

# Charmed hadron molecules

## --what we have learnt from the study of the X(3872) properties--

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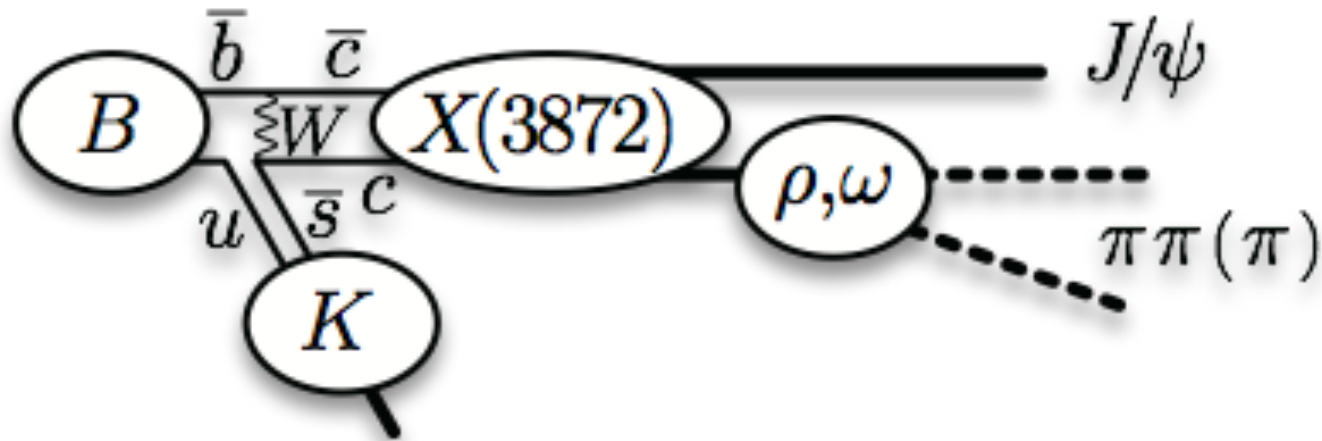
KEK theory center workshop on Hadron physics with high-momentum hadron beams at J-PARC, March 13-16, 2015, KEK, Japan

# Outline

- X(3872): experimental results
- Charmonium-molecule hybrid model of X(3872)
- Summary of X(3872)
- Consideration about other hadronic molecules:  $Z_b^+(10610)$ ,  $Z_b^+(10650)$ ,  $Z_c^+(3900)$
- At J-PARC

# X(3872)

- First observed by B decay



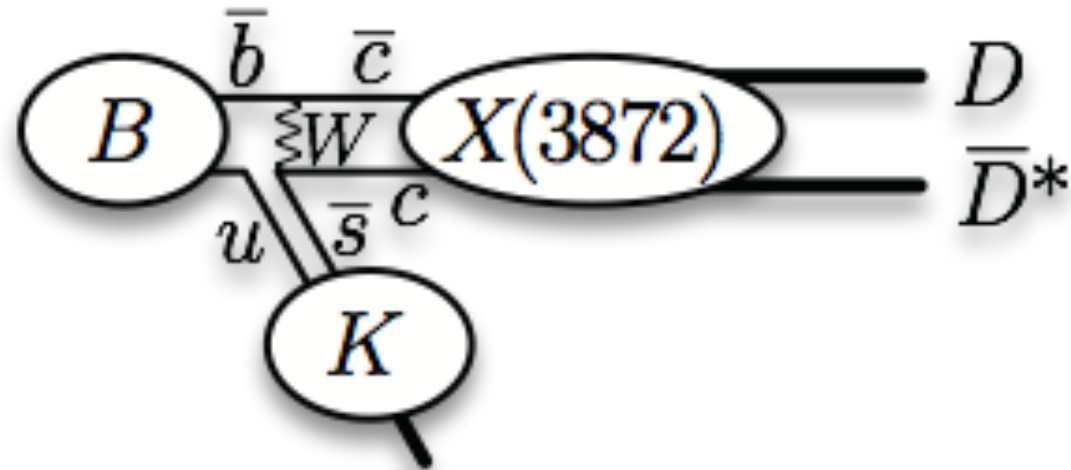
X(3872) should be **Charmonium-like** state.

$J/\psi \pi^n$  Exp:

S.K.Choi et al. [Belle] PRL91, 262001 (2003)

B.Aubert et al. [BaBar] PRD71, 071103 (2005)

# X(3872)



X(3872) decays to  $DD^*_{\text{bar}}$

$$\text{Br}(X \rightarrow D^0 \bar{D}^{*0}) / \text{Br}(X \rightarrow J/\psi \pi^2)$$

$$= 8.92 \pm 2.42, 19.9 \pm 8.05 \text{ (calc from papers)}$$

Belle

BABAR

$D^0 \bar{D}^{*0}$  exp

T.Aushev et al. [Belle] PRD81, 031103 (2010)

B.Aubert et al. [BaBar] PRD77, 011102 (2008)

# Formation process of X(3872)

- X(3872) from

- B<sup>+</sup> decay

- B<sup>0</sup> decay

- pp<sup>bar</sup> collision

p̄p exp

D.Acosta et al. [CDF] PRL93, 072001 (2004)

V.M.Abazov et al. [D0] PRL93, 162002 (2004)

- $\text{Br}(pp^{\text{bar}} \rightarrow X) / \text{Br}(pp^{\text{bar}} \rightarrow \psi(2S)) > 0.046$

C.Bignamini et al. PRL103, 162001 (2009)

- pp collision

- $J^{PC} = 1^{++}$

pp exp

R.Aaij et al. [LHCb] EPJC72, 1972 (2012)

# Decay process of $X(3872)$

- $X(3872)$  decays into
  - $J/\psi \pi\pi$  [Belle][BaBar][CDF][D0][CMS]
  - $J/\psi \pi^3$  [BaBar]([Belle] not published)
  - $D^0 \bar{D}^{*0}$  [Belle][BaBar]
  - $J/\psi \gamma$  [Belle][BaBar][LHCb]
  - $\psi(2S) \gamma$  [BaBar][LHCb]([Belle] not seen)

# X(3872) facts

4 two-meson thresholds are nearby.

- $J^{PC}=1^{++}$
- X(3872) ( $uc\bar{u}\bar{c}$ ,  $dcd\bar{c}$ ) thresholds
  - $D^{\pm}D^{*\mp}$   $3879.87\pm 0.12$  MeV
  - $J/\psi\omega$   $3879.57\pm 0.12$  MeV
  - $J/\psi\rho$   $3872.18\pm 0.25$  MeV
  - $D^0\bar{D}^{*0}$   $3871.80\pm 0.12$  MeV (**Lowest threshold**)

# Mass and Decay to $J/\psi \pi^n$

- $X(3872)$  decays into

- ▷  $J/\psi \pi\pi$  **isovector**

- ▶ mass  $3871.69 \pm 0.17$  MeV

- width  $< 1.2$  MeV

very narrow width

- ▷  $J/\psi \pi^3$  **isoscalar**

- ▶  $\text{Br}(X \rightarrow J/\psi \pi^3) / \text{Br}(X \rightarrow J/\psi \pi^2) =$

- $1.0 \pm 0.4 \pm 0.3, 0.8 \pm 0.3$

Belle

BABAR

A large isospin symmetry breaking

Binding energy  
is **0.11** MeV.



# Decay to $D^0\bar{D}^{*0}$

- X(3872) decays into  $D^0\bar{D}^{*0}$

▷ mass       $3872.9^{+0.6+0.4}_{-0.4-0.5}$  MeV      Belle  
                  $3875.1^{+0.7}_{-0.5\pm0.5}$  MeV      BABAR

mass from  $D^0\bar{D}^{*0}$  > mass from  $J/\psi\pi\pi$

▷ width       $3.9^{+2.8+0.2}_{-1.4-1.1}$  MeV      Belle  
                  $3.0^{+1.9}_{-1.4\pm0.9}$  MeV      BABAR

width from  $D^0\bar{D}^{*0}$  > width from  $J/\psi\pi\pi$

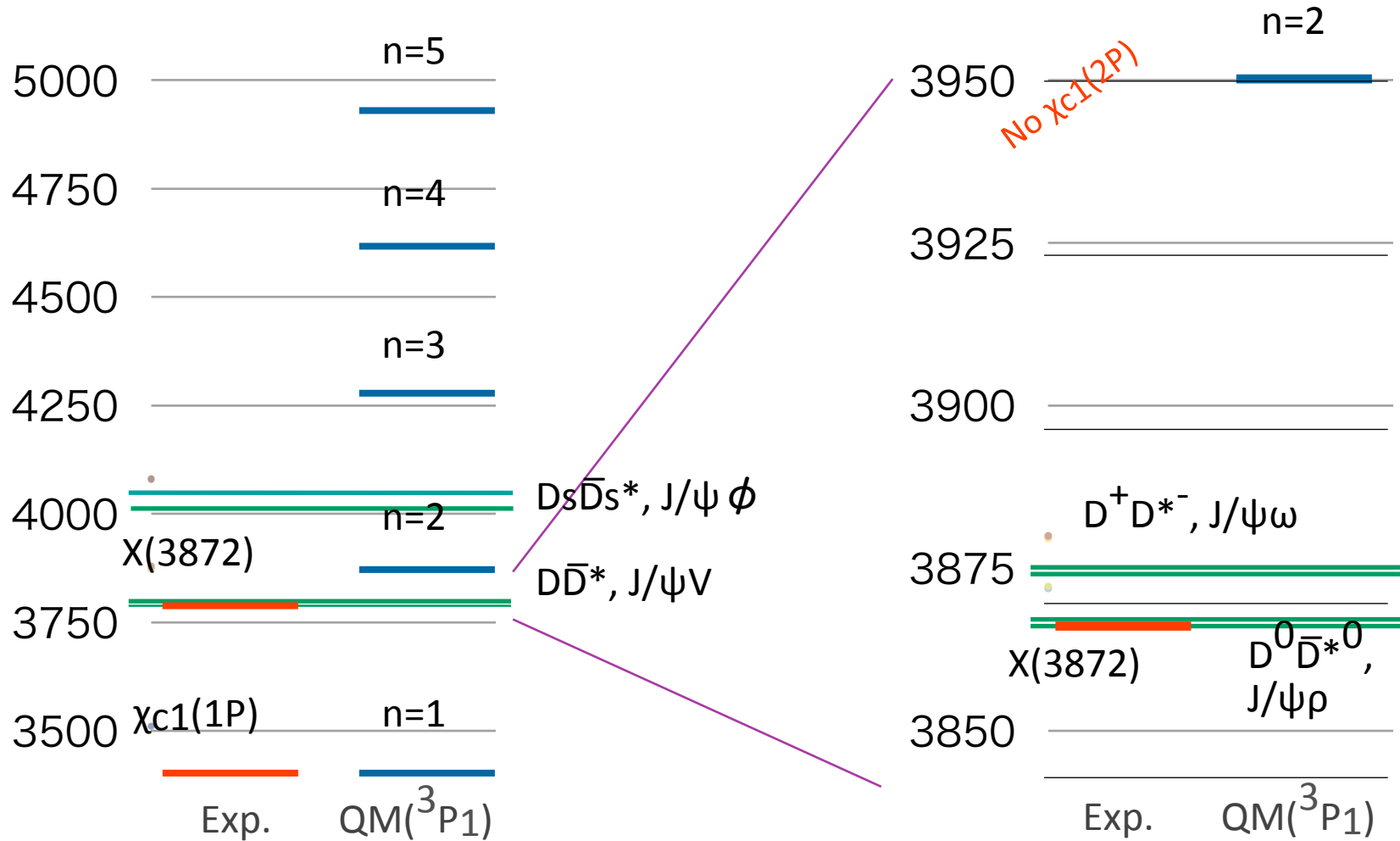
# X(3872)

- Charmonium:  $1^{++}$  corresponds  $^3P_1$  state

$\chi_{c1}(1P)$       Mass:  $3510.66 \pm 0.07$  MeV

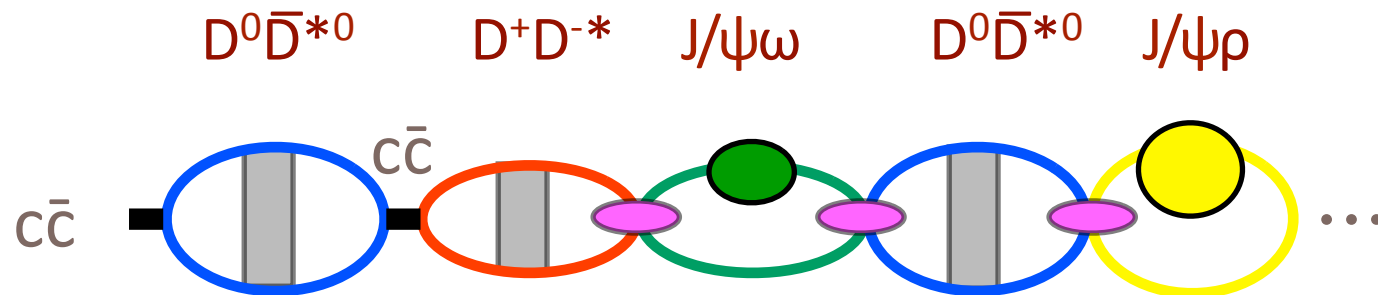
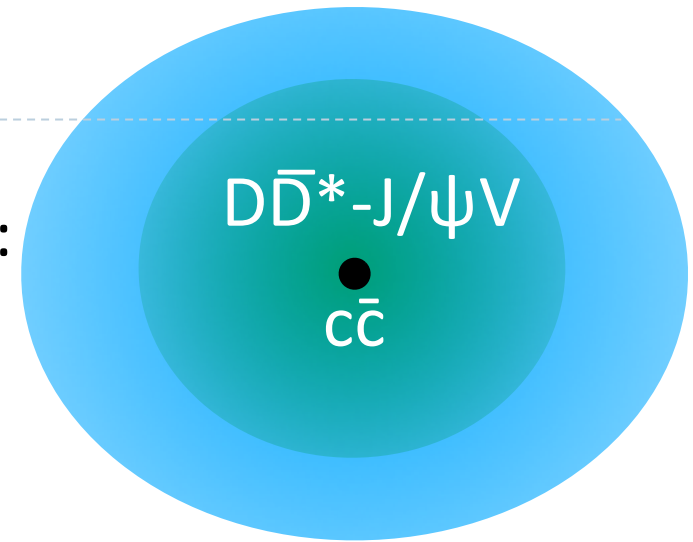
$\chi_{c1}(2P)$       Mass: Quark model prediction  
NR: 3925 MeV  
GI: 3953 MeV      **Not observed**

# 1<sup>++</sup> Spectrum



# Our picture of $X(3872)$

- Two-meson molecule with a  $c\bar{c}$  core:
  - ▷  $c\bar{c} - D^0\bar{D}^{*0} - D^+D^{*-} - J/\psi\omega - J/\psi\rho$
  - ▷  $\omega$  and  $\rho$  have width.
  - ▷  $J/\psi\omega$  and  $J/\psi\rho$  couple to  $c\bar{c}$  only via  $D\bar{D}^*$  channels (OZI).



M. Takizawa and S. Takeuchi, Prog. Theor. Exp. Phys. 2013, 093D01  
 S.Takeuchi, K.Shimizu, and M.Takizawa, Prog. Theor. Exp. Phys. 2014, 123D01

# Model Hamiltonian

- $H = H_0 + V$

$$H_0 = \begin{pmatrix} H_0^{(P)} & 0 \\ 0 & E_0^{(Q)} \end{pmatrix} \quad H_0^{(P)} = \sum_i \left( M_i + m_i + \frac{k_i^2}{2\mu_i} \right)$$

$$V = \begin{pmatrix} V_{PP} & V_{PQ} \\ V_{QP} & 0 \end{pmatrix}$$

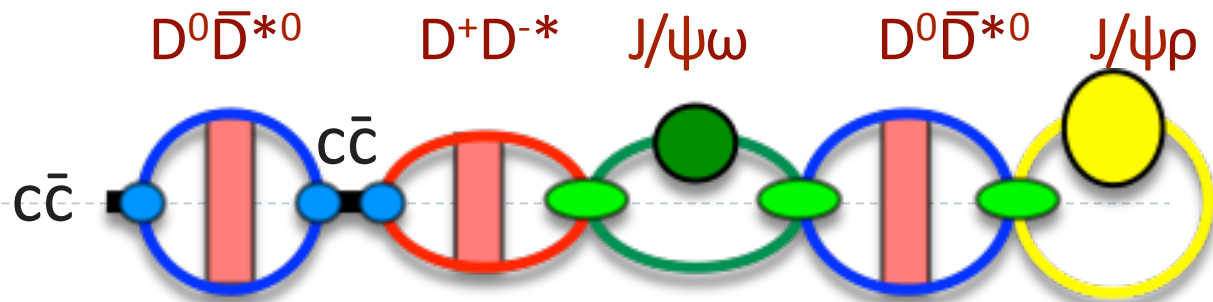
P :  $D^0\bar{D}^{*0}$  -  $D^+D^{*-}$  -  $J/\psi \omega$  -  $J/\psi \rho$ , Q :  $c\bar{c}$

$V_{PP}$  : potential between the two mesons

$V_{PQ}$  : coupling between the two-meson and  $c\bar{c}$  channels

$E_0^{(Q)}$  : QM predicted  $\chi_{c1}(2P)$  mass = 3950 MeV

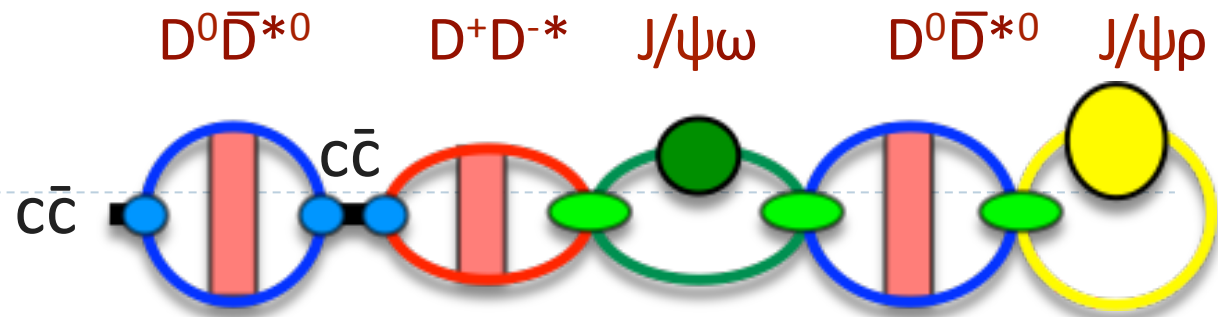
# Potentials



- Interaction strengths are determined by

- ▶  $D \bar{D}^* - J/\psi \omega$  coupling  $u$ : ← the quark model
  - ▶  $q \bar{q}$  interaction to obtain this transfer strength. (Godfrey-Isgur)
  - ▶ RGM for QM  $\rightarrow V_{PP} = N^{-1/2} H N^{-1/2} - H_0$

# Potentials



- Interaction strengths are determined by
  - ▷  $D\bar{D}^*$  interaction  $v$  :
    - ▶ Param set  $A$ : maximum attraction which does not make a bound state in  $B\bar{B}^*$  systems.
    - ▶ Param set  $QM$ : use a quark model value
  - ▷  $D\bar{D}^* - c\bar{c}$  coupling  $g$  :
    - ▶ a free parameter to give the  $X(3872)$  peak energy.

# X(3872) : a bound state

- Density distri. of two mesons

▷  $r^2 |\phi_{MM}(r)|^2$

- Probability

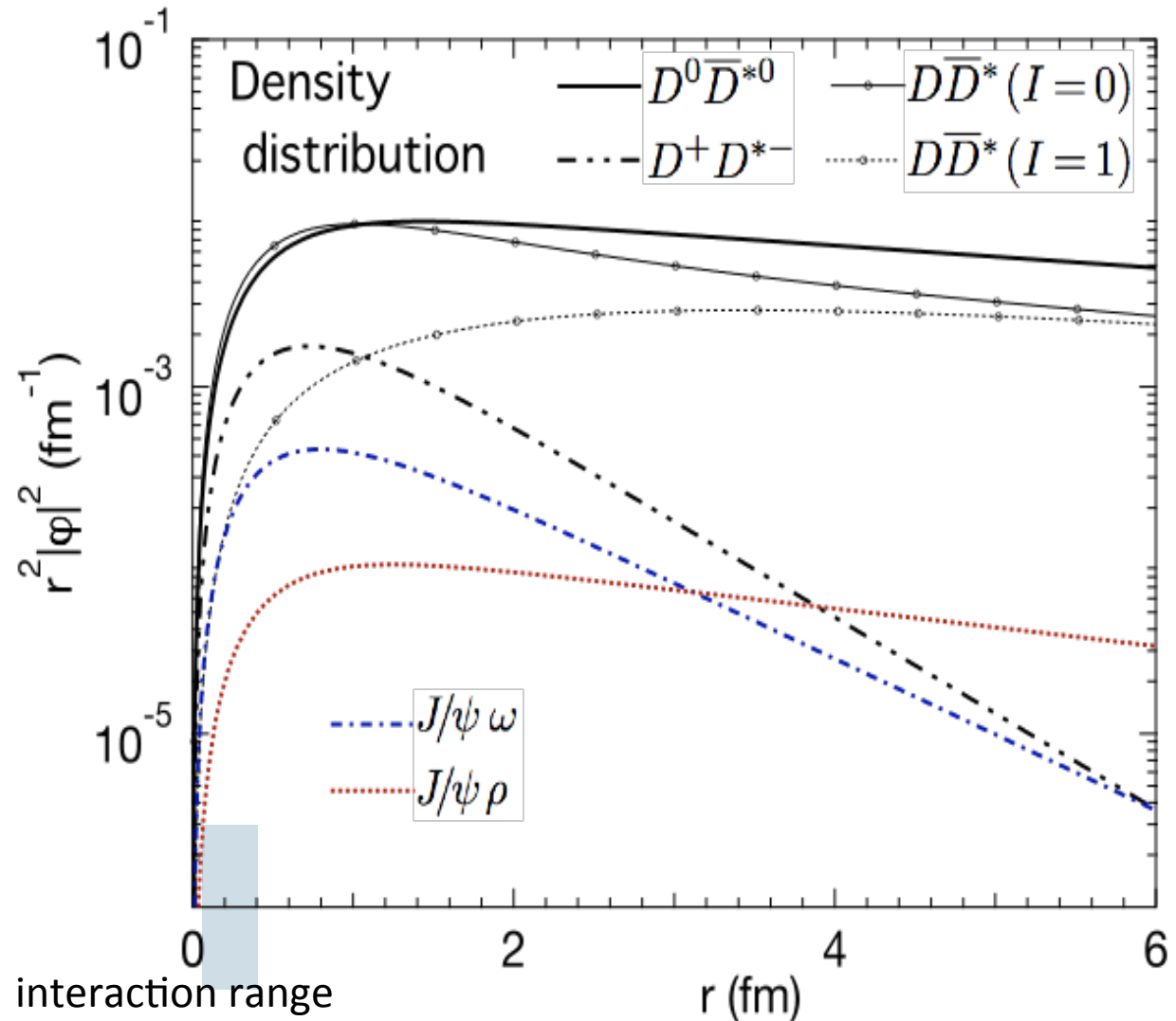
▷ 0.036  $c\bar{c}$

▷ 0.913  $D^0\bar{D}^{*0}$

▷ 0.034  $D^+D^{*-}$

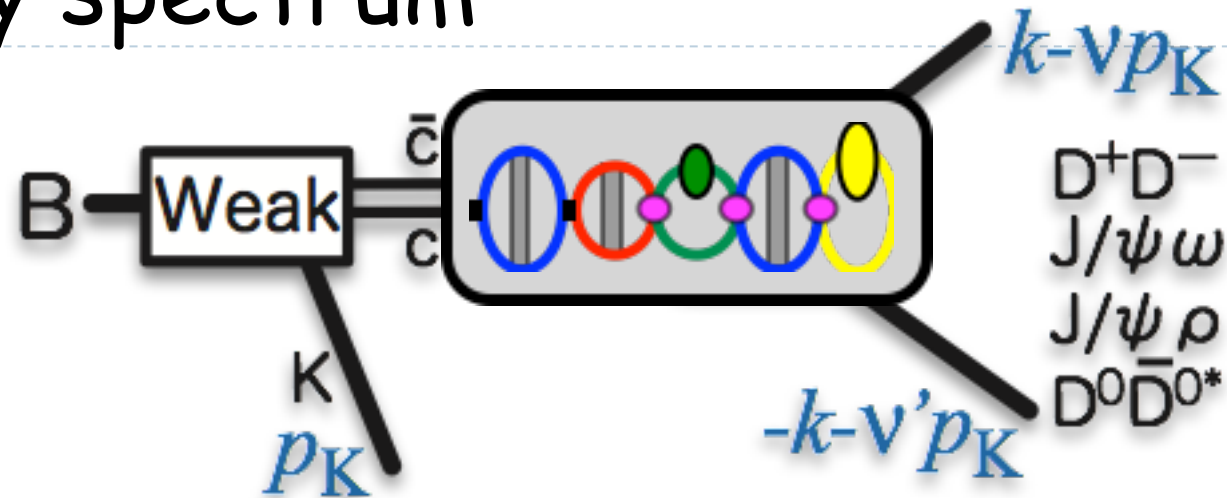
▷ 0.010  $J/\psi\omega$

▷ 0.006  $J/\psi\rho$





# B Decay spectrum



- Decay rate  $\propto c\bar{c}$  self energy

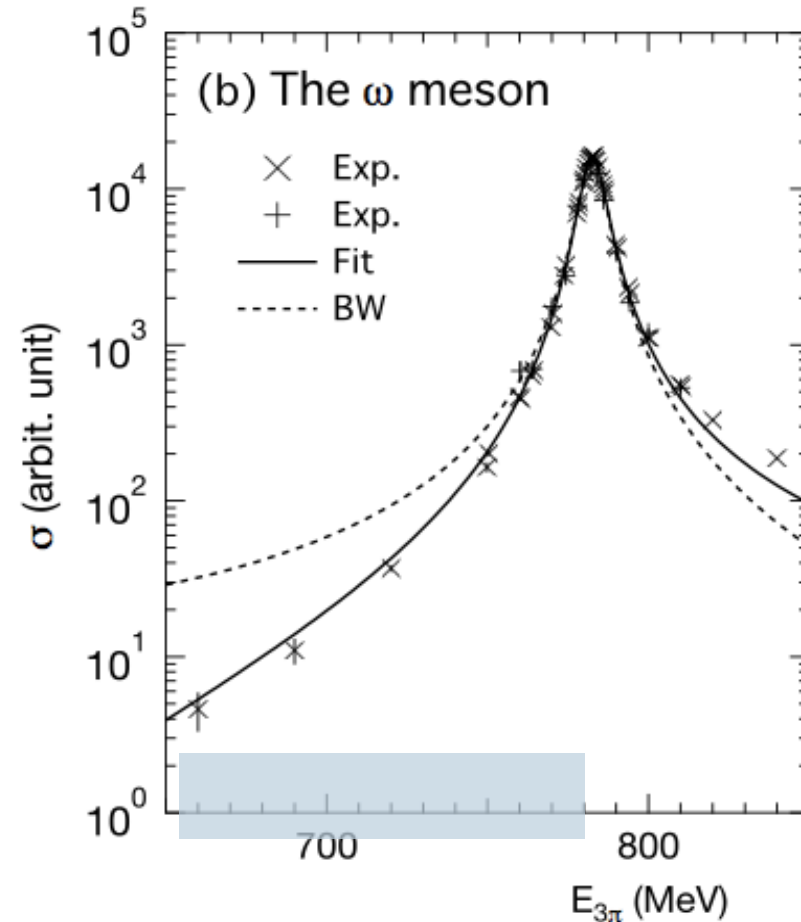
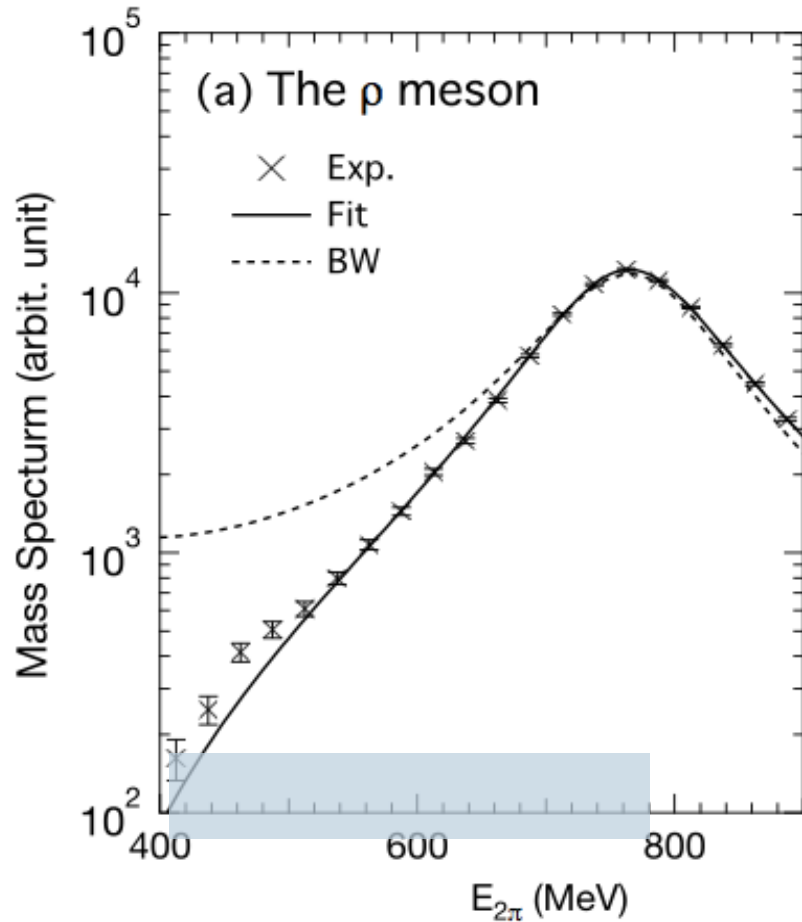
$$\frac{dW}{dE} = -\frac{1}{\pi} \text{Im} \langle c\bar{c} | G_Q(E) | c\bar{c} \rangle \quad \int_{E_{\text{threshold}}}^{\infty} \frac{dW}{dE} dE = 1$$

$$\frac{dW(c\bar{c} \rightarrow f)}{dE} = \frac{2}{\pi} \mu_f \int \frac{k^2 dk \mu_f \Gamma_f}{(k_f^2 - k^2)^2 + (\mu_f \Gamma_f)^2} \left| \langle \langle f; k | V_{PQ} \tilde{G}_Q | c\bar{c} \rangle \right|^2$$

$$\tilde{G} = \frac{1}{E - H + i\Gamma(E)/2} \leftarrow \text{observed } \rho \rightarrow \pi\pi \text{ or } \omega \rightarrow \pi\pi\pi \text{ width}$$

# Vector meson widths

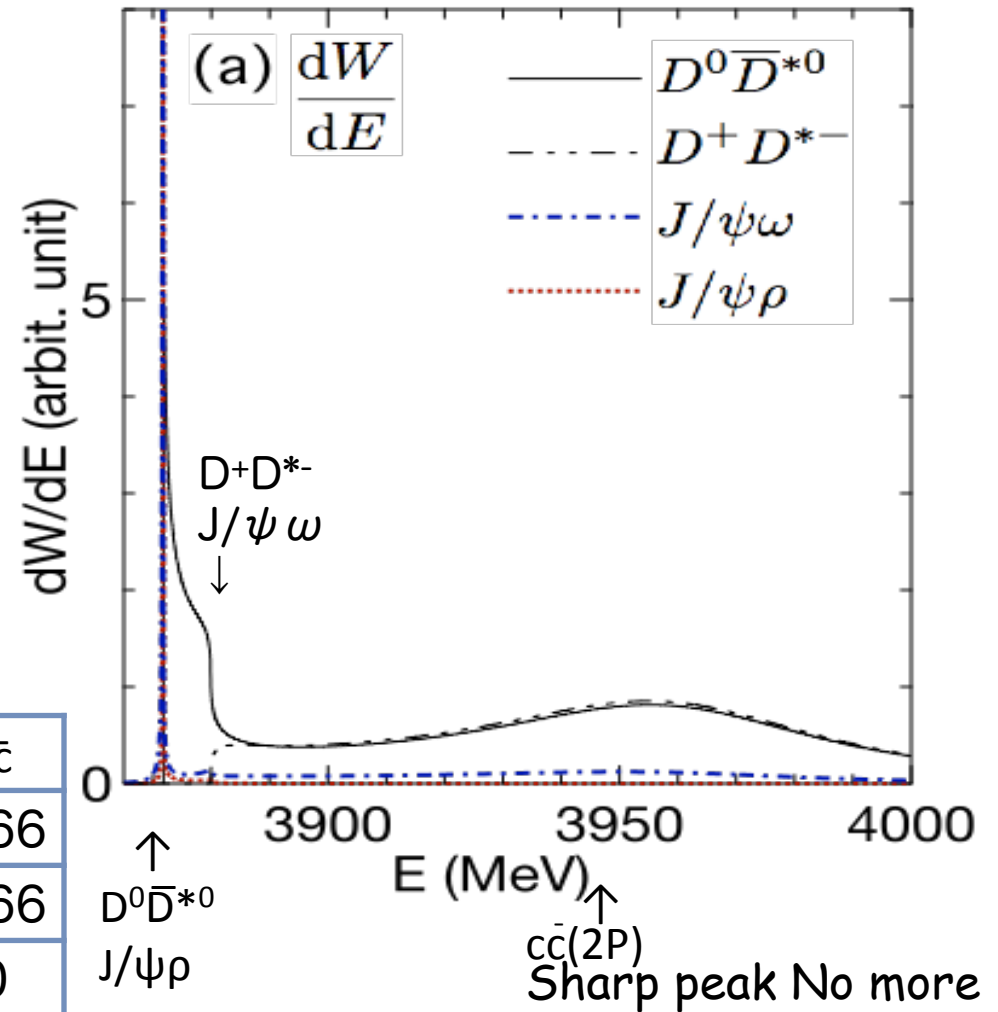
- Energy dependent width for  $\rho$  and  $\omega$



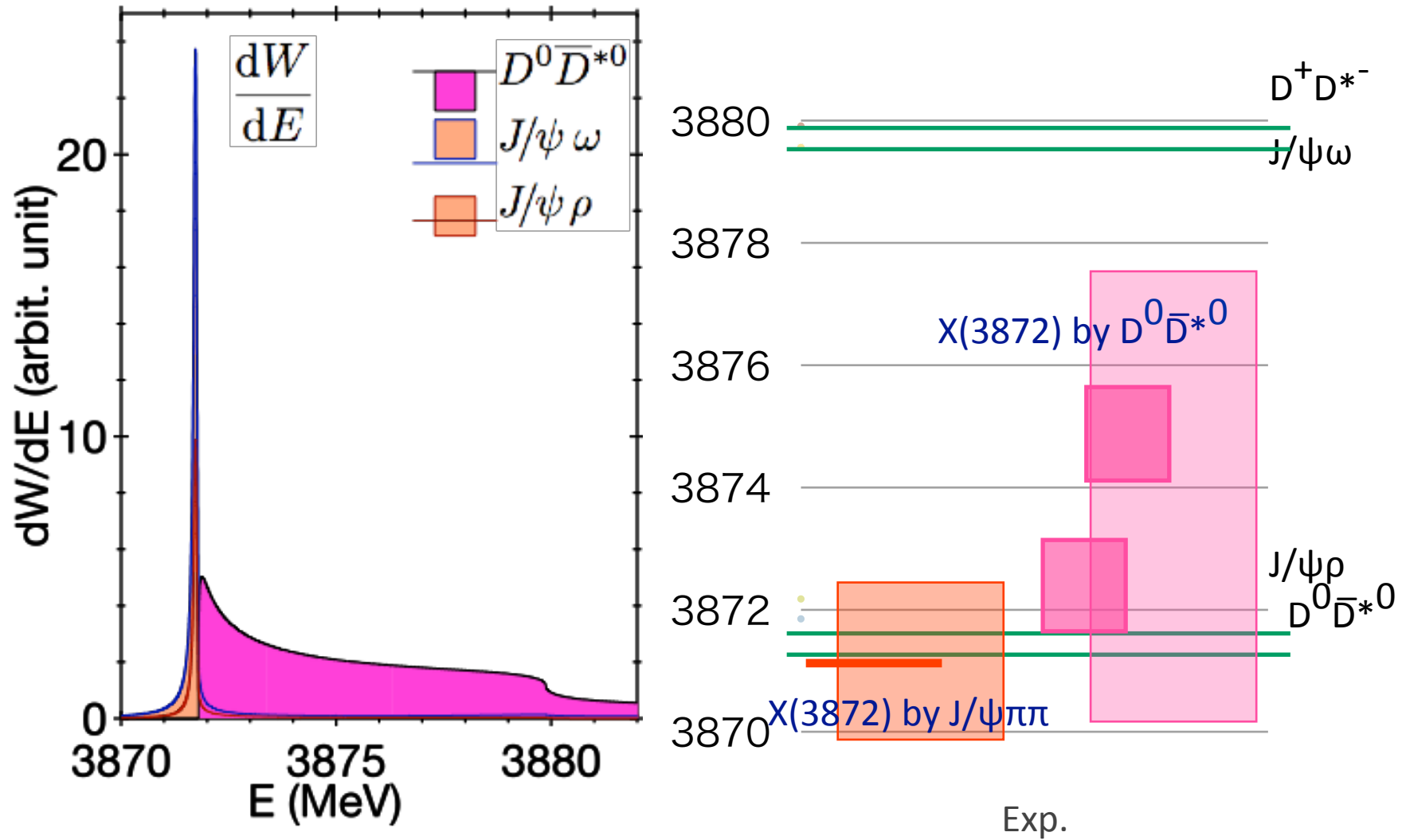
# $c\bar{c} \rightarrow$ two mesons : for 'bound' X(3872)

- $D^0\bar{D}^{0*}$  peak appears at the threshold.
- $J/\psi\omega$  and  $J/\psi\rho$  peaks appear at the threshold.
- $J/\psi\omega$  and  $J/\psi\rho$  are comparable.

$v,u,g^2$	$D^0\bar{D}^{0*}$	$D^+D^{*-}$	$J/\psi\rho$	$J/\psi\omega$	$c\bar{c}$
$D^0\bar{D}^{0*}$	-0.19	0	0.19	0.19	0.66
$D^+D^{*-}$		-0.19	-0.19	0.19	0.66
$J/\psi\rho$			0	0	0
$J/\psi\omega$				0	0

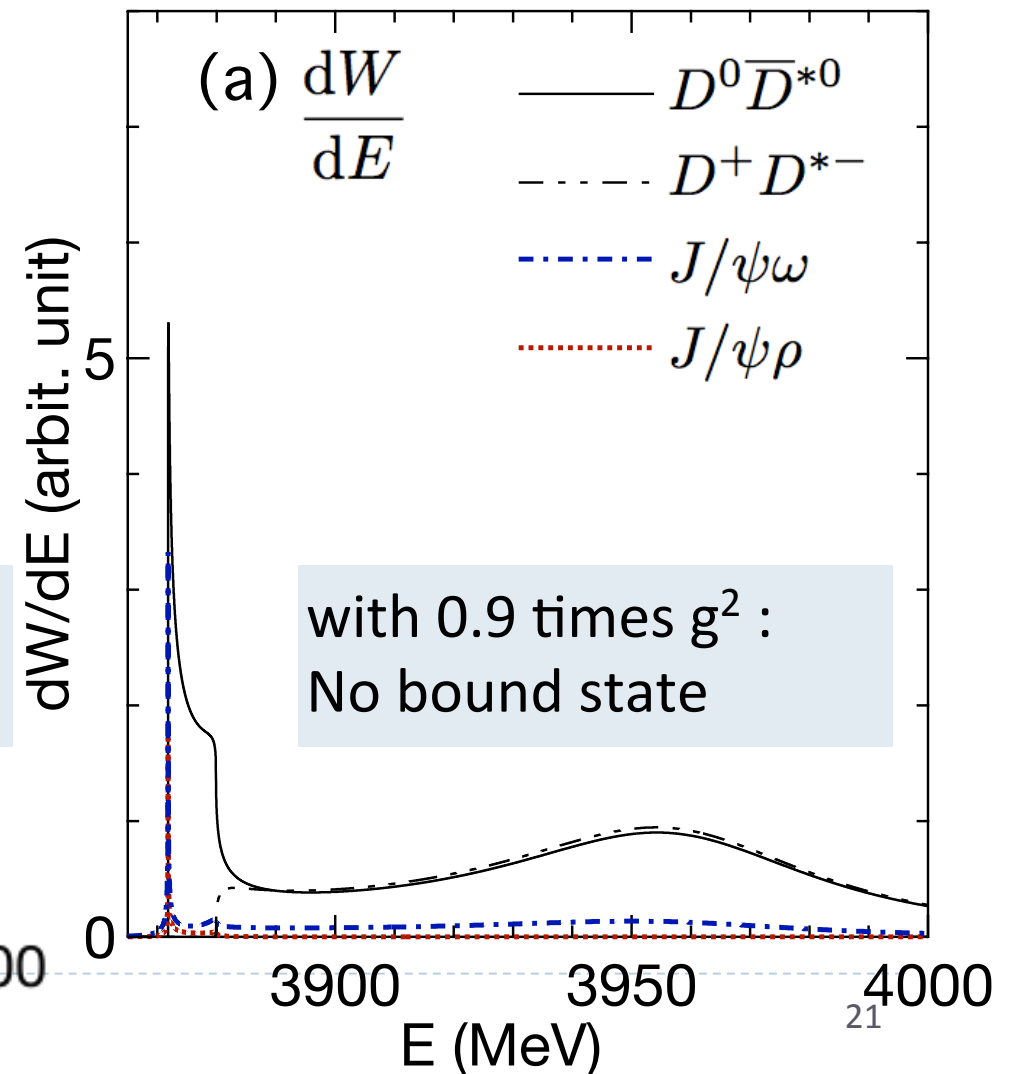
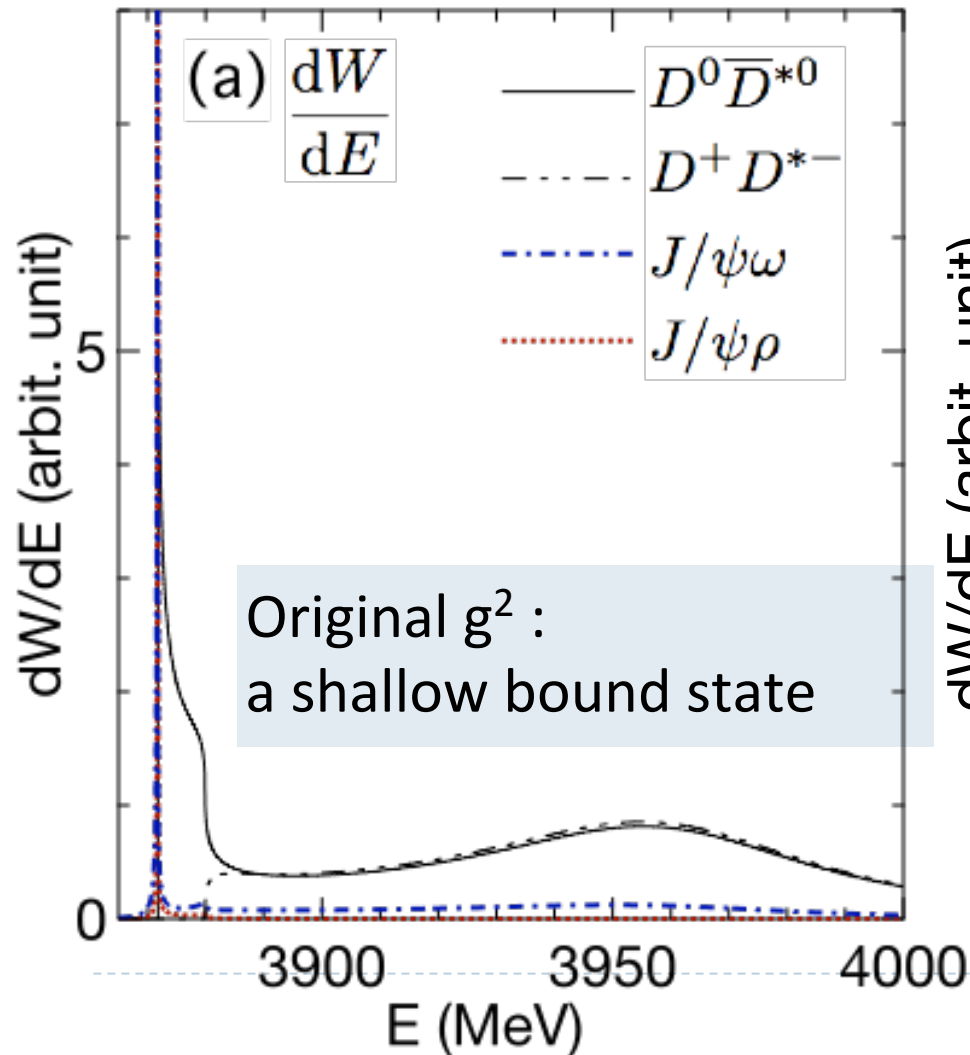


# 1<sup>++</sup> Spectrum



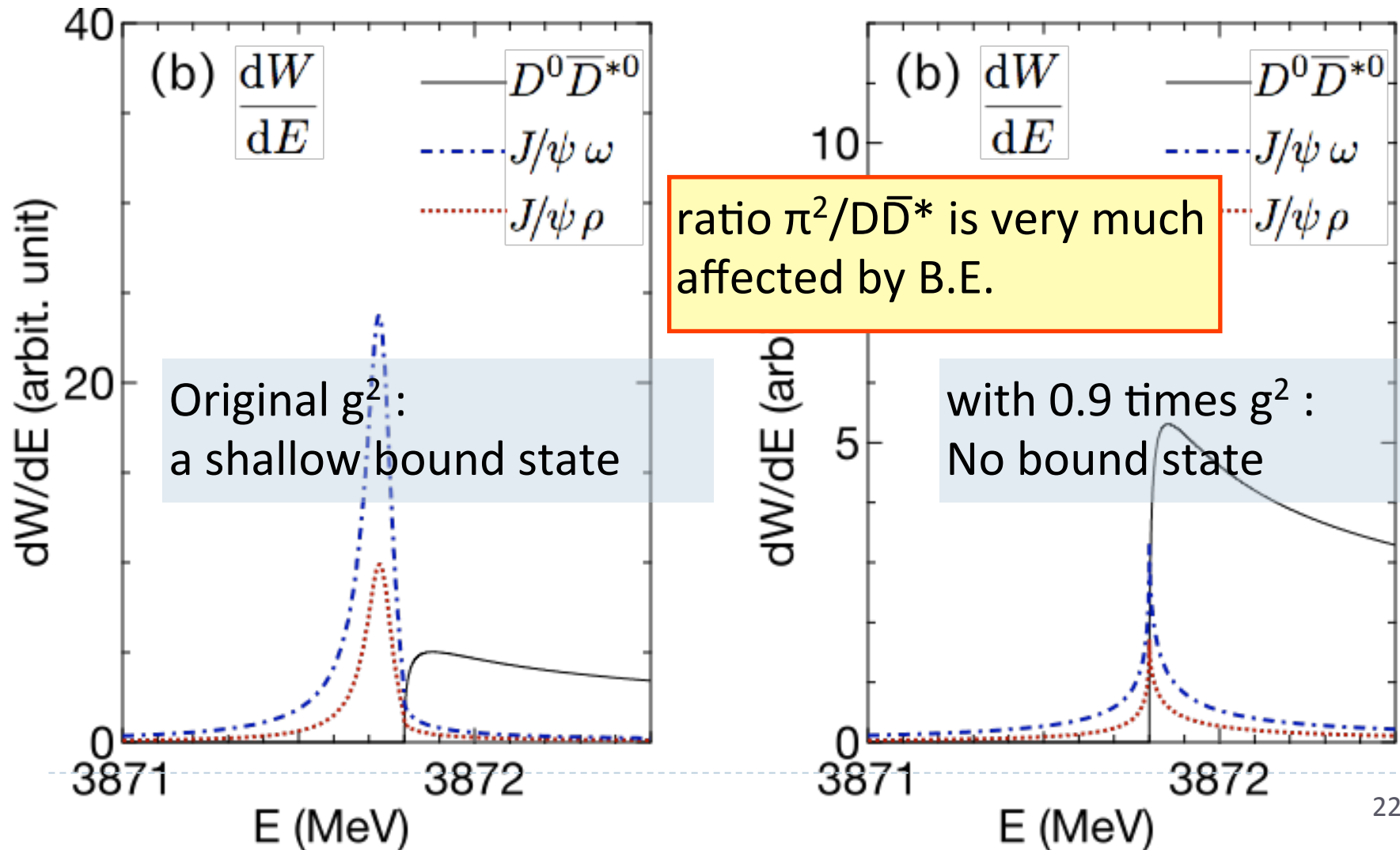
# X(3872) needs to be a bound state?

- No... it can be a virtual state



# X(3872) needs to be a bound state?

- A virtual state can make  $J/\psi \pi^n$  peak.



# Summary of X(3872)

- In order to reproduce the properties of the X(3872), **introduction of the all thresholds near the peak position is important.**
- $c\bar{c} - D^0\bar{D}^{*0} - D^+D^{*-} - J/\psi\omega - J/\psi\rho$  **full coupled channel calculation is necessary.**
- **$c\bar{c}$  component** may explain the production rate of X(3872) in the  $pp^{\text{bar}}$  collisions.
- **$c\bar{c}$  component** gives rise to the additional **attractive interaction** effectively.

# Summary of X(3872)

- Large rho-meson width enhances the isospin symmetry breaking.
- Calculation of the experimentally observed spectrum is important.

There is a mechanism to have narrow peaks of  $J/\psi\rho,\omega$  and broader peaks of  $D\bar{D}^*$ .



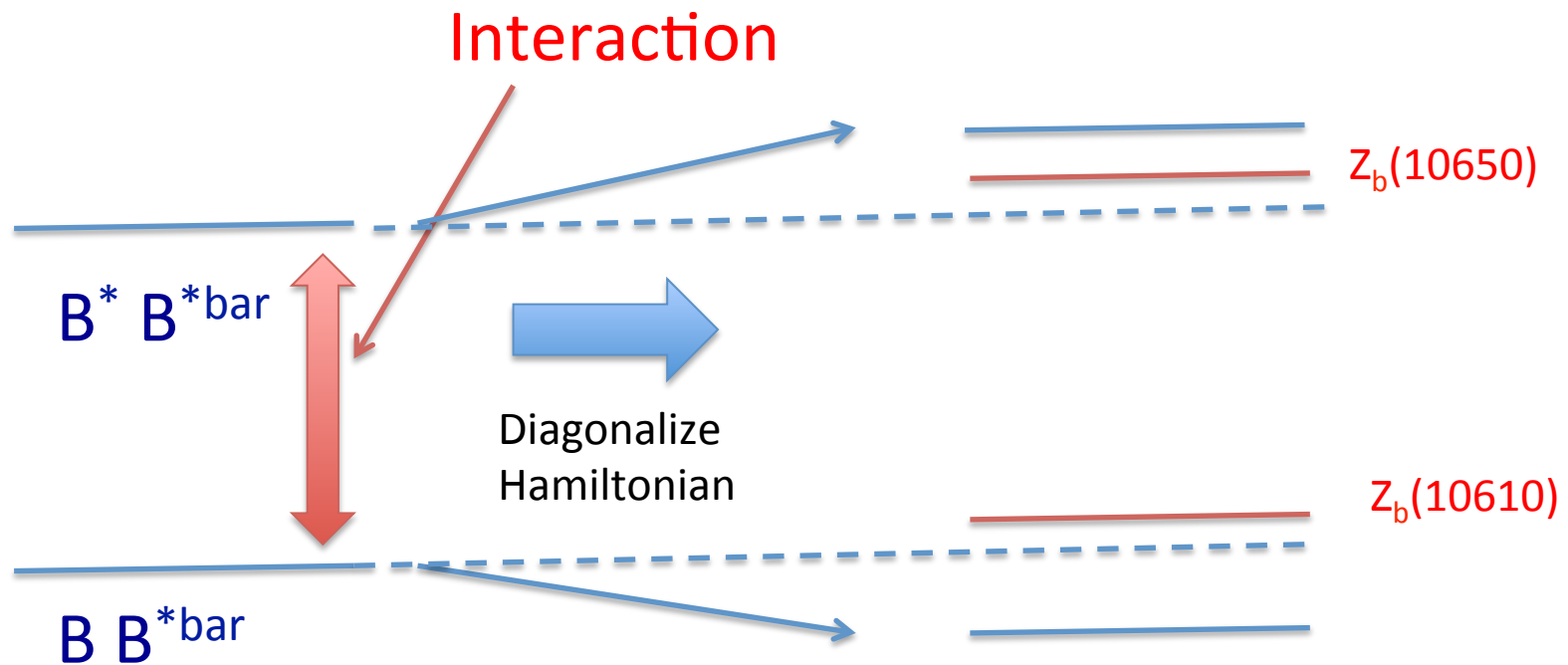
# Summary of X(3872)

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- A large Isospin breaking can be realized.
  - ▷ sensitive to  $(g/g_0)^2$
- $\frac{\Gamma(D^0 \bar{D}^{0*})}{\Gamma(J/\psi \rho)}$  probe of the binding energy of X(3872)
- $D\bar{D}^*$  int, g,  $\rightarrow$  other heavy quark systems  
 $D\bar{D}^*$  int  $\sim$   $B\bar{B}^*$  int (heavy quark symmetry)  
 $\rightarrow$  roughly consistent with  $Z_b^+(10610)$ ,  $Z_b^+(10650)$

# Consideration about other hadronic molecules

- $Z_b^+(10610)$  and  $Z_b^+(10650)$



# Consideration about other hadronic molecules

- $Z_c^+(3900)$
- Simultaneous explanation of  $Z_b^\pm$  and  $Z_c^\pm$  is rather difficult since heavy quark symmetry requires almost same interactions between  $DD^{*\bar{}}$  and  $BB^{*\bar{}}$ , but the kinetic energy of the  $BB^{*\bar{}}$  system is much smaller than that of  $DD^{*\bar{}}$  system.

# At J-PARC

- So far, hadronic molecular states have not been observed at the fixed target experiments.
- Target mass dependence of the production rate and the width of the hadronic molecule state may differ from that of the usual hadrons.

Further theoretical studies are necessary.

# Backup

**Radiative decays of  $X(3872)$**

# Radiative decays of $X(3872)$

$$R_\gamma = \frac{B(X(3872) \rightarrow \psi(2S) \gamma)}{B(X(3872) \rightarrow J/\psi \gamma)}$$

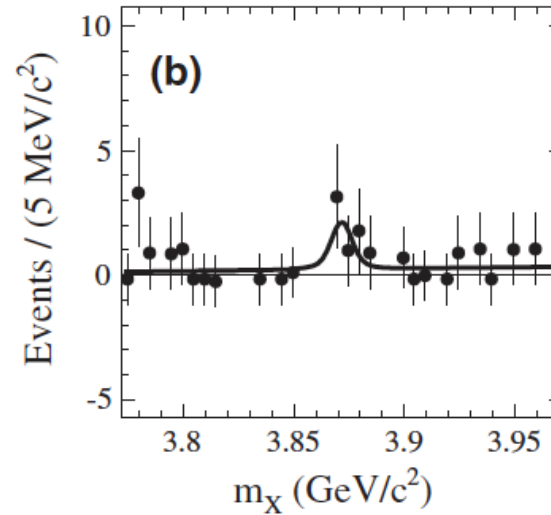
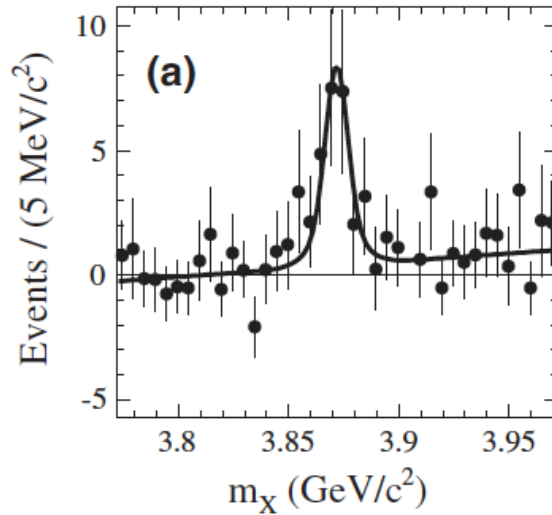
BaBar:  $3.4 \pm 1.4$  ( $3.5\sigma$ ) PRL102, 132001(2009)

Belle:  $< 2.1$  (90%CL) PRL107, 091803(2011)

LHCb:  $2.46 \pm 0.64(\text{stat}) \pm 0.29(\text{sys})$  ( $4.4\sigma$ )  
NP B886, 665 (2014)

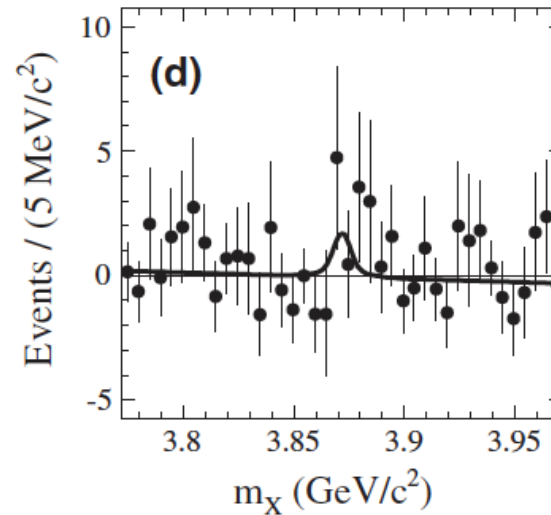
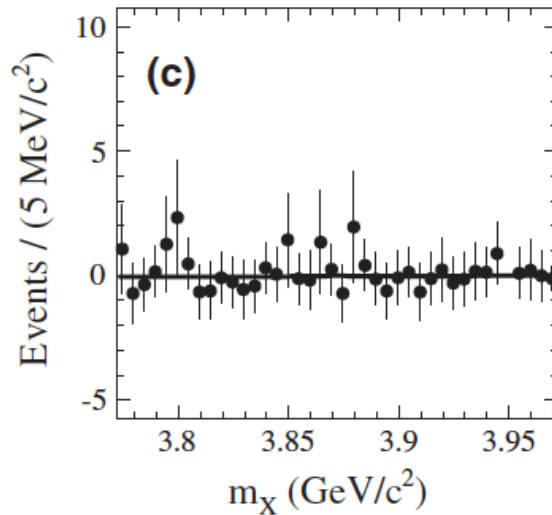
# BaBar's results ( $J/\psi\gamma$ )

$B^{\pm} \rightarrow X K^{\pm}$



$K_S^0$

$K^{*\pm}$

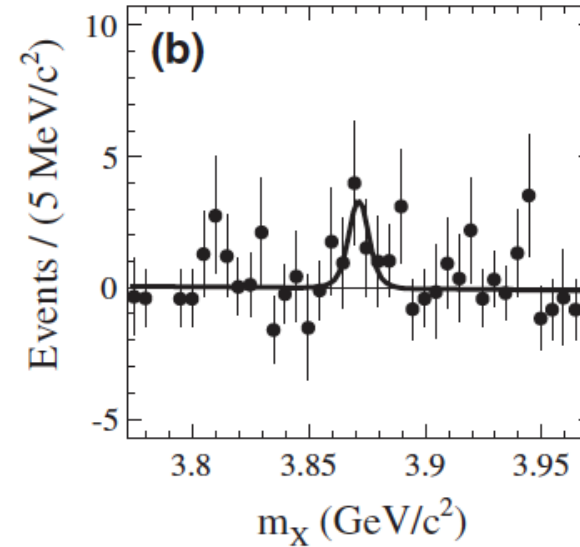
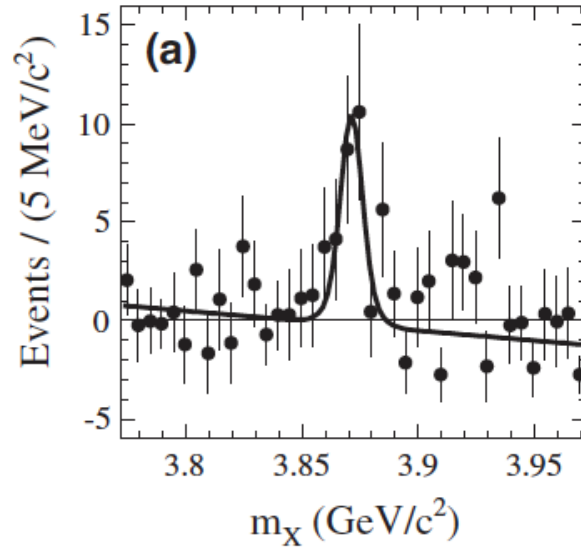


$K^{*0}$



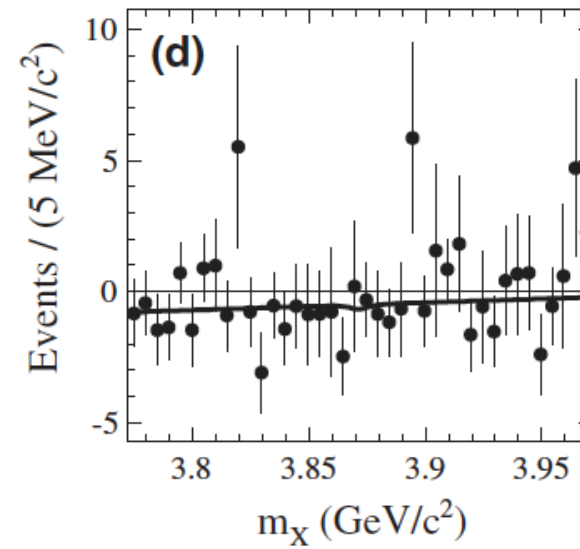
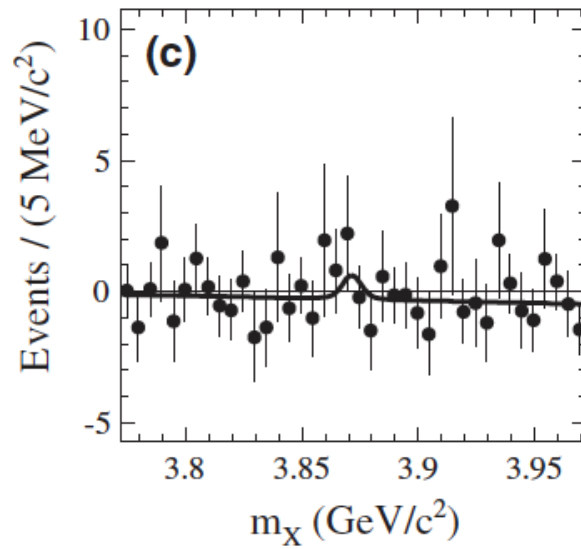
# BaBar's results ( $\psi'\gamma$ )

$B^{\pm} \rightarrow X K^{\pm}$



$K^0$

$K^{*}_{\pm}$



$K^{*0}$

# Belle's results ( $J/\psi\gamma$ )

$$B^\pm \rightarrow X K^\pm$$

$$B^0 \rightarrow X K_S^0$$

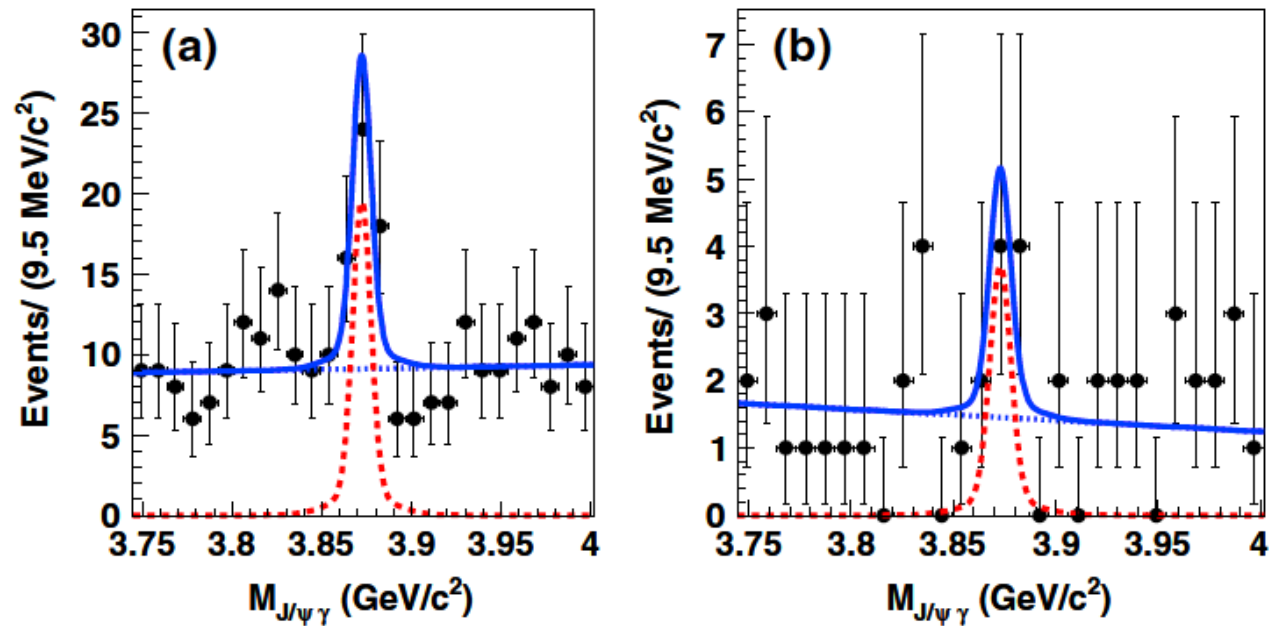


FIG. 2 (color online).  $M_{J/\psi\gamma}$  distributions for (a)  $B^+ \rightarrow X(3872)(\rightarrow J/\psi\gamma)K^+$  and (b)  $B^0 \rightarrow X(3872)(\rightarrow J/\psi\gamma)K_S^0$  decays. The curves show the signal (red dashed) and the background component (blue dotted) as well as the overall fit (blue solid).

# Belle's results ( $\psi'\gamma$ )

$$B^\pm \rightarrow X K^\pm$$

$$B^0 \rightarrow X K_S^0$$

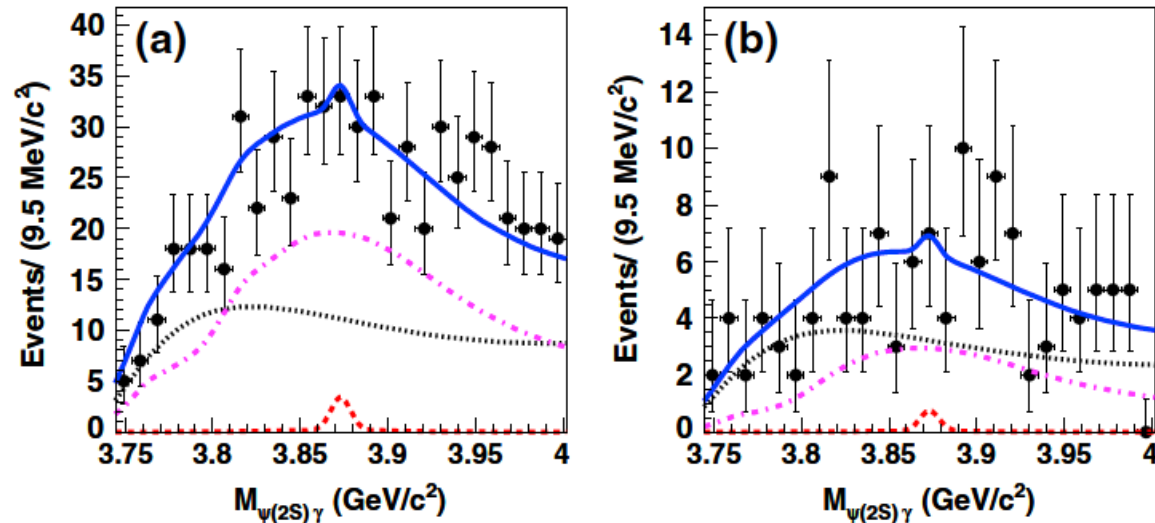
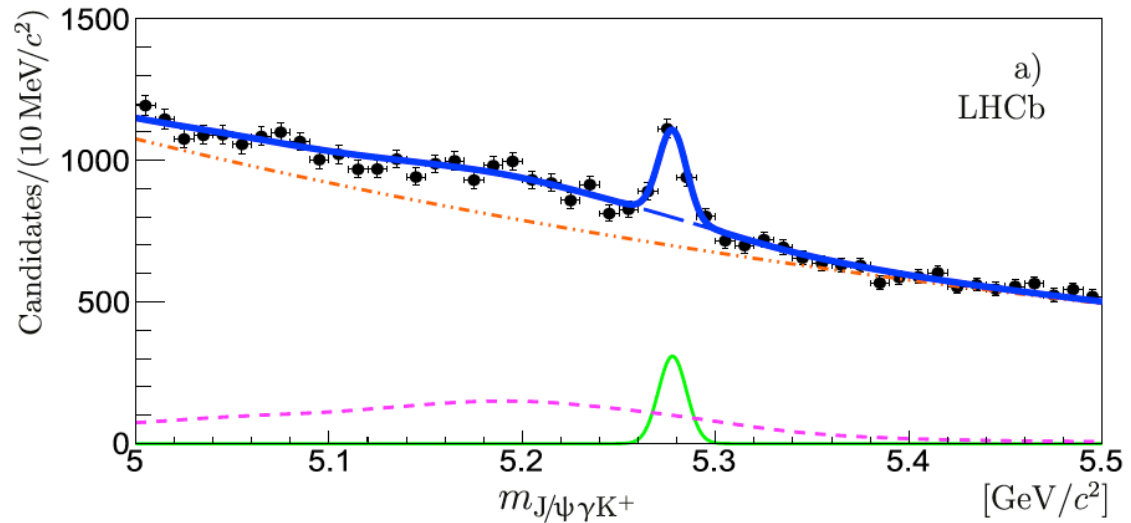


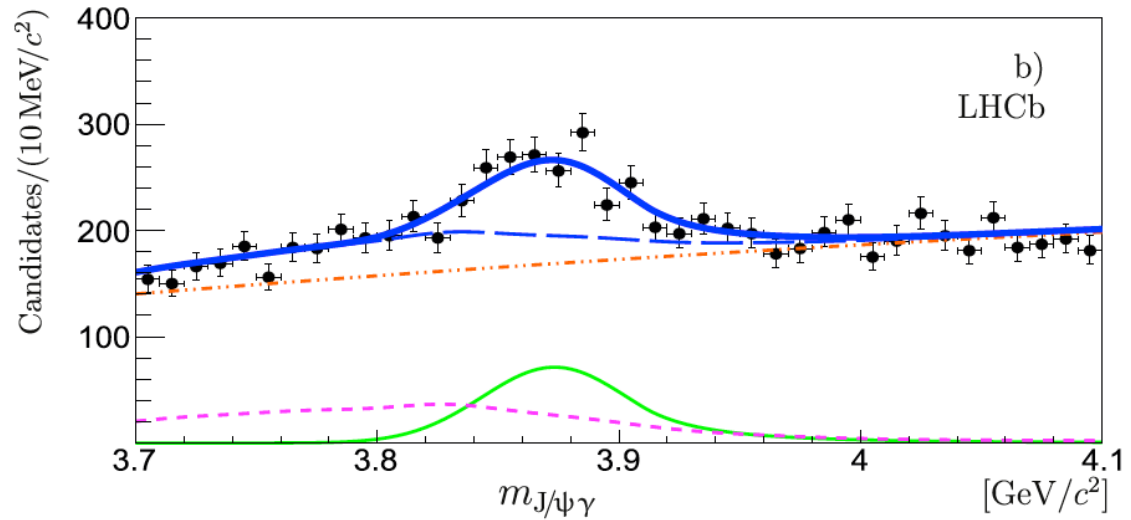
FIG. 3 (color online).  $M_{\psi'\gamma}$  distributions for (a)  $B^+ \rightarrow X(3872)(\rightarrow \psi'\gamma)K^+$  and (b)  $B^0 \rightarrow X(3872)(\rightarrow \psi'\gamma)K^0$ . The curves show the signal [red dashed for  $X(3872)$ ] and the background component [pink dot-dashed for background from  $B \rightarrow \psi'K^*$  and  $B \rightarrow \psi'K$  component, and black dotted for combinatorial background modeled by the threshold function] as well as the overall fit (blue solid).

# LHCb's results ( $J/\psi\gamma$ )

Invariant mass  
of  $J/\psi\gamma K^+$

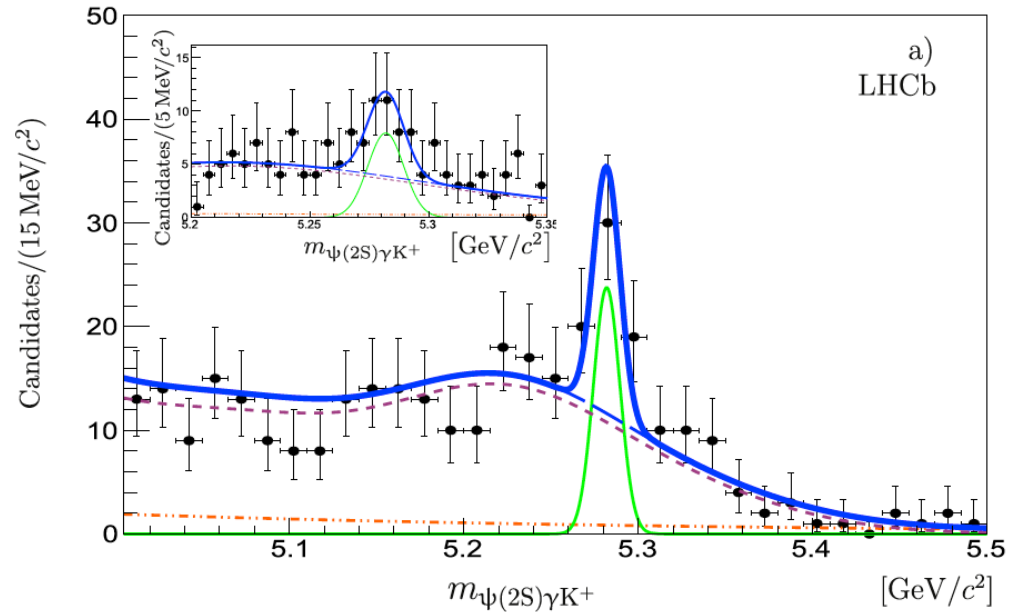


Invariant mass  
of  $J/\psi\gamma$

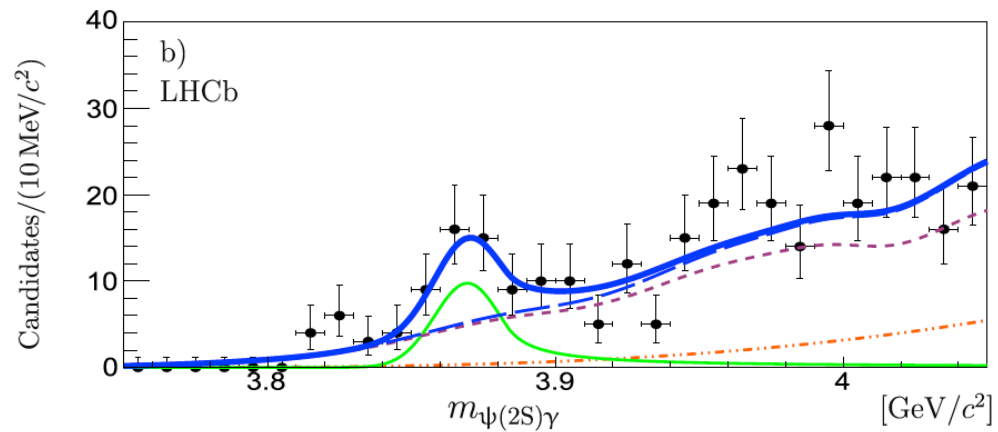


# LHCb's results ( $\psi'\gamma$ )

Invariant mass  
of  $\psi(2S)\gamma K^+$



Invariant mass  
of  $\psi(2S)\gamma$



# Wavefunction of X(3872)

- Probability of the each component of X(3872)

Our numerical results:

Model	$CC^{\text{bar}}$	$D^0 D^{*0\text{bar}}$	$D^+ D^{*-}$	$J/\psi\rho$	$J/\psi\omega$
Case A	0.036	0.913	0.034	0.006	0.010
Case QM	0.061	0.864	0.019	0.007	0.019

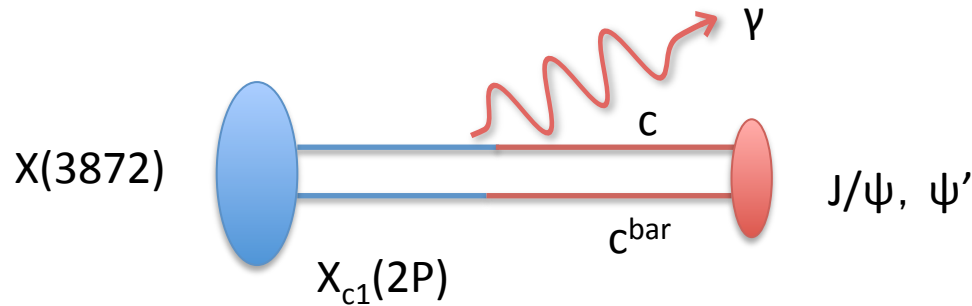
$\rho$  and  $\omega$  widths are not introduced for the radiative decay calculations.

# Wavefunctions of Charmonia

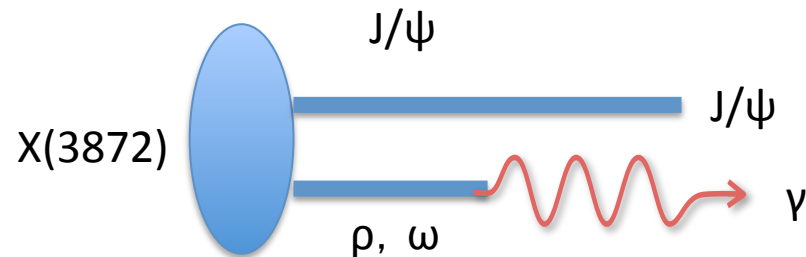
- Charmonia:  $\chi_{c1}(2P)$ ,  $\psi(2S)$ ,  $J/\psi$
- As for the wavefunction of charmonium we used the potential model of S. Godfrey and N. Isgur, PRD 32, 189 (1985) (without the LS and tensor terms for simplicity)
- $\chi_{c1}(2P)$  in the X(3872) is offshell, but we used the onshell wavefunction.

# Radiative decays of X(3872)

- E1 transition of  $cc^{\text{bar}}$  core to  $J/\psi(\psi') \gamma$



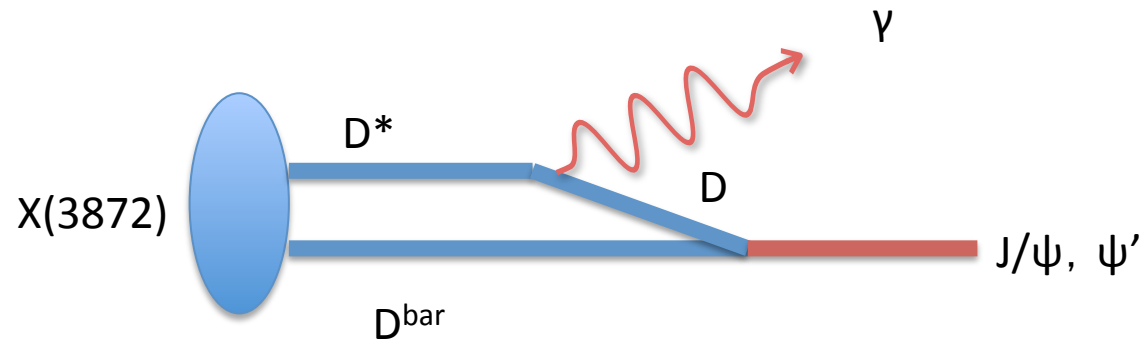
- Vector meson dominance





# Radiative decays of X(3872)

- D meson loop (not considered here)



- Partial decay width  $D^{*+} \rightarrow D^+ \gamma = 1.3 \text{ keV}$
- Overlap of  $DD^{\text{bar}} - J/\psi$  is small

# E1 transitions of $c\bar{c}$ core

$$\Gamma(X(3872) \rightarrow J/\psi(\psi') + \gamma) = \left(\frac{2}{3}\right)^4 \alpha \left| \langle \psi_f | r | \psi_i \rangle \right|^2 E_\gamma^3 \frac{E_f^{c\bar{c}}}{M_i^{c\bar{c}}}$$

$$\langle J/\psi | r | \chi_{c1}(2P) \rangle = 0.0396 \text{ fm}$$

$$\langle \psi' | r | \chi_{c1}(2P) \rangle = 0.5225 \text{ fm}$$

$$\langle J/\psi | r | \chi_{c1}(2P) \rangle = 0$$

for harmonic oscillator wavefunction

# Radiative decay widths ( $cc^{\text{bar}}$ core contribution)

Numerical result: Case A

$$\Gamma(X(3872) \rightarrow J/\psi + \gamma) = 0.7 \text{ keV}$$

$$\Gamma(X(3872) \rightarrow \psi' + \gamma) = 2.6 \text{ keV}$$

$$R = \frac{\Gamma(X(3872) \rightarrow \psi' + \gamma)}{\Gamma(X(3872) \rightarrow J/\psi + \gamma)} = 3.6$$

# Radiative decay widths ( $c\bar{c}$ core contribution)

Numerical result: Case QM

$$\Gamma(X(3872) \rightarrow J/\psi + \gamma) = 1.0 \text{ keV}$$

$$\Gamma(X(3872) \rightarrow \psi' + \gamma) = 3.5 \text{ keV}$$

$$R = \frac{\Gamma(X(3872) \rightarrow \psi' + \gamma)}{\Gamma(X(3872) \rightarrow J/\psi + \gamma)} = 3.5$$

# VMD contribution

$$\Gamma_{VMD} = \frac{4}{27} \alpha \frac{q E_{J/\psi}}{m_X} |\psi_\omega(r=0)|^2 \left( Z_{\omega J/\psi} \phi_{\omega J/\psi}(q) + 3Z_{\rho J/\psi} \phi_{\rho J/\psi}(q) \right)^2$$

$q$ : photon momentum

$Z_\alpha^2$ : probability of finding  $\alpha$  component  
in  $X(3872)$

Numerical result:

Case A

$$\Gamma_{VMD}(X(3872) \rightarrow J/\psi + \gamma) = 2.6 \text{ keV}$$

Case QM

$$\Gamma_{VMD}(X(3872) \rightarrow J/\psi + \gamma) = 6.2 \text{ keV}$$

# Discussions

- Since we cannot determine the relative sign of the  $cc^{\text{bar}}$  core contribution and VMD contribution, rough estimation is:

$$\Gamma(X(3872) \rightarrow J / \psi + \gamma)$$

$$= \left( \sqrt{\Gamma_{VMD}} - \sqrt{\Gamma_{c\bar{c}}} \right)^2 \sim \left( \sqrt{\Gamma_{VMD}} + \sqrt{\Gamma_{c\bar{c}}} \right)^2$$

$$= 0.6 \text{ keV} \sim 6.0 \text{ keV} \quad (\text{Case A})$$

$$= 1.6 \text{ keV} \sim 12.2 \text{ keV} \quad (\text{Case QM})$$

$$R = \frac{\Gamma(X(3872) \rightarrow \psi' + \gamma)}{\Gamma(X(3872) \rightarrow J/\psi + \gamma)}$$

$$= 0.43 \sim 4.3 \quad (\text{Case A})$$

$$= 0.29 \sim 1.18 \quad (\text{Case QM})$$

$3.4 \pm 1.4$  [BaBar]       $< 2.1$  (90%CL) [Belle]

$2.46 \pm 0.64(\text{stat}) \pm 0.29(\text{sys})$  [LHCb]

$$\Psi(4040, 4415) \rightarrow X(3872) + \gamma$$

- BESIII observed the radiative decay  
 $Y(4260) \rightarrow X(3872) + \gamma$   
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$$\frac{B[Y(4260) \rightarrow X(3872) \gamma]}{B[Y(4260) \rightarrow J/\psi \pi^+ \pi^-]} = 0.1$$

- $Y(4260)$  may be not a simple charmonium and the structure of  $Y(4260)$  is unclear.



$$\Psi(4040, 4415) \rightarrow X(3872) + \gamma$$

- We consider well established  $\psi(3S, 4040)$  and  $\psi(4S, 4415)$  instead of  $\Upsilon(4260)$ .
- We consider E1 transition of  $cc^{\text{bar}} 3S_1(n=3,4) \rightarrow cc^{\text{bar}}$  core of  $X(3872) + \gamma$ .

Decay width	$\Psi(4040) \rightarrow X(3872) \gamma$	$\Psi(4415) \rightarrow X(3872) \gamma$
Case A	2.4 keV	0.16 keV
Case QM	4.0 keV	0.27 keV

# Summary

- Estimation of  $\Gamma(X(3872) \rightarrow J/\psi + \gamma)$  is rather difficult.
- $\Gamma(X(3872) \rightarrow \psi' + \gamma)$  is proportional to size of  $cc^{\text{bar}}$  core component in  $X(3872)$ .
- The measurement of absolute value of  $\Gamma(X(3872) \rightarrow \psi' + \gamma)$  is very important to determine the nature of  $X(3872)$ .

# Summary

- We have a chance to observe  $\psi(4040) \rightarrow X(3872) \gamma$  decay.
- $\Gamma(Y(4260) \rightarrow X + \gamma) < \Gamma(\psi(4040) \rightarrow X + \gamma)$   
since  $Y(4260)$  may have the non- $cc^{\text{bar}}$  component.