilde{x})^2 [(F + \text{Re} I(\tilde{x}))^2 + \pi^2 \varphi(\tilde{x})^2] + (4 + \rho^2) \rho^2 \tilde{x}^2 F^2} \bar{\mu}_{\text{nucl}},$$

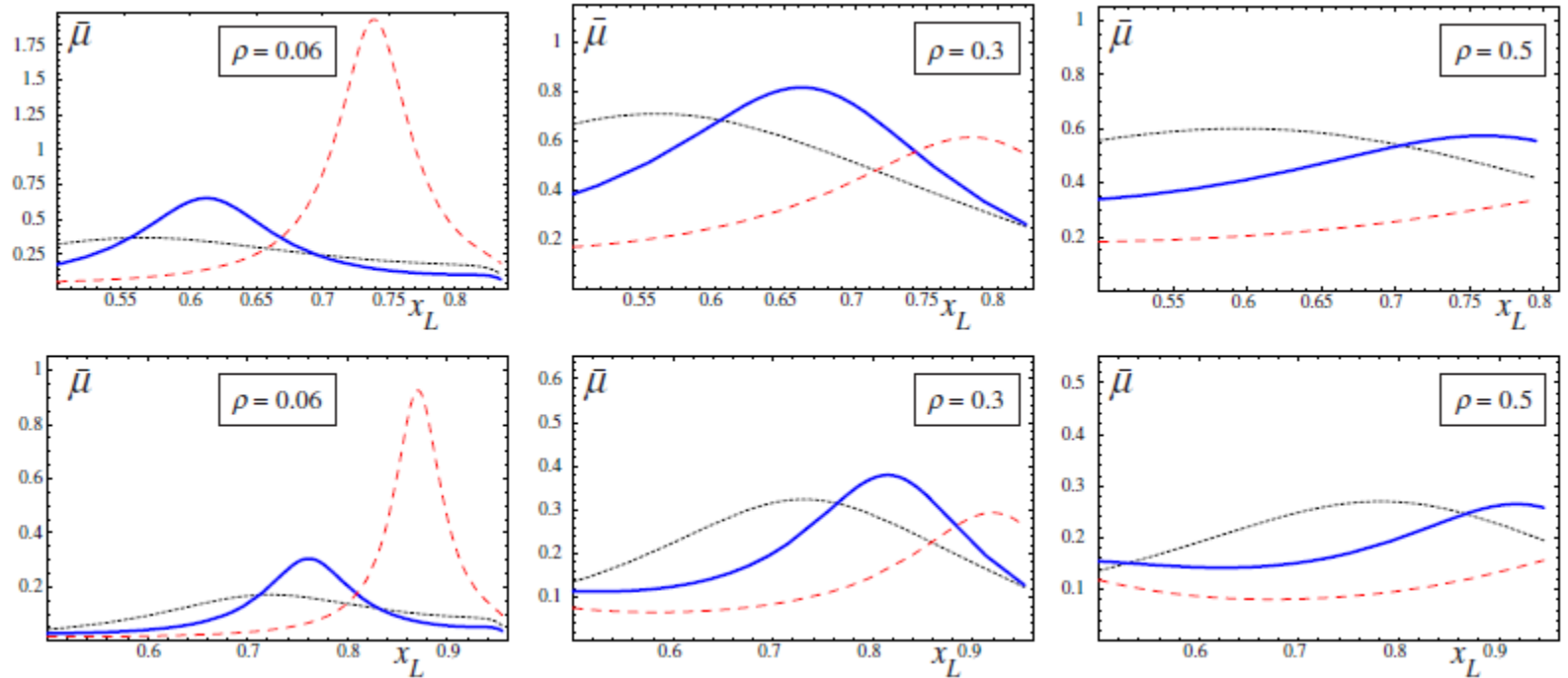
$$\bar{\mu}_{\text{nucl}} \equiv \frac{\frac{4}{9} \Delta q_u^v(x_p; \mu^2) + \frac{4}{9} \Delta q_u^s(x_p; \mu^2) + \frac{1}{9} \Delta q_d^s(x_p; \mu^2)}{\frac{4}{9} q_u^v(x_p; \mu^2) + \frac{4}{9} q_u^s(x_p; \mu^2) + \frac{1}{9} q_d^s(x_p; \mu^2)},$$

$$\bar{\nu}(\tilde{x}, \rho) = 2\rho \bar{\mu}(\tilde{x}, \rho),$$

Asymmetries vs data



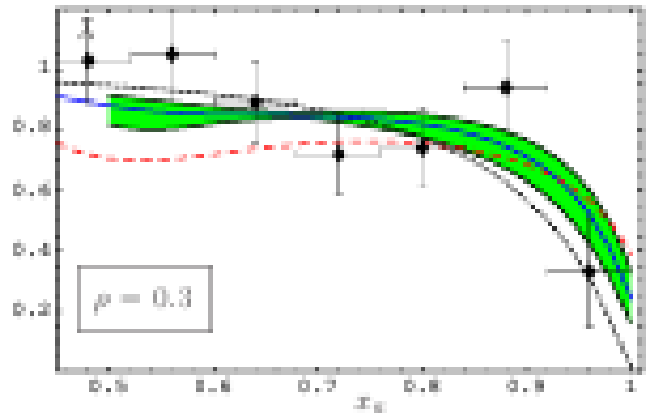
Longitudinal Target Polarization (no at COMPASS, possible at J-PARC?) -> scan of DA



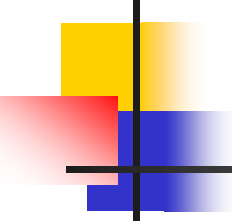
The upper row corresponds to a center-of-mass energy $s = 100 \text{ GeV}^2$.

Properties of exclusive DY

- Polarization T- $\rightarrow$ L (collinear limit)



- More detailed calculations of angular distributions required



CONCLUSIONS

- SE/E – various region of transverse momentum
- Small transverse momentum – angular distributions crucial

Polarizabilities of pions, kaons..?

- Lattice studies for pions, rho (kaons are planned)

[Magnetic polarizabilities of light mesons in \$SU\(3\)\$ lattice gauge theory](#)

[E.V. Lushevskaya](#) (Moscow, ITEP), [O.V. Teryaev](#) (Dubna, JINR), [O.A. Kochetkov](#) (Moscow, ITEP & Regensburg U.). Nov 16, 2014. 8 pp.

e-Print: [arXiv:1411.4284](#) [hep-lat]

- $g \sim 2$
- polarizabilities
- No tachyonic mode

