## Remarks on $\pi^- p \rightarrow l^- l^+ n$

P. Kroll

Fachbereich Physik, Univ. Wuppertal and Univ. Regensburg KEK, Tsukuba, March 2015



the exclusive limit of the Drell-Yan process directly related to leptoproduction of pions - same GPDs

-  $\hat{s}-\hat{u}$  crossed subprocess

$$\mathcal{H}^{\pi^- \to \gamma^*}(\hat{u}, \hat{s}) = -\mathcal{H}^{\gamma^* \to \pi^+}(\hat{s}, \hat{u})$$

or for hard scattering kernel  $\mathcal{F}^{t.l.*}(x,\xi,\mathbf{k}_{\perp}^2,\bar{z}) = \mathcal{F}^{s.l.}(x,\xi,-\mathbf{k}_{\perp}^2,z)$ equivalent to  $Q^2 \to -Q'^2$  Berger-Diehl-Pire (01): leading-twist, LO analysis of long. cross( i.e. exploiting asymp. factorization formula)

we know that leading-twist analysis of  $\pi^+$  production fails with HERMES data by order of magnitude

Therefore a reanalyis of the exclusive Drell-Yan process seems appropriate making use of what we have learned from analysis of pion production

- retaining quark transverse momenta in the subprocess (the MPA)
- treating pion-pole contribution as an OPE term
- take into account transverse photons and transversity GPDs

## What do we know about the GPDs?

- $\widetilde{H}$ : known from form factor analysis Diehl-K(13) and probed in  $\pi$  leptonproduction, spin asymmetries in DVCS, SDME in DVMP
- $\widetilde{E}$ : contribution from pion pole, no clear signal for a non-pole contribution
- $H_T$ : not well-known, information from  $\pi^+$  leptoproduction, spin asymmetries in vector-meson leptoproduction, moments from lattice QCD
- $\bar{E}_T$ : not well-known, information from SDMEs in vector-meson leptoproduction, moments from lattice QCD

## **Cross section**



k momentum of  $l^ \tau = Q'^2/(s-m^2)$ 

$$\frac{d\sigma}{dtdQ'^2d\cos\theta d\phi} = \frac{3}{8\pi} \left\{ \sin^2\theta \frac{d\sigma_L}{dtdQ'^2} + \frac{1+\cos^2\theta}{2} \frac{d\sigma_T}{dtdQ'^2} + \frac{1}{\sqrt{2}} \sin\left(2\theta\right)\cos\phi \frac{d\sigma_{LT}}{dtdQ'^2} + \sin^2\theta\cos\left(2\phi\right)\frac{d\sigma_{TT}}{dtdQ'^2} \right\}$$

PK 4

$$\frac{d\sigma_L}{dt dQ'^2} = \frac{\alpha_{\rm elm}}{48\pi^2} \frac{\tau^2}{Q'^6} \sum_{\nu'} |\mathcal{M}_{0\nu',0+}|^2 \qquad \frac{d\sigma_T}{dt dQ'^2} = \frac{\alpha_{\rm elm}}{48\pi^2} \frac{\tau^2}{Q'^6} \sum_{\mu=\pm 1,\nu'} |\mathcal{M}_{\mu\nu',0+}|^2$$
$$\frac{d\sigma_{LT}}{dt dQ'^2} = \frac{\alpha_{\rm elm}}{48\pi^2} \frac{\tau^2}{Q'^2} \operatorname{Re} \sum_{\nu'} \mathcal{M}_{0\nu'0+}^* (\mathcal{M}_{+\nu'0+} - \mathcal{M}_{-\nu'0+})$$
$$\frac{d\sigma_{TT}}{dt dQ'^2} = \frac{\alpha_{\rm elm}}{48\pi^2} \frac{\tau^2}{Q'^2} \operatorname{Re} \sum_{\nu'} \mathcal{M}_{+\nu'0+}^* \mathcal{M}_{-\nu'0+}$$

 $\phi$  integration

$$\frac{d\sigma}{dtdQ'^2d\cos\theta} = \frac{3}{4}\sin^2\theta\frac{d\sigma_L}{dtdQ'^2} + \frac{3}{8}(1+\cos^2\theta)\frac{d\sigma_T}{dtdQ'^2}$$

 $\boldsymbol{\theta}$  integration

$$\frac{d\sigma}{dtdQ'^2} = \frac{d\sigma_T}{dtdQ'^2} + \frac{d\sigma_T}{dtdQ'^2}$$

What do we learn on the GPDs (or convolutions) from the Drell-Yan process:

- $d\sigma_L$ :  $\langle \widetilde{H} \rangle$  and  $\langle \widetilde{E} \rangle$  (essentially the pion pole)
- $d\sigma_T$ :  $\langle H_T \rangle$  ( $\langle \bar{E}_T \rangle$  small)
- $\sin 2\theta \cos \phi$ :  $\operatorname{Re}\left(\langle \widetilde{E} \rangle^* \langle H_T \rangle\right)$
- $\sin^2 \theta \cos 2\phi$ :  $|\langle \bar{E}_T \rangle|^2$  (small)

can exclusive Drell-Yan help to fix these GPDs better? measurements at which values of  $s, Q'^2, t$ ?

**Polarization?**